

$$= \int_0^{\infty} \exp(-st) \int_0^{\infty} H(t-\tau) f(\tau) g(t-\tau) d\tau dt$$

where $H(t-\tau)$ is the unit step function. Now changing the order of integration, we have

$$L\{f * g\} = \int_0^{\infty} \left[\int_0^{\infty} \exp(-st) H(t-\tau) g(t-\tau) dt \right] f(\tau) d\tau$$

Let $t - \tau = t'$, then on substitution, we have

$$\begin{aligned} L\{f * g\} &= \int_0^{\infty} \left[\int_{-\tau}^{\infty} \exp[-s(t' + \tau)] H(t') g(t') dt' \right] f(\tau) d\tau \\ &= \int_0^{\infty} e^{-s\tau} \left[\int_0^{\infty} \exp(-st') H(t') g(t') dt' \right] f(\tau) d\tau \end{aligned}$$

where using the property of step function that $H(t') = 0$ for $t' < 0$ and $H(t') = 1$ for $t' \geq 0$, we have

$$\int_{-\tau}^0 \exp(-st') H(t') g(t') dt' = 0$$

Therefore

$$\begin{aligned} L\{f * g\} &= \int_0^{\infty} e^{-s\tau} \left[\int_0^{\infty} e^{-st'} g(t') dt' \right] f(\tau) d\tau \\ &= \int_0^{\infty} e^{-s\tau} f(\tau) d\tau \int_0^{\infty} e^{-st'} g(t') dt' \\ &= F(s) G(s) \end{aligned}$$

6.6.4 Illustrative examples

Example 1 ✓

Calculate the following convolutions.

- (a) $t * \exp(t)$, (b) $t * \sin 2t$, (c) $\sin 2t * \exp(-t)$

Solution

(a)

$$\begin{aligned} t * \exp(t) &= \int_0^t \tau \exp(t-\tau) d\tau = \exp(t) \int_0^t \tau \exp(-\tau) d\tau \\ &= \exp(t) [\tau \{-\exp(-\tau)\} - \{\exp(-\tau)\} \cdot 1] \Big|_0^t \\ &= \exp(t) [1 - \exp(-t)(1+t)] \end{aligned}$$

(b)

$$\begin{aligned}
 t \star \sin 2t &= \int_0^t \tau \sin 2(t - \tau) d\tau \\
 &= \tau \left[\frac{\cos 2(t - \tau)}{2} - \frac{\sin 2(t - \tau)}{-2} \right] \Big|_0^t \\
 &= \frac{1}{2} t \cos t + \frac{1}{2} \sin t - \frac{1}{2} \sin 2t \\
 &= \frac{1}{2} (t \cos t - \sin 2t + \sin t)
 \end{aligned}$$

(c)

Use the formula for the integral of $\exp(ax) \cos bx$ given in appendix A.

Example 2

Using the convolution theorem, calculate the inverse Laplace transform of the function

$$(a) \frac{3}{s^2(s^2 + 9)}, \quad (b) \frac{s}{(s^2 + 9)^2}$$

Solution

(a) Writing $H(s) = F(s)G(s)$. With $F(s) = 1/s^2$ and $G(s) = 3/(s^2 + 9)$, we have $f(t) = t$ and $g(t) = \sin 3t$.

Therefore by convolution theorem

$$\begin{aligned}
 h(t) &= f(t) \star g(t) = \int_0^t f(t - \tau) g(\tau) d\tau \\
 &= \int_0^t (t - \tau) \sin 3\tau d\tau \\
 &= \left(-\frac{t \cos 3\tau}{3} + \frac{\tau \cos 3\tau}{3} - \frac{\sin 3\tau}{9} \right) \Big|_0^t \\
 &= -\frac{\sin 3t}{9} + \frac{t}{3} = \frac{1}{9} (3t - \sin 3t)
 \end{aligned}$$

(b) Here $H(s) = s/(s^2 + 9)^2$. With $F(s) = s/(s^2 + 9)$, and $G(s) = 1/(s^2 + 9)$, we have $f(t) = \cos 3t$, $g(t) = (1/3) \sin 3t$. Therefore by convolution theorem

$$\begin{aligned}
 h(t) &= L^{-1}\{H(s)\} = f(t) \star g(t) = \frac{1}{3} \int_0^t \cos 3(t - \tau) \sin 3\tau d\tau \\
 &= \frac{1}{3} \int_0^t (\cos 3t \cos 3\tau + \sin 3t \sin 3\tau) \sin 3\tau d\tau
 \end{aligned}$$

$$\begin{aligned}
&= \frac{1}{3} \int_0^t (\cos 3t \cos 3\tau \sin 3\tau + \sin 3t \sin^2 3\tau) d\tau \\
&= \frac{1}{6} \cos 3t \int_0^t \sin 6\tau d\tau + \frac{1}{3} \sin 3t \int_0^t \frac{1 - \cos 6\tau}{2} d\tau \\
&= \frac{1}{36} \cos 3t (1 - \cos 6t) + \frac{1}{6} \sin 3t \left(\tau - \frac{\sin 6\tau}{6} \right) \Big|_0^t \\
&= \frac{1}{36} [-\cos 3t + \cos 3t] + \frac{1}{6} t \sin 3t \\
&= \frac{1}{6} t \sin 3t
\end{aligned}$$

Example 3

Use convolution theorem to calculate the Laplace transform of

$$f(t) = \int_0^t (t - \beta)^3 e^{\beta} \sin \beta d\beta$$

Solution

By definition of convolution, $f(t) = t^3 \star (e^t \sin t)$. Therefore

$$\begin{aligned}
L\{f(t)\} &= L\{t^3\} L\{e^t \sin t\} = \frac{3!}{s^4} (L\{\sin t\}) \Big|_{s \rightarrow s-1} \\
&= \frac{6}{s^4} \frac{1}{(s-1)^2 + 1} = \frac{6}{s^4(s^2 - 2s + 2)}
\end{aligned}$$

6.6.5 Exercises*Complete*

1. Find the Laplace transform of the given integrals

(a) $\int_0^t (t - \beta) \sin 3\beta d\beta$, (b) $\int_0^t \exp[-(t - \beta)] \sin \beta d\beta$.

