

Example A string is stretched and fixed between two points $(0,0)$ and $(l,0)$. Motion is initiated by displacing the string in the form $u = \lambda \sin(\pi x/l)$ and released from rest at time $t=0$. Find the displacement of any point on the string at any time t .

Soln The displacement $u(x,t)$ is governed by $u_{tt} = c^2 u_{xx}$, $0 < x < l$, $t > 0$,

BCs: $u(0,t) = u(l,t) = 0$ ————— (1)

ICs: $u(x,0) = \lambda \sin \frac{\pi x}{l}$, $u_t(x,0) = 0$.

A Taking Laplace transform, we get

$$c^2 \frac{d^2 \bar{u}}{dx^2} = s^2 \bar{u}(x,s) - s u(x,0) - u_t(x,0)$$

$$\boxed{\frac{d^2 \bar{u}}{dx^2} - \frac{s^2}{c^2} \bar{u} = -\frac{\lambda s}{c^2} \sin \frac{\pi x}{l}}$$

$$\frac{1}{s^2 + \frac{\pi^2 c^2}{l^2}}$$

Its general soln is

$$\bar{u}(x,s) = A e^{\frac{s}{c}x} + B e^{-\frac{s}{c}x} + \frac{\lambda s \sin(\frac{\pi}{l})x}{s^2 + \pi^2 c^2 / l^2}$$

Laplace transform of (*) yields

$$\bar{u}(0,s) = 0, \bar{u}(l,s) = 0 \text{ ————— (2)}$$

in (1) in (2)

$$A + B = 0 \quad \text{--- (3)}$$

$$A e^{s\ell/c} + B e^{-s\ell/c} = 0 \quad \text{--- (4)}$$

Solving (3) and (4) we have $A = B = 0$.

Thus from (1) $\bar{u}(x, s) = \frac{\lambda s \sin(\pi x/\ell)}{s^2 + \pi^2 c^2/\ell^2}$.

Laplace inversion yields.

$$u(x, t) = \mathcal{L}^{-1}[\bar{u}(x, s)]$$

$$= \lambda \mathcal{L}^{-1}\left[\frac{s}{s^2 + \frac{\pi^2 c^2}{\ell^2}}\right] \sin \frac{\pi x}{\ell}$$

$$= \lambda \cos\left(\frac{\pi c}{\ell} t\right) \sin \frac{\pi x}{\ell}$$

which gives the displacement of the string for a given time.

Example An infinitely long string having one end at $x=0$ is initially at rest on the x -axis. The end $x=0$ undergoes a periodic transverse displacement described by $A_0 \sin \omega t, t > 0$. Find the displacement of any point on the string at any time t .

Soln Let $u(x, t)$ be the transverse displacement of the string at any point x and at any time t . Then the transverse displacement of the string is described by the

PDE: $u_{tt} = c^2 u_{xx}, \quad x > 0, \quad t > 0 \quad \text{--- (1)}$

(1) BCS: $u(0, t) = A_0 \sin \omega t, t > 0$ ($A_0 = \text{constant}$) (2)

ICS: $u(x, 0) = 0, u_t(x, 0) = 0, x > 0.$

The Laplace transform of (1) is

$$c^2 \frac{d^2 \bar{u}}{dx^2} = s^2 \bar{u} - s u(x, 0) - u_t(x, 0)$$

using initial conditions we have.

$$\frac{d^2 \bar{u}}{dx^2} - \frac{s^2}{c^2} \bar{u}(x, s) = 0 \quad (3)$$

Laplace transform of (2) is

$$\bar{u}(0, s) = \frac{A_0 \omega}{s^2 + \omega^2} \quad (4)$$

$\frac{f(x) = A_0 \sin \omega t}{\mathcal{L}\{f(x)\} = \frac{A_0 \omega}{s^2 + \omega^2}}$

