

The eigenvalue equation is

$$H\Psi_n = E_n \Psi_n.$$

Assuming that the  $n$ -th eigenstate  $\Psi_n$  and eigenvalue  $E_n$  can be expanded as a power series in  $\lambda$ :

$$E_n = \sum_i \lambda^i E_n^{(i)} = E_n^{(0)} + \lambda E_n^{(1)} + \lambda^2 E_n^{(2)} + \dots$$

$$\Psi_n = \sum_i \lambda^i \Psi_n^{(i)} = \Psi_n^{(0)} + \lambda \Psi_n^{(1)} + \lambda^2 \Psi_n^{(2)} + \dots$$

where

$E_n^{(0)}$  —  $n$ -th eigenenergy of  $H_0$ .

$\Psi_n^{(0)}$  —  $n$ -th eigenstate of  $H_0$ .

The other terms  $E_n^{(1)}, E_n^{(2)}, E_n^{(3)}, \dots$  are the first, second, third, ... order

corrections to the unperturbed eigenenergy  $E_n^{(0)}$ . Similarly  $\Psi_n^{(1)}, \Psi_n^{(2)}, \Psi_n^{(3)}, \dots$

are the first, second, third, ... order corrections to  $\Psi_n^{(0)}$ . For physically exactable problems the perturbation series expansion should be convergent.

## 1) Non-Degenerate Case.

The time-independent Schrodinger equation is

$$H \psi_n = E_n \psi_n.$$

$$(H_0 + \lambda H_1) \psi_n = E_n \psi_n.$$

$$(H_0 + \lambda H_1) (\psi_n^{(0)} + \lambda \psi_n^{(1)} + \lambda^2 \psi_n^{(2)} + \dots).$$

$$= (E_n^{(0)} + \lambda E_n^{(1)} + \lambda^2 E_n^{(2)} + \dots) \psi_n^{(0)}$$

$$(\psi_n^{(0)} + \lambda \psi_n^{(1)} + \lambda^2 \psi_n^{(2)} + \dots)$$

By comparing the coefficients of  $\lambda^0, \lambda^1, \lambda^2, \dots$  one obtains

$$O(\lambda^0): H_0 \psi_n^{(0)} = E_n^{(0)} \psi_n^{(0)}$$

$$O(\lambda^1): H_0 \psi_n^{(1)} + H_1 \psi_n^{(0)} = E_n^{(0)} \psi_n^{(1)} + E_n^{(1)} \psi_n^{(0)}$$

$$O(\lambda^2): H_0 \psi_n^{(2)} + H_1 \psi_n^{(1)} = E_n^{(0)} \psi_n^{(2)} + E_n^{(1)} \psi_n^{(1)} + E_n^{(2)} \psi_n^{(0)}$$

It is convenient to fix the normalization of  $\psi_n$  by

$$(\psi_n^{(0)} | \psi_n^{(0)}) = \int \psi_n^{(0)*} \psi_n^{(0)} dx = 1$$

$$\lambda (\Psi_n^{(0)}, \Psi_n) + \lambda^2 (\Psi_n^{(0)}, \Psi_n^{(2)}) + \dots = 0.$$

or equivalently

$$(\Psi_n^{(0)}, \Psi_n^{(1)}) = (\Psi_n^{(0)}, \Psi_n^{(2)}) = \dots = 0.$$

### First Order Correction:

The first order perturbation is obtained by the equation

$$H_0 \Psi_n^{(1)} + H_1 \Psi_n^{(0)} = E_n^{(1)} \Psi_n^{(1)} + E_n^{(2)} \Psi_n^{(0)}$$

Since the exact states  $\Psi_n^{(0)}$  form a complete orthonormal set, the states  $\Psi_n^{(1)}$   $i \geq 1$ , can be expanded as a linear combination of  $\Psi_n^{(0)}$ .

$$\Psi_n^{(i)} = \sum_m C_{nm}^{(i)} \Psi_m^{(0)} \quad i \geq 1$$

For  $i=1$ , we have

$$\Psi_n^{(1)} = \sum_m C_{nm}^{(1)} \Psi_m^{(0)}$$

So the first order approximation is now expressed as

$$\sum_m C_{nm}^{(1)} \underbrace{(H_0 \psi_m^{(0)})}_{H_0 \psi_m^{(0)}} + H_1 \psi_n^{(0)} = \sum_m C_{nm}^{(1)} E_m^{(0)} \psi_m^{(0)} + E_n^{(1)} \psi_n^{(0)}$$

$$H_0 \psi_m^{(0)} = E_m^{(0)} \psi_m^{(0)}$$

$$\sum_m C_{nm}^{(1)} E_m^{(0)} \psi_m^{(0)} + H_1 \psi_n^{(0)} = \sum_m C_{nm}^{(1)} E_n^{(1)} \psi_m^{(0)} + E_n^{(1)} \psi_n^{(0)}$$

or

$$\sum_m C_{nm}^{(1)} (E_m^{(0)} - E_n^{(1)}) \psi_m^{(0)} + H_1 \psi_n^{(0)} = E_n^{(1)} \psi_n^{(0)} \quad (1)$$