2. Find the inverse Laplace transform using the convolution theorem or otherwise.

(a) $4/[s^2(s - 2)]$, (b) $1/(s^2 + 1)^2$, (c) $1/(s^2 - 1)^2$.

3. Show that

$$L^{-1} \left\{ \frac{F(s)}{s} \right\} = \int_0^t f(\tau) d\tau \rightarrow \text{Soln by LHS}$$

(Hint: use convolution theorem with $G(s) = 1/s$).

4. Show that

$$L^{-1} \left\{ \frac{F(s)}{s^2} \right\} = \int_0^t \left\{ \int_0^\tau f(\lambda) d\lambda \right\} d\tau$$

$$= L^{-1} \left\{ \frac{1}{s} \frac{F(s)}{s} \right\} = 1 \star \int_0^t f(\lambda) d\lambda \text{ by Q3.}$$

(Hint Apply the convolution theorem to $1/s$ and $F(s)/s$.)

5. In problem 4, show that

$$\text{L.H.S.} = t \int_0^t f(\lambda) d\lambda - \int_0^t t' f(t') dt'$$

6. Solve the integral equation

$$y(t) = at + \int_0^t y(\tau) \sin(t - \tau) d\tau$$

$$(\text{Ans.: } y(t) = at + (a/3!)t^3 \quad).$$

7. Solve the integral equation by using convolution theorem:

$$y(t) = f(t) + \int_0^t g(t - \tau) y(\tau) d\tau$$

$$(\text{Ans. } y(t) = L^{-1} \{F(s)/[1 - G(s)]\} \quad).$$

8. Solve the D.Es. by Laplace transform method.

$$(a) \quad y''(t) + k^2 y(t) = f(t)$$

$$(b) \quad y''(t) - 2ky'(t) + k^2 y(t) = f(t)$$

$$(c) \quad y''(t) + \lambda y'(t) + k^2 y(t) = f(t)$$

In each case discuss the physical significance.

$$[\text{Ans.: } (b) \quad e^{-kt} y(t) = c_1 + c_2 t + \int_0^t (t - \tau) e^{-k\tau} d\tau].$$

9. Solve the problem

$$y'' + \omega^2 y = f(t), \quad y(0) = y_0, \quad y'(0) = y_1$$

and discuss the case when

$$f(t) = \begin{cases} f_0, & t_0 < t < t_1 \\ 0, & \text{for other } t \text{ values} \end{cases}$$

$$(\text{Ans. } y(t) = y_0 \cos \omega t + (y_1/\omega) \sin \omega t + (1/\omega) f(t) * \sin \omega t \quad).$$

10. Solve the inhomogeneous problems with zero initial conditions i.e. $u(0) = 0$ and $u'(0) = 0$.

$$(a) \quad u'' + au = 1$$

$$(b) \quad u'' + u = t$$

$$(c) \quad u'' + 2u' = 1 - \exp(-t)$$

$$(d) \quad u'' - u = 1$$

$$(e) \quad u'' + 4u = \sin t$$

11. Solve the I.V.P.

$$y''(t) + ty'(t) - y(t) = 0, \quad y(0) = 0, \quad y'(0) = 1$$

(Hint: On applying L.T. and using I.Cs. we obtain

$$Y'(s) + (2 - s^2)/s Y(s) = -1/s.$$

The integrating factor for this DE is $s^2 \exp(-s^2/2)$, and the solution is given by $Y(s) = 1/s^2 + c \exp(s^2/2)$.

→ formula to study on page 243.

6.7 Laplace and Inverse Laplace Transforms of some other Functions

6.7.1 Some useful results

The following results will be used in the sequel.

The Gaussian Integral

$$\int_{-\infty}^{+\infty} \exp(-x^2) dx = \sqrt{\pi}$$

Proof

Let

$$I = \int_{-\infty}^{+\infty} \exp(-x^2) dx = \int_{-\infty}^{+\infty} \exp(-y^2) dy$$

Therefore

$$\begin{aligned} I^2 &= \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} \exp(x^2) \exp(-y^2) dx dy \\ &= \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} \exp[-(x^2 + y^2)] dx dy \end{aligned}$$

Reverting to the polar coordinates (r, θ) , we have

$$\begin{aligned} I^2 &= \int_{r=0}^{r=\infty} \int_{\theta=0}^{\theta=2\pi} \exp(-r^2) r dr d\theta \\ &= 2\pi \int_0^{\infty} \exp(-r^2) r dr \end{aligned}$$

$$= \pi \int_0^{\infty} \exp(-p) dp, \quad (p = r^2)$$

Hence

$$I = \int_{-\infty}^{+\infty} \exp(-x^2) dx = \sqrt{\pi} \text{ and } \int_0^{\infty} \exp(-x^2) dx = \sqrt{\pi}/2$$

