

→ Reasoning as before the number of molecules simultaneously having velocities in the range v_x and $v_x + dv_x$, v_y and $v_y + dv_y$, v_z and $v_z + dv_z$ is given by

$$dN^3(v_x, v_y, v_z) = N f(v_x) f(v_y) f(v_z) dv_x dv_y dv_z \quad \text{--- (4)}$$

where $dN^3(v_x, v_y, v_z)$ is a differential of 3rd order eq (4) the number of endpoints in the volume elements $(dv_x dv_y dv_z)$ and (density of points in (points is unit volume). in)

the velocity space is given by

$$\rho = \frac{dN^3(v_x, v_y, v_z)}{dv_x \cdot dv_y \cdot dv_z} = N f(v_x) f(v_y) f(v_z) \quad \text{--- (5)}$$

Since velocity is isotropic, ^{same structure} density should be constant at all points on the sphere

$v_x^2 + v_y^2 + v_z^2 = v^2$ of radius v is the velocity space.

Hence $N f(v_x) f(v_y) f(v_z) = \text{constant}$ --- (6)

when $v_x^2 + v_y^2 + v_z^2 = \text{constant}$ --- (7)

The eq (7) is called the eq of constraint as it reduces independent variables from ~~3~~ three to two.

Important: partial differentiation of the eq (6) & (7)

gives $f(v_x) f(v_y) f(v_z) \frac{\partial f(v_x)}{\partial v_x} dv_x + f(v_x) f(v_y) \frac{\partial f(v_y)}{\partial v_y} dv_y + f(v_x) f(v_y) \frac{\partial f(v_z)}{\partial v_z} dv_z = 0$

Representing the derivatives by prime^{notation}ly
dividing through out by $f(v_x) f(v_y) f(v_z)$

$$\frac{f'(v_x)}{f(v_x)} dv_x + \frac{f'(v_y)}{f(v_y)} dv_y + \frac{f'(v_z)}{f(v_z)} dv_z = 0 \quad \text{--- (7)}$$

Now from eq (7)

$$2v_x dv_x + 2v_y dv_y + 2v_z dv_z = 0 \quad \text{--- (8)}$$

$$(v_x dv_x + v_y dv_y + v_z dv_z) = 0$$

Let us multiply the eq (8) with $m\beta$
called Lagrange undetermined multiplier.
(where m is mass of molecule and β
some arbitrary unknown fraction) and
add to equation (7) (8)

$$\left[\frac{f'(v_x) + m\beta v_x}{f(v_x)} \right] dv_x + \left[\frac{f'(v_y) + m\beta v_y}{f(v_y)} \right] dv_y + \left[\frac{f'(v_z) + m\beta v_z}{f(v_z)} \right] dv_z = 0 \quad \text{--- (9)}$$

Since eq (9) is an identity and for this
identity to hold^{hold} to all the terms on the L.H.S
of above eq should separately vanish
 $i = 1$

$$\left[\frac{f'(v_x) + m\beta v_x}{f(v_x)} \right] dv_x = 0 \quad \text{--- (10)}$$

Important: $\left[\frac{f'(v_y) + m\beta v_y}{f(v_y)} \right] dv_y = 0 \quad \text{--- (11)}$

$$\left[\frac{f'(v_2) + m\beta v_2}{f(v_2)} \right] dv_2 = 0 \quad \text{--- (13)}$$

Let us consider eq (13)

$$\left[\frac{f'(v_1) + m\beta v_1}{f(v_1)} \right] dv_1 = 0$$

$$\frac{f'(v_1)}{f(v_1)} dv_1 + m\beta v_1 dv_1 = 0$$

$$\frac{f'(v_1)}{f(v_1)} dv_1 = -m\beta v_1 dv_1 \quad \text{--- (14)}$$

Integrating it we get

$$\int \frac{f'(v_1)}{f(v_1)} dv_1 = - \int m\beta v_1 dv_1 \quad \leftarrow$$

$$\ln |f(v_1)| = - \frac{m\beta v_1^2}{2} + \ln A$$

$$\ln |f(v_1)| - \ln A = - \frac{m\beta v_1^2}{2}$$

$$\ln \left| \frac{f(v_1)}{A} \right| = - \frac{m\beta v_1^2}{2}$$

$$\frac{f(v_1)}{A} = e^{-\frac{1}{2}m\beta v_1^2}$$

$$f(v_1) = A e^{-\frac{1}{2}m\beta v_1^2}$$

--- (15)

