

## lecture-06

Devisson & Germer Experiment to prove wave particle properties:

To discuss critical results of wave particle duality, a new dynamical system is formed.

Wave particle duality was first proved by De Broglie in 1924. In 1927, Devisson & Germer proved wave-particle properties experimentally. They passed a beam of electrons through a metal crystal grating.

Under certain conditions, beam of light travels in form of waves & under certain conditions, it travel like photons as particles.

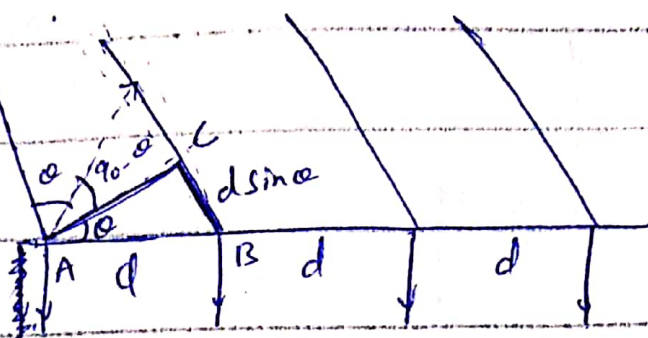
By considering light as wave, wave mechanics was established

of light consists of waves, it should satisfy two wave properties, diffraction, interference



Beam waves have wave length  $\lambda$  which is also property of waves.

fig(a)



We have considered a nikel crystal grating. Beam of electrons is incident on plate at perpendicular.

Then waves of Beam are diffracted at an angle ' $\theta$ '. Distance b/w two nikel atoms is separated by ' $d$ '. In right angled triangle  $\Delta ABC$ , we have

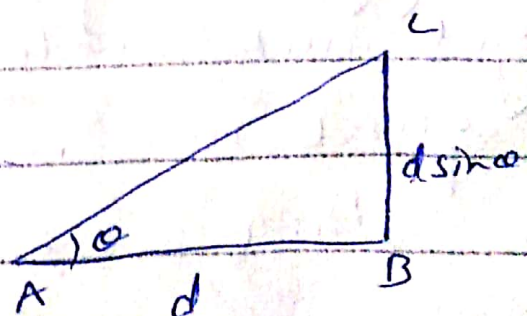
$$BC = d \sin \theta$$

so electrons beam

consists of particles

but obey wave property,

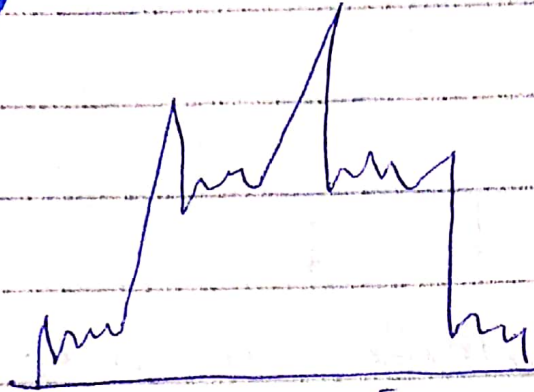
which is diffraction. From where





We can say that interference property also holds to prove wave particle dual nature.

Fig (b)



a typical diffraction pattern, showing maxima & min arising from constructive or destructive interference.

**Bragg's Diffraction Peak:-**

These peaks were observed as Bragg's diffraction peaks. We can get max intensity if  $BC \perp$  equal to integral multiple of  $\lambda$ .

$$d \sin \theta = n\lambda, \text{ where}$$

$$d = \frac{h}{p}$$

(It's condition on  $\theta$  for constructive interference to be at its strongest)



## G.P Thompson Experiment:-

Similar diffraction patterns were obtained by G.P Thompson when positively charged particles (proton, oxygen) were incident on thin metal.

## Summerfield-Whitson Principle:-

Due to this principle, we have to quantize a particle. This is generalization of third law of Bohr. This rule is also applied for non-circular (elliptic) orbits.