The Gamma Function

The gamma function $\Gamma(x)$ is defined by

$$\Gamma(x) = \int_0^{\infty} e^{-t} t^{x-1} dt$$

From this definition it follows from integration by parts that

$$\begin{aligned} \Gamma(x+1) &= \int_0^{\infty} e^{-t} t^x dt \\ &= t^x (-e^{-t}) \Big|_0^{\infty} + \int_0^{\infty} e^{-t} x t^{x-1} dt \\ &= 0 + x \int_0^{\infty} e^{-t} t^{x-1} dt = x \Gamma(x) \end{aligned}$$

Hence $\Gamma(x+1) = x\Gamma(x)$.

$$\text{Also } \Gamma(1) = \int_0^{\infty} e^{-t} t^{1-1} dt = \int_0^{\infty} e^{-t} dt = 1.$$

Therefore with $x = n$, n a positive integer, using the above relation, we obtain

$$\begin{aligned} \Gamma(n+1) &= n\Gamma(n) = n(n-1)(n-2)\cdots 2 \cdot 1 \cdot \Gamma(1) \\ &= n(n-1)(n-2)\cdots 2 \cdot 1 \cdot 1 = n! \end{aligned}$$

where we have used the result $\Gamma(1) = 1$.

From this result it is clear that the gamma function can be regarded as the generalization of the factorial function.

6.7.2 Laplace transform of the step function

By definition

$$\begin{aligned} L\{H(t-t_0)\} &= \int_0^{\infty} e^{-st} H(t-t_0) dt \\ &= \int_0^{t_0} e^{-st} H(t-t_0) dt + \int_{t_0}^{\infty} e^{-st} H(t-t_0) dt \end{aligned}$$

$$\begin{aligned}
 &= 0 + \int_{t_0}^{\infty} e^{-st} dt, \quad (H(t - t_0) = 1) \\
 &= \left. \frac{e^{-st}}{-s} \right|_{t_0}^{\infty} = \frac{e^{-st_0}}{s}
 \end{aligned}$$

Hence $L\{H(t - t_0)\} = e^{-st_0}/s$.

In particular when $t_0 = 0$, $L\{H(t)\} = 1/s$.

6.7.3 Laplace transform of the logarithmic function

To calculate $L\{\ln t\}$, proceed as follows.

$$L\{\ln t\} = \int_0^{\infty} e^{-st} \ln t dt$$

Now we put $st = u$, in the integral on the right side, and obtain

$$\begin{aligned}
 L\{\ln t\} &= \int_0^{\infty} e^{-u} (\ln u - \ln s) \frac{du}{s} \\
 &= \frac{1}{s} \int_0^{\infty} e^{-u} \ln u du - \frac{\ln s}{s} \int_0^{\infty} e^{-u} du \\
 &= \frac{1}{s} \int_0^{\infty} e^{-u} \ln u du - \frac{\ln s}{s} \times 1
 \end{aligned}$$

By definition $\Gamma(x + 1) = \int_0^{\infty} e^{-u} u^x du$. On differentiating

$$\Gamma'(x + 1) = \int_0^{\infty} e^{-u} u^x \ln u du$$

or on putting $x = 0$, we have

$$\Gamma'(1) = \int_0^{\infty} e^{-u} \ln u du$$

Therefore on substitution

$$L\{\ln t\} = (\Gamma'(1) - \ln s)/s$$

where $\Gamma'(1)$ is a constant called *Euler's constant*, whose value is approximately 0.577215665.

6.7.4 Laplace transform of functions of the form $t^n f(t)$

Theorem

If $f(t)$ is a function of exponential order c , then

$$L\{t^n f(t)\} = (-1)^n \left(\frac{d}{ds}\right)^n F(s), \text{ for } s > a$$

Proof

We know that

$$F(s) = \int_0^{\infty} \exp(-st) f(t) dt$$

Differentiating both sides w.r.t. s , we have

$$\begin{aligned} F'(s) &= \int_0^{\infty} (d/ds) \exp(-st) f(t) dt \\ &= \int_0^{\infty} (-t) \exp(-st) f(t) dt = (-1) L\{t f(t)\} \end{aligned}$$

which is equivalent to the statement

$$L\{t f(t)\} = -F'(s) = \left(-\frac{d}{ds}\right) F(s)$$

Repeating the same process, we have

$$L\{t^2 f(t)\} = \left(-\frac{d}{ds}\right)^2 F(s)$$

and finally

$$L\{t^n f(t)\} = \left(-\frac{d}{ds}\right)^n F(s) = (-1)^n \left(\frac{d}{ds}\right)^n F(s)$$

which is equivalent to the statement

$$L\{(-t)^n f(t)\} = \left(-\frac{d}{ds}\right)^n F(s) = \left(\frac{d}{ds}\right)^n F(s)$$

Corollary

The following result can be deduced as a corollary from the above.

$$L\{p(t) f(t)\} = p(-D) F(s)$$

where $p(t)$ is a polynomial in t , $D = d/ds$ and $F(s) = L\{f(t)\}$.

Let

$$p(t) = a_0 + a_1 t + a_2 t^2 + a_3 t^3 + \dots + a_n t^n = \sum_i a_i t^i$$

Then

$$\begin{aligned} L\{p(t) f(t)\} &= L\left\{\sum_i (a_i t^i) f(t)\right\} \\ &= \sum_i a_i L\{t^i f(t)\} = \sum_i a_i \left(-\frac{d}{ds}\right)^i L\{f(t)\} \\ &= \sum_i a_i (-D)^i F(s) = p(-D) F(s) \end{aligned}$$

6.7.5 Laplace transforms of Bessel functions

In many physical problems, we have to calculate Laplace transforms of Bessel functions. The calculation for different Bessel functions of the first kind of order 0 i.e. $J_0(t)$ is illustrated in examples 1 and 2 below.