The general solution of (3) is

$$\bar{u}(x, s) = A e^{(s/c)x} + B e^{-(s/c)x} \quad (5)$$

Since the displacement $\bar{u}(x, s)$ is bounded as $x \rightarrow \infty$, we get $A = 0$ and, therefore, from

$$(5) \quad \bar{u}(x, s) = B e^{-(s/c)x} \quad (6)$$

Using (4) in (6) we obtain

$$B = \bar{u}(0, s) = \frac{A_0 \omega}{s^2 + \omega^2} \quad (7)$$

(7) in (6) yields

$$\bar{u}(x, s) = \frac{A_0 \omega}{s^2 + \omega^2} e^{-(s/c)x}$$

Taking inverse Laplace we get

$$u(x, t) = \mathcal{L}^{-1}[\bar{u}(x, s)] = \mathcal{L}^{-1}\left[\frac{A_0 \omega}{s^2 + \omega^2} e^{-(s/c)x}\right]$$

$$= \begin{cases} A_0 \sin \omega(t - x/c) & \text{if } t > x/c \\ 0 & \text{if } t < x/c \end{cases}$$

plc: Find the soln of

$$u_t = k u_{xx}, \quad 0 < x < \infty, \quad t > 0$$

$$u(x, t) \rightarrow 0 \text{ as } x \rightarrow \infty$$

$$u(0, t) = f(t)$$

$$u(x, 0) = 0$$

assuming $L^{-1} \left[\exp\left(-\sqrt{\frac{s}{k}} x\right) \right] = \frac{x}{2\sqrt{k\pi t^3}} \exp\left(-\frac{x^2}{4kt}\right)$.

Taking Laplace transform, we have

$$k \frac{d^2 \bar{u}}{dx^2} = s \bar{u}(x, s) - u(x, 0)$$

$$u(x, 0) = 0 \quad \therefore$$

$$\boxed{\frac{d^2 \bar{u}}{dx^2} - \frac{s}{k} \bar{u} = 0}$$

Soln is

$$\bar{u}(x, s) = A \exp\left(\sqrt{\frac{s}{k}} x\right) + B \exp\left(-\sqrt{\frac{s}{k}} x\right) \quad \text{--- (1)}$$

The Laplace transform of the BCS give

$$\bar{u}(x, s) \rightarrow 0 \text{ as } x \rightarrow \infty \quad \text{--- (2)}$$

$$\bar{u}(0, s) = \bar{f}(s) \quad \text{--- (3)}$$

Using (2) in (1) we get $A = 0$

$$\therefore \bar{u}(x, s) = B \exp\left(-\sqrt{\frac{s}{k}} x\right) \quad \text{--- (4)}$$

From (3) & (4)

$$B = \bar{f}(s)$$

$$\therefore \boxed{\bar{u}(x, s) = \bar{f}(s) \exp\left(-\sqrt{\frac{s}{k}} x\right)} \quad \text{--- (5)}$$

Laplace inversion of (5) yields

$$u(x,t) = \mathcal{L}^{-1} \left[\bar{g}(s) e^{-\frac{\sqrt{s} x}{\sqrt{k}}} \right]$$

$$= \mathcal{L}^{-1} \left[\left[\mathcal{L}(g(t)) \right] \mathcal{L} \left[\frac{x}{2\sqrt{k\pi t^3}} e^{-\frac{x^2}{4kt}} \right] \right]$$

By convolution theorem we have

$$u(x,t) = \int_0^t \frac{x \exp \left[-\frac{x^2}{4k(t-u)} \right]}{2\sqrt{\pi k} (t-u)^{3/2}} g(u) du$$

Example Solve $u_t(x,t) = u_{xx}(x,t) - hu(x,t)$, $h = \text{constant}$, $0 < x < \pi$, $t > 0$

Bcs: $u(0,t) = 0, t > 0$, $u(\pi,t) = 1, t > 0$ } (2)

IC: $u(x,0) = 0, t = 0$. (3)