### Statement:-

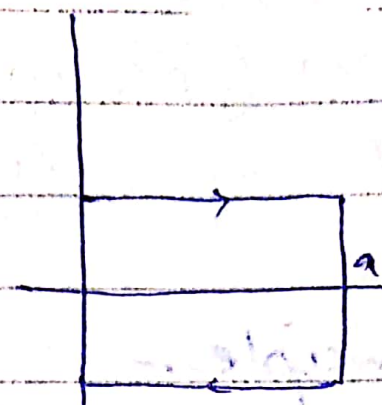
Let  $P_i$  denotes momenta of particles  $i = (1, 2, \dots, N)$ . Also consider  $q_i$ , their coordinates, this rule states,

$\oint P_i dq_i = n_i h$ , where  $\oint$  denotes a closed cycle. we have to prove that:-

$$E = \frac{n^2 h^2}{8m a^2}$$



To prove it, we consider a box of width  $a$ .



$$2\pi m v \lambda = nh$$

$$m v \lambda = \frac{nh}{2\pi} = n \frac{h}{2\pi}$$

In this box, a particle is moving freely in one dimension. In this box, particle moves back & forth motion. First we write formula of total energy where  $P.E = 0$  (particle is moving freely).

$$\text{Total energy} = E = K.E + P.E = \frac{1}{2} m v^2$$

$$E = \frac{m^2 v^2}{2m} = \frac{p^2}{2m}$$

Now, we apply de Broglie's principle; as

$$2 \int_0^a p dq = nh$$

$$2 p a = nh \Rightarrow p = \frac{nh}{2a}$$

Now formula of energy; as

$$E = \frac{1}{2m} \left( \frac{nh}{2a} \right)^2 = \frac{n^2 h^2}{8ma^2}$$



## Sh Schrodinger equation :-

we will consider wave motion of a single particle. let  $\psi(x, t)$  be a wave f. we know that wave equation is

$$\psi_{xx} - c^2 \psi_{tt} = 0. \text{ It is clear that}$$

its solution is

$$\psi(x, t) = e^{i(kx - \omega t)} \quad (\text{justify sol})$$

In next theory, we will take

frequency as  $\nu$ .  $\omega$

$$\omega = 2\pi\nu, \quad E = h\nu = \frac{h}{2\pi} \cdot 2\pi\nu = \hbar\omega$$

$$\omega \quad \text{and} \quad k = \frac{2\pi}{\lambda} = \frac{2\pi\nu}{c} = \frac{\omega}{c}$$

$$\Rightarrow \omega = kc \quad \text{or} \quad k^2 c^2 = \omega^2$$

we form wave equation of a packet when its localized (waves are present in specific space in specific time).

By inverse Fourier we get

$$\psi(x, t) = \int e^{i(kx - \omega t)} f(k) d^3k \quad \text{--- (1)}$$

Differentiating it w.r.t 't' twice  
Time in vector form :-



$$\nabla^2 \psi(x, t) = -k^2 \int e^{i(kx - \omega t)} f(k) d^3k$$

multiplying by  $-\frac{\hbar^2}{2m}$  in b.s; as

$$-\frac{\hbar^2}{2m} \nabla^2 \psi(x, t) = \frac{\hbar^2 k^2}{2m} \int e^{i(kx - \omega t)} f(k) d^3k$$

$$-\frac{\hbar^2}{2m} \nabla^2 \psi(x, t) = \hbar \omega \int e^{i(kx - \omega t)} f(k) d^3k \quad (2)$$

F

$$\left\{ \begin{array}{l} E = \frac{p^2}{2m} = \frac{\hbar^2 k^2}{2m} \\ p = \frac{h}{\lambda} = \frac{h}{2\pi} \frac{2\pi}{\lambda} = \hbar k \quad \text{or} \quad E = \hbar \omega \quad \text{so} \\ \hbar \omega = \frac{\hbar^2 k^2}{2m} \end{array} \right.$$

Differentiating (1) w.r.t (t); as

$$\frac{\partial \psi(x, t)}{\partial t} = -i\omega \int e^{i(kx - \omega t)} f(k) d^3k$$

Multiplying by  $i\hbar$  in b.s; as

$$i\hbar \frac{\partial \psi}{\partial t} = \hbar \omega \int e^{i(kx - \omega t)} f(k) d^3k \quad (3)$$

Comparing (2) & (3); we get

$$-\frac{\hbar^2}{2m} \nabla^2 \psi = i\hbar \frac{\partial \psi}{\partial t}$$

which is Schrodinger equation.