6.7.6 The second shifting/translation theorem

Just as the first shifting theorem enables us to calculate Laplace transforms of products of functions of the type $e^{kt} f(t)$, the second shifting theorem in a similar fashion enables us to compute inverse Laplace transforms of functions of the form $e^{-as} F(s)$. It can be stated as

$$L^{-1}\{e^{-as} F(s)\} = H(t-a) f(t-a), \quad (a > 0)$$

Proof

By definition

$$L\{H(t-a) f(t-a)\} = \int_0^\infty \exp(-st) H(t-a) f(t-a) dt$$

If we let $t-a = t'$, then the limits of integration on the right side will vary from $-a$ to ∞ . Therefore

$$\begin{aligned} L\{H(t-a) f(t-a)\} &= \int_{-a}^{\infty} \exp[-s(t'+a)] H(t') f(t') dt' \\ &+ \int_0^\infty \exp[-s(t'+a)] H(t') f(t') dt' \\ &= 0 + \exp(-as) \int_0^\infty \exp(-st') H(t') f(t') dt' \\ &= \exp(-as) L\{f\} = \exp(-as) F(s) \end{aligned}$$

where we have used the defining properties of the step function.

6.7.7 Illustrative examples

Example 1

Find the Laplace transform of $J_0(t)$ where

$$J_0(t) = (1/\pi) \int_0^\pi \cos(t \sin \theta) d\theta$$

Solution

$$L\{J_0(t)\} = \int_0^\infty \exp(-st) \frac{1}{\pi} \left\{ \int_0^\pi \cos(t \sin \theta) d\theta \right\} dt$$

Reversing the order of integration, we have

$$L\{J_0(t)\} = \frac{1}{\pi} \int_0^\pi \left\{ \int_0^\infty e^{-st} \cos(t \sin \theta) dt \right\} d\theta$$

Now let I denote the integral in braces. Then using the formula

$$\int e^{ax} \cos bx dx = \frac{e^{ax}}{a^2 + b^2} (a \cos bx + b \sin bx)$$

we have

$$\begin{aligned} I &= \int_0^\infty \exp(-st) \cos(t \sin \theta) dt \\ &= \frac{\exp(-st)}{s^2 + \sin^2 \theta} [-s \cos(t \sin \theta) + \sin \theta \cdot \sin(t \sin \theta)] \Big|_0^\infty \\ &= 0 - \frac{1}{s^2 + \sin^2 \theta} \cdot [-s] = \frac{s}{s^2 + \sin^2 \theta} \end{aligned}$$

On substitution, we obtain

$$\begin{aligned} L\{J_0(t)\} &= \frac{s}{\pi} \int_0^\pi \frac{d\theta}{s^2 + \sin^2 \theta} \\ &= \frac{s}{\pi} \int_{-\pi/2}^{+\pi/2} \frac{d\theta'}{s^2 + \cos^2 \theta'}, \quad (\text{where } \theta = \theta' + \pi/2) \\ &= \frac{2s}{\pi} \int_0^{\pi/2} \frac{d\theta}{s^2 + \cos^2 \theta} \\ &= \frac{2s}{\pi} \int_0^{\pi/2} \frac{\sec^2 \theta d\theta}{1 + s^2 \sec^2 \theta} \\ &= \frac{2s}{\pi} \int_0^{\pi/2} \frac{\sec^2 \theta d\theta}{1 + s^2(1 + \tan^2 \theta)} \end{aligned}$$

Letting $u = \tan \theta$, $\sec^2 \theta d\theta = du$, we obtain

$$\begin{aligned}
 L\{J_0(t)\} &= \frac{2s}{\pi} \int_0^\infty \frac{du}{1+s^2(1+u^2)} \\
 &= \frac{2s}{\pi} \int_0^\infty \frac{du}{(1+s^2)+s^2u^2} \\
 &= \frac{2}{\pi s} \int_0^\infty \frac{du}{(1+s^2)/s^2+u^2} \\
 &= \frac{2}{\pi s} \int_0^\infty \frac{du}{a^2+u^2}, \quad a^2 = \frac{1+s^2}{s^2} \\
 &= \frac{2}{\pi s} \cdot \frac{1}{a} \cdot \tan^{-1} \frac{u}{a} \Big|_0^\infty \\
 &= \frac{2}{\pi} \cdot \frac{1}{\sqrt{1+s^2}} \cdot \frac{\pi}{2} \\
 &= \frac{1}{\sqrt{1+s^2}}
 \end{aligned}$$

Example 2

Given that Bessel functions of the first kind and positive integral order satisfy the recurrence relations

$$J_1 = -J'_0, \quad J_{n+1} = J_{n-1} - 2J'_n, \quad n \geq 1$$

with $J_0(0) = 1$, $J_n(0) = 0$, $n > 0$, show that

$$L\{J_n(t)\} = \frac{(\sqrt{s^2+1}-s)^n}{\sqrt{s^2+1}}$$

Also find the Laplace transform for $J_0(at)$, $a > 0$.