Taking Laplace transform of (1) we get

$$\frac{d^2 \bar{u}}{dx^2} - (h+s) \bar{u}(x,s) = 0 \quad \left[\begin{array}{l} \text{As } \bar{u}(x,0) = u(x,0) \\ = \frac{d^2 \bar{u}}{dx^2} - h\bar{u} \end{array} \right]$$

The general soln is

$$\bar{u}(x,s) = A \cosh(\sqrt{h+s} x) + B \sinh(\sqrt{h+s} x) \quad (4)$$

Taking Laplace transform of (2) we have

$$\bar{u}(0,s) = 0, \bar{u}(\pi,s) = \frac{1}{s} \quad (5)$$

$\bar{u}(0,s) = 0 \Rightarrow A = 0$ from (4) Hence

$$\bar{u}(x,s) = B \sinh(\sqrt{h+s} x) \quad (6)$$

$$f(s) = \frac{1}{s} = B \operatorname{sinh}(\sqrt{h+s} \pi)$$

$$\Rightarrow B = \frac{1}{s \operatorname{sinh}(\sqrt{h+s} \pi)} \quad \text{--- (7)}$$

Using (7) in (6) we get:

$$\bar{u}(x, s) = \frac{\operatorname{sinh}(\sqrt{h+s} x)}{s \operatorname{sinh}(\sqrt{h+s} \pi)}$$

By means of complex inversion formula

$$u(x, t) = \frac{1}{2\pi i} \int_{\gamma-i\infty}^{\gamma+i\infty} \frac{e^{st} \operatorname{sinh}(\sqrt{h+s} x)}{s \operatorname{sinh}(\sqrt{h+s} \pi)} \cdot L \rightarrow \text{(8)}$$

= sum of all the contributions from all the poles of the integrand.

From (8), the poles are given by $s=0$, and $\operatorname{sinh}(\sqrt{h+s} \pi) = 0 = \sin(i \sqrt{h+s} \pi)$. Therefore,

$$i \sqrt{h+s} \pi = n\pi, \text{ implying } h+s = -n^2.$$

Thus, the integrand has poles at $s=0$, and $s = -n^2 - h, n=1, 2, 3, \dots$

The residue of the expression $\frac{e^{st} \operatorname{sinh}(\sqrt{h+s} x)}{s \operatorname{sinh}(\sqrt{h+s} \pi)}$ at $s=0$ is $\frac{\operatorname{sinh}(\sqrt{h} x)}{\operatorname{sinh}(\sqrt{h} \pi)}$ --- (9)

The residue of the integrand at $s = -n^2 - h$ is $\frac{e^{-(n^2+h)t} \operatorname{sinh}(\sqrt{-n^2} x)}{\operatorname{sinh} \sqrt{-n^2} \pi + \frac{(-n^2-1)\pi}{2\sqrt{-n^2}} \operatorname{cosh}(\sqrt{-n^2} \pi)}, n=1, 2, 3, \dots$

Using the relations $\operatorname{sinh} x = i \sin ix$, and $\operatorname{cosh} x = \cos ix$, the above residues become

$$\frac{[-(n^2+h)t] i \sin nx}{-(n^2+h)\pi \cos n\pi} = \frac{2n \exp[-(n^2+h)t] \sin nx}{(n^2+h)\pi \cos n\pi}$$

$$= \frac{2n \exp[-(n^2+h)t] \sin nx}{(n^2+h)\pi (-1)^n} \quad \text{--- (10)}$$

Hence from (8) --- (10) we get

$$u(x,t) = \frac{\sinh(\sqrt{h}x)}{\sinh(\sqrt{h}\pi)} + \frac{2e^{-ht}}{\pi} \sum_{n=1}^{\infty} \frac{n(-1)^n \exp(-n^2 t) \sin nx}{(n^2+h)}$$