By taking the inner product with  $\psi_k^{(0)}$ , we get

$$\sum_m C_{nm}^{(1)} (E_m^{(0)} - E_n^{(1)}) (\psi_k^{(0)}, \psi_m^{(0)}) + (\psi_k^{(0)}, H_1 \psi_n^{(0)}) = E_n^{(1)} (\psi_k^{(0)}, \psi_n^{(0)})$$

$$\rightarrow (2)$$

when  $k = n$ , equation (2) gives

$$E_n^{(1)} = \left( \psi_n^{(0)}, H_1 \psi_n^{(0)} \right)$$

$$E_n^{(1)} = (H_1)_{nn}$$

where  $(H_1)_{nn}$  is the matrix element of  $H_1$ .

The first order correction in the  $n$ -th energy eigenvalue due to perturbation  $H_1$  is the expectation value of  $H_1$  in the unperturbed  $n$ -th eigenstate. (4)

when  $k \neq n$ , the only non-zero terms are  $(k = m)$ .

$$c_{nk}^{(1)} (E_k^{(0)} - E_n^{(0)}) = - (\psi_k^{(0)}, H_1 \psi_n^{(0)})$$

this gives.

$$c_{nk}^{(1)} = \frac{(\psi_k^{(0)}, H_1 \psi_n^{(0)})}{(E_n^{(0)} - E_k^{(0)})} \quad (k \neq n)$$

Let us consider the state.

$$\psi_n \approx \psi_n^{(0)} + \lambda \psi_n^{(1)}$$

The normalization of  $\psi_n$  gives.

$$1 = (\psi_n, \psi_n)$$

$$= 1 + \lambda (c_{nn}^{(1)} + c_{nn}^{(1)})$$

or.

$$c_{nn}^{(1)} + c_{nn}^{(1)} = 0$$

Let

$$C_{nn}^{(1)} = \alpha + i\beta.$$

$$C_{nn}^{(1)*} + C_{nn}^{(1)} = 2\alpha$$

$$\boxed{\alpha = 0.}$$

The normalization implied that the real part of  $C_{nn}^{(1)}$  vanishes i.e.

$$C_{nn}^{(1)} = i\beta.$$

The first order  $n$ th ~~correct~~ wavefn.  $\psi_n$  can be written as:

$$\psi_n = \psi_n^{(0)} + \lambda \psi_n^{(1)}$$

$$= \psi_n^{(0)} + \lambda \left[ \sum_{m \neq n} C_{nm}^{(1)} \psi_m^{(0)} + C_{nn}^{(1)} \psi_n^{(0)} \right]$$

$$= \psi_n^{(0)} + \lambda \left[ \sum_{m \neq n} C_{nm}^{(1)} \psi_m^{(0)} + i\beta \psi_n^{(0)} \right]$$

$$= (1 + i\lambda\beta) \psi_n^{(0)} + \lambda \sum_{m \neq n} C_{nm}^{(1)} \psi_m^{(0)}$$

Since we are interested in normalization to the first order in  $\lambda$ , therefore, we may write.

$$1 + i\lambda\beta \approx \exp(i\beta\lambda)$$

so that.

$$\psi = e^{i\beta\lambda} \psi_n^{(0)} + \lambda \sum_{m \neq n} C_{nm}^{(1)} \psi_m^{(0)} \quad (5)$$

Let us look at the probability density:

$$\begin{aligned} P &= |\psi \psi^*| \\ &= \left( e^{i\beta\lambda} \psi_n^{(0)} + \lambda \sum_{m \neq n} C_{nm}^{(1)} \psi_m^{(0)} \right) \left( e^{-i\beta\lambda} \psi_n^{(0)*} + \lambda \sum_{m \neq n} C_{nm}^{*(1)} \psi_m^{(0)*} \right) \\ &= \psi_n^{(0)} \psi_n^{(0)*} + \lambda \sum_{m \neq n} C_{nm}^{(1)*} e^{i\beta\lambda} \psi_n^{(0)} \psi_m^{(0)*} \\ &\quad + \lambda \sum_{m \neq n} C_{nm}^{(1)} e^{-i\beta\lambda} \psi_m^{(0)} \psi_n^{(0)*} + \lambda^2 \end{aligned}$$

or.

$$(\psi_n, \psi_n) = (\psi_n^{(0)}, \psi_n^{(0)}) + 0 + 0 + \dots$$

This says that the total probability does not depend on the phase  $\beta$ . Therefore we can make any choice of phase,

say,  $\beta = 0$ .