Solution

From the first recurrence relation, using the formula for Laplace transform of a derivative, we obtain

$$\begin{aligned}
 F_1(s) &= L\{J_1(t)\} = -L\{J'_0(t)\} \\
 &= -\{sL\{J_0(t)\} - J_0(0)\} \\
 &= -sL\{J_0(t)\} + J_0(0), \quad (\text{where } J_0(0) = 1) \\
 &= -sF_0(s) + 1 \tag{1}
 \end{aligned}$$

But from example 1

$$F_0(s) = L\{J_0(t)\} = 1/\sqrt{s^2+1}$$

Therefore the given formula for $L\{J_n(t)\}$ is true for $n = 0$. Also

$$\begin{aligned} F_1(s) &= L\{J_1(t)\} = -sL\{J_0(t)\} + 1 \\ &= \frac{-s}{\sqrt{s^2+1}} + 1 = \frac{\sqrt{s^2+1} - s}{\sqrt{s^2+1}} \end{aligned}$$

which shows that the given formula is true for $n = 1$. Next we use the second recurrence relation, viz. $J_{n+1} = J_{n-1} - 2J'_n$, which gives $J_2 = J_0 - 2J'_1$. Therefore

$$\begin{aligned} F_2(s) &= L\{J_2(t)\} = L\{J_0(t)\} - 2L\{J'_1(t)\} \\ &= F_0(s) - 2\{sF_1(s) - J_1(0)\} \\ &= F_0 - 2sF_1(s) + J_1(0) \\ &= F_0(s) - 2sF_1(s), \text{ because } J_1(0) = 0 \\ &= \frac{1}{\sqrt{s^2+1}} - \frac{2s(\sqrt{s^2+1} - s)}{\sqrt{s^2+1}} \\ &= \frac{1 - 2s\sqrt{s^2+1} + 2s^2}{\sqrt{s^2+1}} \\ &= \frac{(\sqrt{s^2+1} - s)^2}{\sqrt{s^2+1}} \end{aligned}$$

which shows that the formula is true for $n = 2$. Now we use mathematical induction to establish the formula for the general index $n + 1$, assuming that it is true for n and $n - 1$.

From the second recurrence relation

$$\begin{aligned} L\{J_{n+1}(t)\} &= L\{J_{n-1}(t)\} - 2L\{J'_n(t)\} \\ &= L\{J_{n-1}(t)\} - 2[sL\{J_n(t)\} - J_n(0)] \\ &= L\{J_{n-1}(t)\} - 2sL\{J_n(t)\}, \quad J_n(0) = 0, \quad n \geq 1 \\ &= \frac{(\sqrt{s^2+1} - s)^{n-1}}{\sqrt{s^2+1}} - \frac{2s(\sqrt{s^2+1} - s)^n}{\sqrt{s^2+1}} \\ &= \frac{(\sqrt{s^2+1} - s)^{n-1} \cdot [1 - 2s\sqrt{s^2+1} + 2s^2]}{\sqrt{s^2+1}} \\ &= \frac{(\sqrt{s^2+1} - s)^{n-1} \cdot (\sqrt{s^2+1} - s)^2}{\sqrt{s^2+1}} \\ &= \frac{(\sqrt{s^2+1} - s)^{n+1}}{\sqrt{s^2+1}} \end{aligned}$$

Hence the proof. To obtain the same formula when the argument is at , we use the rule of scales, viz. $L\{f(at)\} = (1/a)F(s/a)$, $a > 0$ and obtain

$$L\{J_n(at)\} = \frac{(\sqrt{s^2+a^2} - s)^n}{a^n \sqrt{s^2+a^2}}$$

Example 3 *H.W.*

Evaluate $L\{\exp(at) - (\cos bt)/t\}$ and deduce that

$$L\left\{\frac{\sin^2 t}{t}\right\} = (1/2) \ln\left(\frac{\sqrt{s^2+4}}{s}\right), \quad s > 1$$

Solution

Here we will use the result

$$L\left\{\frac{f(t)}{t}\right\} = \int_s^\infty F(s') ds' \quad (1)$$

provided the limit of $(f(t)/t)$ as $t \rightarrow 0$ exists. In this problem

$$\lim_{t \rightarrow 0} \frac{f(t)}{t} = \lim_{t \rightarrow 0} \frac{\exp(at) - \cos(bt)}{t} = \lim_{t \rightarrow 0} \frac{ae^{at} + b \sin(bt)}{1} = a$$

Also

$$F(s) = L\{f(t)\} = L\{e^{at} - \cos(bt)\} = \frac{1}{s-a} - \frac{s}{s^2+b^2}, \quad (s > a)$$

Hence from (1)

$$\begin{aligned} L\left\{\frac{f(t)}{t}\right\} &= L\left\{\frac{e^{at} - \cos bt}{t}\right\} \\ &= \int_s^\infty \left[\frac{1}{s'-a} - \frac{s'}{s'^2+b^2}\right] ds' \\ &= \left\{\ln(s'-a) - \frac{1}{2} \ln(s'^2+b^2)\right\} \Big|_s^\infty \\ &= \ln \frac{s'-a}{(s'^2+b^2)^{1/2}} \Big|_s^\infty = \ln \frac{(1-0)}{(1+0)} - \ln \frac{s-a}{\sqrt{s^2+b^2}} \\ &= 0 - \ln \frac{s-a}{\sqrt{s^2+b^2}} = \ln \frac{\sqrt{s^2+b^2}}{s-a}, \quad s > a \end{aligned}$$

To deduce the second part, we put $a = 0$, $b = 2$, so that $1 - \cos 2t = 2 \sin^2 t$.