Ex Find the inverse Laplace transform of

$$\frac{e^{-1/s}}{\sqrt{s}}$$

$$\frac{1}{\sqrt{s}} e^{-1/s} = \frac{1}{\sqrt{s}} \sum_{n=0}^{\infty} \frac{(-1)^n}{n! s^n} = \sum_{n=0}^{\infty} \frac{(-1)^n}{n! s^{n+1/2}}$$

$$\mathcal{L}^{-1} \left[\frac{e^{-1/s}}{\sqrt{s}} \right] = \sum_{n=0}^{\infty} \frac{(-1)^n}{n!} \mathcal{L}^{-1} \left[\frac{1}{s^{n+1/2}} \right] = \sum_{n=0}^{\infty} \frac{(-1)^n}{n!} \frac{t^{n-1/2}}{\Gamma(n+1/2)} \quad \text{--- (1)}$$

$$\Gamma(n+1/2) = \frac{(2n)!}{2^{2n} n!} \sqrt{\pi}$$

using in (1) we get

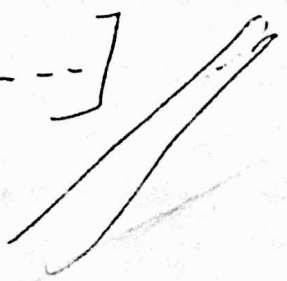
$$\mathcal{L}^{-1} \left[\frac{e^{-1/s}}{\sqrt{s}} \right] = \sum_{n=0}^{\infty} \frac{(-1)^n}{(2n)! \sqrt{\pi}} \frac{2^{2n} t^{n-1/2}}{2^{2n} n!}$$

$$= \frac{t^{-1/2}}{\sqrt{\pi}} - \frac{4 t^{1/2}}{2! \sqrt{\pi}} + \frac{4 t^{3/2}}{4! \sqrt{\pi}} - \dots$$

$$= \frac{1}{\sqrt{\pi t}} \left[1 - \frac{(2\sqrt{t})^2}{2!} + \frac{(2\sqrt{t})^4}{4!} - \dots \right]$$

$$= \frac{\cos 2\sqrt{t}}{\sqrt{\pi t}}$$

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Ex Find the Laplace transform of $\text{erf}(1/\sqrt{t})$.

$$L[\text{erf}(1/\sqrt{t})] = \int_0^{\infty} e^{-st} \text{erf}(1/\sqrt{t}) dt$$

$$= \int_0^{\infty} e^{-st} \left[\frac{2}{\sqrt{\pi}} \int_0^{1/\sqrt{t}} e^{-u^2} du \right] dt$$

$$= \frac{2}{\sqrt{\pi}} \int_0^{\infty} \int_0^{1/\sqrt{t}} e^{-st-u^2} du dt. \quad \text{--- (1)}$$

Now, changing the order of integration, we get.

$$L[\text{erf}(1/\sqrt{t})] = \frac{2}{\sqrt{\pi}} \int_0^{\infty} e^{-u^2} du \int_0^{1/u^2} e^{-st} dt.$$

$$= \frac{2}{\sqrt{\pi}} \int_0^{\infty} e^{-u^2} du \left[\frac{e^{-st}}{-s} \right]_0^{1/u^2}$$

$$= -\frac{2}{\sqrt{\pi}} \int_0^{\infty} \left[e^{-u^2} - \frac{e^{-s/u^2}}{-s} \right] du \quad \text{--- (2)}$$

Consider

$$I(a, b) = \int_0^{\infty} e^{-a^2 u^2 - b^2/u^2} du \quad \text{--- (2a)}$$

$$I(a, 0) = \int_0^{\infty} e^{-a^2 u^2} du = \frac{\sqrt{\pi}}{2a} \quad \text{--- (3)}$$

$$\left(\text{since } \frac{2}{\sqrt{\pi}} \int_0^{\infty} e^{-u^2} du = 1 \right)$$