This gives

$$C_{nn}^{(1)} = 0$$

So, the  $n$ -th wavefunction  $\psi_n$  and energy eigenvalue  $E_n$  with first order correction are respectively.

$$\psi_n = \psi_n^{(0)} + \sum_{m \neq n} \frac{(H_1)_{mn}}{E_n^{(0)} - E_m^{(0)}} \psi_m^{(0)}$$

$\lambda=1$

$$E_n = E_n^{(0)} + (H_1)_{nn}$$

### Second Order Correction:

For the 2nd order correction, we

take

$$H_0 \psi_n^{(2)} + H_1 \psi_n^{(1)} = E_n^{(2)} \psi_n^{(2)} + E_n^{(1)} \psi_n^{(1)} + E_n^{(0)} \psi_n^{(0)}$$

we know

$$\psi_n^{(z)} = \sum_m C_{nm}^{(z)} \psi_m^{(0)}, \quad z=2$$

$$\sum_m C_{nm}^{(2)} H_0 \psi_m^{(0)} + \sum_m C_{nm}^{(1)} H_1 \psi_m^{(0)}$$

$$E_n^{(0)} \psi_n^{(2)} = \sum_m C_{nm}^{(2)} E_n^{(0)} \psi_m^{(0)} + \sum_m C_{nm}^{(1)} E_n^{(1)} \psi_m^{(0)} + E_n^{(2)} \psi_n^{(2)}$$

or

$$\sum_m C_{nm}^{(2)} (E_m^{(0)} - E_n^{(0)}) \psi_m^{(0)} = \sum_m C_{nm}^{(1)} (E_n^{(1)} - H_1) \psi_m^{(0)} + E_n^{(2)} \psi_n^{(2)}$$



(6)

Multiplying from the left by  $\psi_n^{(0)*}$  and integrate over whole space:

$$\sum_m C_{nm}^{(2)} (E_m^{(0)} - E_n^{(0)}) (\psi_n^{(0)}, \psi_m^{(0)}) = \sum_m C_{nm}^{(1)} E_n^{(1)} (\psi_n^{(0)}, \psi_m^{(0)}) - \sum_m C_{nm}^{(1)} (\psi_n^{(0)}, H_1 \psi_m^{(0)}) + E_n^{(2)} (\psi_n^{(0)}, \psi_n^{(0)})$$

⇒

$$\sum_m C_{nm}^{(2)} (E_m^{(0)} - E_n^{(0)}) \delta_{nm} = \sum_m C_{nm}^{(1)} E_n^{(1)} \delta_{nm} - \sum_m C_{nm}^{(1)} (\psi_n^{(0)}, H_1 \psi_m^{(0)}) + E_n^{(2)}$$

we know, when

$$C_{nn}^{(1)} = 0, \quad m = n$$

$$C_{nm}^{(1)} = \frac{(H_1)_{mn}}{(E_n^{(0)} - E_m^{(0)})}, \quad m \neq n.$$

therefore, when  $m \neq n$ .

$$E_n^{(2)} = \sum_{m \neq n} C_{nm}^{(1)} (\psi_n^{(0)}, H_1 \psi_m^{(0)})$$

$$E_n^{(2)} = \sum_{m \neq n} C_{nm}^{(1)} (H_1)_{nm}$$

$$E_n^{(2)} = \sum_{m \neq n} \frac{(H_1)_{nm} (H_1)_{mn}}{E_n^{(0)} - E_m^{(0)}}$$

The required 2nd order correction in  $E_n$ .

The 2nd order correction to  $\psi_n$

we know the eq.

$$\sum_m C_{nm}^{(2)} (E_m^{(0)} - E_n^{(0)}) \psi_m^{(0)} = \sum_m C_{nm}^{(1)} (E_n^{(1)} - H_1) \psi_m^{(0)} + E_n^{(2)} \psi_n^{(0)}$$