Hence

$$L\left\{\frac{\sin^2 t}{t}\right\} = \frac{1}{2} \ln \frac{\sqrt{s^2+4}}{s}, \quad s > 0$$

6.7.8 Exercises

- Use the result $L\{\exp(-kt) f(t)\} = F(s)$, where $f(t)$ is a real function of the real variable t to show that

$$L\{\exp(-\alpha t) \cos \beta t f(t)\} = \operatorname{Re} F(s + \alpha + i\beta)$$

and

$$L\{\exp(-\alpha t) \sin \beta t f(t)\} = -\operatorname{Im} F(s + \alpha + i\beta)$$

(Hint: Take $k = \alpha + i\beta$, and equate real and imaginary parts on both sides.)

2. Evaluate $L\{\exp(kt) t^n\}$ when n is an integer, and deduce the expressions for $L\{t^n \cos kt\}$ and $L\{t^n \sin kt\}$.

Hint: We have

$$L\{\exp(kt) t^n\} = \frac{n!}{s^{n+1}} \Big|_{s \rightarrow s-k} = \frac{n!}{(s-k)^{n+1}}, \operatorname{Re}(s-k) > 0$$

Now let $k = i\alpha$, where α is real and positive, then

$$\begin{aligned} L\{\exp(i\alpha t) t^n\} &= L\{(\cos \alpha t + i \sin \alpha t) t^n\} \\ &= \frac{n!}{(s - i\alpha)^{n+1}}, \operatorname{Re} s > 0 \end{aligned}$$

Now let $s = r \cos \theta$, $\alpha = r \sin \theta$, so that

$$r = \sqrt{s^2 + \alpha^2}, \tan \theta = \alpha/s, 0 \leq \theta < \pi/2$$

Since α and s are positive, $s - i\alpha = r(\cos \theta - i \sin \theta)$, and therefore

$$\begin{aligned} \frac{1}{(s - i\alpha)^{n+1}} &= \frac{1}{r^{n+1}} (\cos \theta - i \sin \theta)^{-n-1} \\ &= \frac{1}{r^{n+1}} [\cos(n+1)\theta + i \sin(n+1)\theta] \end{aligned}$$

Hence

$$L\{\exp(i\alpha t) t^n\} = \frac{n!}{r^{n+1}} [\cos(n+1)\theta + i \sin(n+1)\theta]$$

Equating real and imaginary parts,

$$L\{\cos \alpha t t^n\} = n! (s^2 + \alpha^2)^{-(n+1)/2} \cos(n+1)\theta$$

and

$$L\{\sin \alpha t t^n\} = n! (s^2 + \alpha^2)^{-(n+1)/2} \sin(n+1)\theta$$

where $\theta = \tan^{-1} \alpha/s$.

3. Show that

$$L\left\{\frac{f(t)}{t+a}\right\} = \exp(as) \int_a^\infty \exp(-st') F(t') dt'$$

(Hint: Let $I = L\{f(t)/(t+a)\}$, then with $D = d/ds$

$$\begin{aligned}(D-a)I &= (D-a) \int_0^{\infty} \exp(-st) \frac{f(t)}{t+a} dt \\ &= \int_0^{\infty} (D-a) \exp(-st) \frac{f(t)}{t+a} dt \\ &= \int_0^{\infty} (-t-a) \exp(-st) \frac{f(t)}{t+a} dt \\ &= - \int_0^{\infty} \exp(-st) f(t) dt = -F(s)\end{aligned}$$

Also

$$dI/ds - aI = -F(s)$$

which is a nonhomogeneous linear DE with integrating factor $\mu = \exp(-as)$. The solution of this equation is given by

$$\mu I = - \int \mu F(s) ds + \text{constant}$$

or

$$\exp(-as) I = - \int \exp(-as) F(s) ds + c_1$$

To calculate the constant c_1 , we take the limit as $s \rightarrow \infty$,

$$0 = - \int_{s_0}^{\infty} \exp(-as) F(s) ds + c_1$$

which gives

$$c_1 = \int_{s_0}^{\infty} \exp(-as) F(s) ds$$

Therefore on substitution

$$\exp(-as) I = \int_s^{\infty} \exp(-as) F(s) ds$$

or finally

$$I = \exp(as) \int_s^{\infty} \exp(-as') F(s') ds'$$

4. Prove that

$$L \left\{ \exp(-at) \int_0^t f(t') \exp(at') dt' \right\} = \frac{F(s)}{s+a}$$

(Hint: Note that this is general form of a previous result.

$$\text{L.H.S.} = \int_0^{\infty} \exp[-(s+a)t] \left\{ \int_0^t f(t') \exp(at') dt' \right\} dt$$

Integrating by parts w.r.t. t , we have

$$\begin{aligned} \text{L.H.S.} &= \frac{\exp[-(s+a)t]}{-(s+a)} \cdot \int_0^t f(t') \exp(at') dt' \Big|_0^\infty \\ &\quad + \frac{1}{(s+a)} \int_0^\infty \exp[(s+a)t] \cdot e^{-at} f(t) dt \\ &= 0 + \frac{1}{(s+a)} \int_0^\infty \exp(-st) \cdot f(t) dt = \frac{F(s)}{s+a} \end{aligned}$$

5. Prove that

$$L\{\text{erfc}(t^{1/2})\} = \frac{1}{s} - \frac{1}{\sqrt{s+1}}, \text{ where } \text{erfc } t = \frac{2}{\sqrt{\pi}} \int_t^\infty \exp(-x^2) dx$$