Differentiating (2a) w.r.t. b partially, we have

$$\frac{\partial I}{\partial b} = -2b \int_0^{\infty} u^{-2} e^{-a^2 u^2 - b^2/u^2} du. \quad \text{--- (4)}$$

Let $au = x$ then $du = dx/a$ and so (2a) gives

$$I(a,b) = \frac{1}{a} \int_0^{\infty} \exp\left(-x^2 - \frac{a^2 b^2}{x^2}\right) dx$$

$$= \frac{1}{a} I(1, ab) \quad \text{--- (5)} \quad \text{(by 2a)}$$

Let $b/a = x$ then (4) gives

$$\frac{\partial I(a,b)}{\partial b} = 2b \int_0^{\infty} \frac{x^2}{b^2} \exp\left(-\frac{a^2 b^2}{x^2} - x^2\right) \frac{b}{x^2} dx$$

$$= -2 \int_0^{\infty} \exp\left(-\frac{x^2 + a^2 b^2}{x^2}\right) dx$$

$$= -2 I(1, ab) \quad \text{--- (6)}$$

Eliminating $I(1, ab)$ from (5) & (6) we get

$$\frac{\partial I}{\partial b} = -2a I$$

On integrating $\ln I = -2ab + \ln c$

$$I(a,b) = c e^{-2ab} \quad \text{--- (7)}$$

(8) $I(a,0) = L$ --- (8)

in (7) gives $I(a,b) = I(a,0) e^{-2ab}$ --- (9)

or using (2a) and (3)

Using (10) in (2) with $u = 0$

$$\int_0^{\infty} e^{-a^2 u^2 - b^2/u^2} du = \frac{\sqrt{\pi}}{2a} e^{-2ab} \quad \text{--- (10)}$$

$$L[\operatorname{erf}(\sqrt{t})] = \frac{-2}{s\sqrt{\pi}} \left[\frac{\sqrt{\pi}}{2} e^{-2\sqrt{s}} - \frac{\sqrt{\pi}}{2} \right]$$

$$= \frac{1}{s} (1 - e^{-2\sqrt{s}})$$

Ex Solve

$$\begin{aligned} x' + ky &= a \sin kt, \\ y' - kx &= a \cos kt. \end{aligned} \quad ; \quad x(0) = 0, \quad y(0) = b.$$

Taking Laplace transform we have.

$$s \bar{x}(s) - x(0) + k \bar{y} = \frac{ak}{s^2 + k^2} \quad \text{--- (1)}$$

$$s \bar{y}(s) - y(0) - k \bar{x} = \frac{as}{s^2 + k^2} \quad \text{--- (2)}$$

using conditions we get from (1) + (2) as $skx + k^2 \bar{y} =$

$$s \bar{x} + k \bar{y} = \frac{ak}{s^2 + k^2} \quad \text{--- (3)}$$

$$s \bar{y} - k \bar{x} = \frac{as}{s^2 + k^2} + b \quad \text{--- (4)}$$

Solving (3) and (4) we get

$$\bar{x} = \frac{\begin{vmatrix} \frac{ak}{s^2 + k^2} & k \\ \frac{as}{s^2 + k^2} + b & s \end{vmatrix}}{s^2 + k^2} = \frac{\begin{vmatrix} \frac{ak}{s^2 + k^2} & k \\ \frac{as}{s^2 + k^2} + b & s \end{vmatrix}}{s^2 + k^2} = \frac{\frac{aks}{s^2 + k^2} - \frac{ask}{s^2 + k^2} + bk}{s^2 + k^2} \quad \text{--- (5)}$$

Similarly $\bar{y} = \frac{a}{s^2 + k^2} + \frac{bs}{s^2 + k^2}$ from (4) + (5)

Laplace inversion yields

$$\boxed{x(t) = -b \sin kt} \quad ; \quad y(t) = \frac{a}{k} \sin kt + b \cos kt.$$

~~14/12/05~~