Taking the inner product with  $\psi_k^{(0)}$

$$\Rightarrow \sum_m C_{nm}^{(2)} (E_m^{(0)} - E_n^{(0)}) (\psi_k^{(0)}, \psi_m^{(0)}) = 1 \quad \begin{matrix} k \neq n \\ k \neq m \end{matrix}$$

$$= \sum_m C_{nm}^{(1)} E_n^{(1)} (\psi_k^{(0)}, \psi_m^{(0)}) - \sum_m C_{nm}^{(1)} (\psi_k^{(0)}, H_1 \psi_m^{(0)}) + E_n^{(2)} (\psi_k^{(0)}, \psi_n^{(0)}) = 0$$

or

$$C_{nk}^{(2)} (E_k^{(0)} - E_n^{(0)}) = \sum_m C_{nm}^{(1)} E_n^{(1)} \delta_{km} - \sum_m C_{nm}^{(1)} (H_1)_{km}$$

$$C_{nk}^{(2)} = \frac{1}{(E_k^{(0)} - E_n^{(0)})} \sum_m C_{nm}^{(1)} (E_n^{(1)} \delta_{km} - (H_1)_{km})$$

Since

$$C_{nm}^{(1)} = \frac{(H_1)_{mn}}{E_n^{(0)} - E_m^{(0)}} \quad m \neq n$$

$$E_n^{(1)} = (H_1)_{nn}$$

One can therefore get

$$C_{nk}^{(2)} = \frac{1}{(E_k^{(0)} - E_n^{(0)})} \sum_{m \neq n} \frac{(H_1)_{mn}}{E_n^{(0)} - E_m^{(0)}} \left[ (H_1)_{nn} \delta_{km} - (H_1)_{km} \right]$$

For  $n = k$ .

Let us consider the state

$$\psi_n \approx \psi_n^{(0)} + \lambda \psi_n^{(1)} + \lambda^2 \psi_n^{(2)}$$

It is convenient to

we know from the normalization

condition that

$$(\psi_n^{(0)}, \psi_n) = 1$$

$$\lambda (\psi_n^{(1)}, \psi_n) + \lambda^2 (\psi_n^{(2)}, \psi_n) + \dots = 0$$

The normalization condition gives

$$1 = (\Psi_n, \Psi_n) \\ = 1 + \lambda (C_{nn}^{(1)} + C_{nn}^{(1)}) + \lambda^2 ( \sum_{m \neq n} |C_{nm}^{(1)}|^2 + C_{nn}^{(2)} + C_{nn}^{(2)*} )$$

Also when  $n = k \neq m$ . then  
the term  $\sum_{m \neq n} |C_{nm}^{(1)}|^2$

The co-efficient of  $\lambda^2$  is zero, i.e

$$C_{nn}^{(2)} + C_{nn}^{(2)*} = - \sum_{m \neq n} |C_{nm}^{(1)}|^2$$

The phase is chosen such that

$$C_{nn}^{(2)} = C_{nn}^{(2)*}, \quad C_{nn}^{(2)} \text{ - real}$$

So

$$C_{nn}^{(2)} = -\frac{1}{2} \sum_{m \neq n} \left| \frac{(H_{nm}^{(1)})}{E_n^{(0)} - E_m^{(0)}} \right|^2$$

Since the 2nd order correction in wavefn  $\Psi_n$  is

$$\Psi_n^{(2)} = \sum_{m \neq n} C_{nm}^{(2)} \Psi_m^{(0)} + C_{nn}^{(2)} \Psi_n^{(0)}$$

$$\psi_n^{(2)} = \sum_{m \neq n} \left[ \sum_{k \neq n} \frac{(H_1)_{nk} (H_1)_{km}^*}{(E_n^{(0)} - E_m^{(0)}) (E_n^{(0)} - E_k^{(0)})} - \frac{(H_1)_{nn} (H_1)_{nm}}{(E_n^{(0)} - E_m^{(0)})^2} \right] \psi_m^{(0)} - \frac{1}{2} \sum_{m \neq n} \left| \frac{(H_1)_{mn}}{E_n^{(0)} - E_m^{(0)}} \right|^2 \psi_n^{(0)}$$

The n-th wavefn  $\psi_n$  to the 2nd order correction is

$$\psi_n = \psi_n^{(0)} + \sum_{m \neq n} \left[ \frac{(H_1)_{nm}}{E_n^{(0)} - E_m^{(0)}} - \frac{(H_1)_{nn} (H_1)_{nm}}{(E_n^{(0)} - E_m^{(0)})^2} + \sum_{k \neq n} \frac{(H_1)_{nk} (H_1)_{km}}{(E_n^{(0)} - E_n^{(0)}) (E_n^{(0)} - E_k^{(0)})} \right] \psi_m^{(0)} - \frac{1}{2} \sum_{m \neq n} \left| \frac{(H_1)_{mn}}{E_n^{(0)} - E_m^{(0)}} \right|^2 \psi_n^{(0)}$$

and the energy is

$$E_n = E_n^{(0)} + (H_1)_{nn} + \sum_{m \neq n} \frac{(H_1)_{nm} (H_1)_{mn}}{E_n^{(0)} - E_m^{(0)}}$$