(Hint:

$$\begin{aligned} \text{erfc}(t) &= \frac{2}{\sqrt{\pi}} \int_t^\infty \exp(-x^2) dx \\ &= \frac{2}{\sqrt{\pi}} \left[\int_0^\infty \exp(-x^2) dx - \int_0^t \exp(-x^2) dx \right] \end{aligned}$$

or

$$\text{erfc}(t) = \frac{2}{\sqrt{\pi}} \cdot \frac{\sqrt{\pi}}{2} - \frac{2}{\sqrt{\pi}} \int_0^t \exp(-x^2) dx = 1 - \text{erf}(t)$$

where

$$\text{erf}(t) = (2/\sqrt{\pi}) \int_0^t \exp(-x^2) dx$$

Therefore

$$\begin{aligned} L\{\text{erfc}(\sqrt{t})\} &= L\{1\} - L\{\text{erf}(\sqrt{t})\} \\ &= \frac{1}{s} - \frac{2}{\sqrt{\pi}} L\left\{ \int_0^{\sqrt{t}} \exp(-x^2) dx \right\} \\ &= \frac{1}{s} - \frac{2}{\sqrt{\pi}} L\left\{ \int_0^t \exp(-v) \frac{dv}{2\sqrt{v}} \right\}, \quad (x^2 = v) \\ &= \frac{1}{s} - \frac{1}{\sqrt{\pi}} L\left\{ \int_0^t \frac{\exp(-t')}{\sqrt{t'}} dt' \right\} \\ &= \frac{1}{s} - \frac{1}{s\sqrt{\pi}} L\left\{ \frac{1}{\sqrt{t}} \right\} \Big|_{s \rightarrow s+1} \\ &= \frac{1}{s} - \frac{1}{s\sqrt{s+1}} \end{aligned}$$

$$\text{erf}(t) = \frac{2}{\sqrt{\pi}} \int_0^t e^{-x^2} dx$$

$$\text{erfc}(t) = \frac{2}{\sqrt{\pi}} \int_t^\infty e^{-x^2} dx$$

$$\therefore L\left\{ \frac{1}{\sqrt{t}} \right\} = \sqrt{\frac{\pi}{s}}$$

Therefore

$$L\{\text{erfc} \sqrt{t}\} = \frac{1}{s} \left(1 - \frac{1}{\sqrt{s+1}} \right)$$

using $L\left\{ \int_0^t f(x) dx \right\} = \frac{1}{s} L\left\{ \frac{f(t)}{t} \right\}$
 $\Rightarrow L\left\{ \int_0^t \frac{e^{-x^2}}{\sqrt{x}} dx \right\} = \frac{1}{s} L\left\{ \frac{e^{-t}}{\sqrt{t}} \right\}$

w that

$$(a) \quad L \left\{ \frac{\exp(-t)}{\sqrt{t}} \right\} = \sqrt{\frac{\pi}{s+1}}$$

$$(b) \quad L \left\{ \int_0^t \exp(-v) \frac{1}{\sqrt{v}} dv \right\} = \frac{1}{s} \sqrt{\frac{\pi}{s+1}}$$

$$(c) \quad L \left\{ \exp(t) \operatorname{erfc} \sqrt{t} \right\} = \frac{1}{s + \sqrt{s}}$$

Further Discussion of Inverse Laplace Transform

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1 Uniqueness of the inverse Laplace transform

Wh's theorem guarantees the uniqueness of inverse Laplace transform up to certain additive function called the *null function*. A function $N(x)$ is called a null function if $\int_0^\infty N(x) dx = 0$.

The theorem asserts that if $f(t)$ and $g(t)$ are inverse Laplace transforms of $F(s)$, then $f(t) - g(t) = N(t)$ where $N(t)$ is a null function. If $f(t)$ is continuous, then it will be unique inverse Laplace transform of $F(s)$.

3.2 Heaviside expansion theorem

If $M(s)$ and $N(s)$ are polynomials of degree m and n respectively with $m < n$, and $N(s)$ has n distinct zeros α_i , $i = 1, 2, 3, \dots, n$, none of which is a zero of $M(s)$, then

$$L^{-1} \left\{ \frac{M(s)}{N(s)} \right\} = \sum_{i=1}^n \frac{M(\alpha_i)}{N'(\alpha_i)} \exp(\alpha_i t)$$

If $N(s)$ has a repeated root α_1 of multiplicity r while other roots at $\alpha_2, \alpha_3, \dots, \alpha_n$ are not repeated, the corresponding formula is given by

$$L^{-1} \left\{ \frac{M(s)}{N(s)} \right\} = \sum_{i=2}^n \frac{M(\alpha_i)}{N'(\alpha_i)} \exp(\alpha_i t) + \sum_{j=1}^r \frac{1}{(j-1)!} \left\{ \frac{d^{j-1}}{ds^{j-1}} (s - \alpha_1)^j F(s) \right\} \exp(\alpha_1 t) \Big|_{s=\alpha_1}$$

$L\left(\frac{1}{\sqrt{t}}\right) = \sqrt{\frac{\pi}{s}}$ and then $L\left(\frac{1}{\sqrt{t}}\right) = \sqrt{\frac{\pi}{s}}$

$$c_2 = \lim_{s \rightarrow s_2} (s - s_2) F'(s) = \lim_{s \rightarrow s_2} \frac{1}{s^2(s - s_1)}$$

$$= \frac{1}{s_2^2(s_2 - s_1)}$$

Next we express c_1, c_2 in terms of α and β .

$$s_2 = -\alpha + i\beta = \sqrt{\alpha^2 + \beta^2} e^{i\theta}$$

where $\tan \theta = -\beta/\alpha$.

$$s_1 = -\alpha - i\beta = \sqrt{\alpha^2 + \beta^2} \exp(-i\theta)$$

Therefore

$$s_1^2 = (\alpha^2 + \beta^2) \exp(2i\theta) = b^2 \exp(2i\theta)$$

and

$$s_2^2 = (\alpha^2 + \beta^2) \exp(-2i\theta) = b^2 \exp(-2i\theta)$$

Also $s_1 - s_2 = 2i\beta$. Hence we can write

$$c_1 = \frac{1}{b^2(2i\beta)} e^{-2i\theta}, \quad c_2 = \frac{1}{b^2(-2i\beta)} e^{2i\theta}$$

Finally

$$f(t) = \frac{t}{b^2} - \frac{2\alpha}{b^4} + c_1 e^{s_1 t} + c_2 e^{s_2 t}$$

where the last two terms (T_{34}) can be further simplified as follows.

$$\begin{aligned} T_{34} &= \frac{1}{2i\beta b^2} e^{-2i\theta} e^{(\alpha+i\beta)t} + \frac{1}{-2i\beta b^2} e^{2i\theta} e^{(-\alpha-i\beta)t} \\ &= \frac{1}{2i\beta b^2} e^{-\alpha t} \left[e^{i(2\theta+\beta)t} - e^{-i(2\theta+\beta)t} \right] \\ &= \frac{e^{-\alpha t}}{2i\beta b^2} 2i \sin(\beta - 2\theta) \\ &= \frac{e^{-\alpha t}}{\beta b^2} \sin(\beta + 2\theta) \end{aligned}$$

Example 3

Find the general solution of the differential equation

$$y''(t) + k^2 y(t) = f(t)$$

Solution

Taking the Laplace transform of both sides

$$L\{y''(t)\} + k^2 L\{y(t)\} = L\{f(t)\}$$

Now we take Laplace transform of derivatives, assuming that all functions are of class A) and obtain

$$s^2 Y(s) - sy(0) - y'(0) + k^2 Y(s) = F(s)$$

or

$$(s^2 + k^2) Y(s) = F(s) + sy(0) + y'(0)$$

or

$$Y(s) = [c_1 + c_2 s + F(s)] / (s^2 + k^2)$$

where $c_1 = y'(0)$ and $c_2 = y(0)$ are constants. Therefore

$$\begin{aligned} y(t) &= L^{-1}\{Y(s)\} \\ &= L^{-1}\left\{\frac{c_1}{s^2 + k^2}\right\} + L^{-1}\left\{\frac{c_2 s}{s^2 + k^2}\right\} \\ &\quad + L^{-1}\left\{\frac{F(s)}{s^2 + k^2}\right\} \\ &= \frac{c_1}{k} \sin kt + c_2 \cos kt + L^{-1}\left\{\frac{F(s)}{s^2 + k^2}\right\} \\ &= \frac{c_1}{k} \sin kt + c_2 \cos kt + \frac{1}{k} \sin kt \star f(t) \end{aligned}$$

where we have used the convolution theorem. Finally we have

$$y(t) = \frac{c_1}{k} \sin kt + c_2 (\cos kt) + \frac{1}{k} \int_0^t \sin(t - \tau) f(\tau) d\tau$$

Example 4

Solve the I.V.P.

$$y''(t) + ty'(t) - y(t) = 0, \quad y(0) = 0, \quad y'(0) = 1$$

Solution

Taking the Laplace transform of both the sides

$$L\{y''\} + L\{ty'(t)\} - L\{y(t)\} = 0$$

or

$$s^2 Y(s) - sy(0) - y'(0) + \left(-\frac{d}{ds}\right) L\{y'(t)\} - Y(s) = 0$$

or on substituting for the initial values, we have

$$s^2 Y(s) - 1 - \frac{d}{ds} \{sY(s) - y(0)\} - Y(s) = 0$$

or

$$s^2 Y(s) - 1 - \{sY'(s) + Y(s)\} - Y(s) = 0$$

or

$$(s^2 - 1)Y(s) - sY'(s) - Y(s) = 1$$

or

$$-sY'(s) + (s^2 - 2)Y(s) = 1$$

or

$$Y'(s) + \frac{2 - s^2}{s} Y(s) = \frac{-1}{s} \quad (1)$$

which is of the form $Y' + p(x)Y = q(x)$, (i.e. it is first order linear nonhomogeneous DE). Its integrating factor $\mu(x)$ is given by

$$\mu(x) = \exp \int \left(\frac{2}{s} - s\right) ds = \exp(2 \ln s - s^2/2) = s^2 \exp(-s^2/2)$$

Therefore general solution of (1) is given by

$$\begin{aligned} Y(s)s^2 e^{-s^2/2} &= \int \left(\frac{-1}{s}\right) s^2 \exp(-s^2/2) ds + \text{constant} \\ &= \int \{-s \exp(-s^2/2)\} ds + \text{constant} \\ &= \exp(-s^2/2) + c \end{aligned}$$

where c is a constant. Therefore $Y(s) = 1/s^2 + c \exp(s^2/2)$

Now taking limit as $s \rightarrow \infty$, we have $c = 0$. Therefore $Y(s) = 1/s^2$.

Taking inverse Laplace transform, we obtain $y(t) = t$, as the required solution.

Example 5

Use Laplace transform to solve the problem.

$$y'' - ay = f(t), \quad y(0) = y_0, \quad y'(0) = y_1$$