Similarly

$$f(v_y) = A e^{-\frac{1}{2}m\beta v_y^2}$$

$$f(v_2) = A e^{-\frac{1}{2}m\beta v_2^2}$$

setting these values in eq (5)

$$\begin{aligned}
 P(v) &= NA e^{-\frac{1}{2}m\beta v_x^2} \cdot A e^{-\frac{1}{2}m\beta v_y^2} \cdot A e^{-\frac{1}{2}m\beta v_z^2} \\
 &= NA^3 e^{-\frac{1}{2}m\beta(v_x^2 + v_y^2 + v_z^2)} \\
 &= NA^3 e^{-\frac{1}{2}m\beta v^2} \\
 &= NA^3 e^{-\beta(\frac{1}{2}mv^2)} \quad \text{--- (6)} \\
 &= NA^3
 \end{aligned}$$

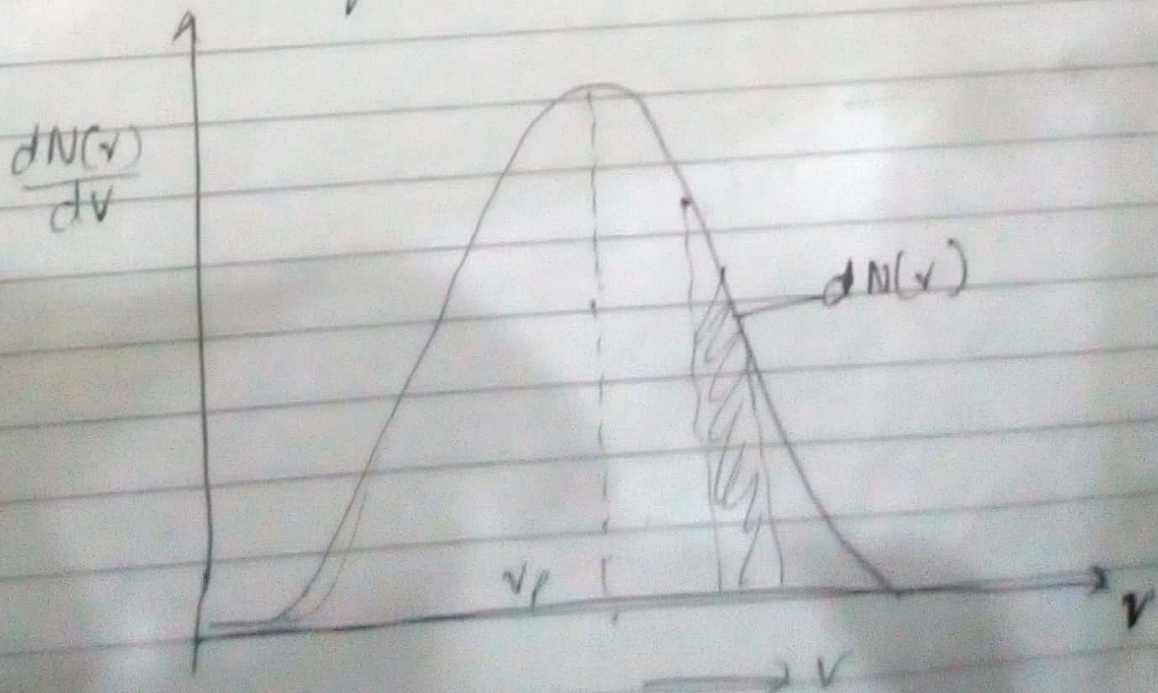
But P is constant at all points with in the infinitesimal spherical shell enclosed by the surface of radii v and $v+dv$ in the velocity space.

The volume of shell is $(4\pi v^2)dv$ which contains velocity points of $dN(v)$ molecules hence eq (6) can be put in the form

$$P(v) = \frac{dN(v)}{4\pi v^2 dv} = NA^3 e^{-\beta(\frac{1}{2}mv^2)}$$

$$\frac{dN(v)}{dv} = 4\pi NA^3 v^2 e^{-\beta(\frac{1}{2}mv^2)}$$

which is the Maxwell law of velocities distribution plotted as shown in fig



The shaded area gives the number of molecules in the range $v \text{ to } v + dv$.

N_v

→ The number of molecules having x components of velocity in the range $v_x \text{ to } v_x + dv_x$ can be obtained by combining eq (1) & (15) Thus

$$dN(v_x) = NA e^{-\beta(\frac{1}{2}mv_x^2)} dv_x \quad (18)$$

Evaluation of β and A :

We use that fact

$$\int_0^\infty dN(v) = N \quad (\text{total number of molecules})$$

Substituting for $dN(v)$ for eq (18) (17)

$$\int_0^\infty 4\pi N A^3 v^2 e^{-\beta(\frac{1}{2}mv^2)} dv = N$$

$$\int_0^\infty 4\pi A^3 v^2 e^{-\beta(\frac{1}{2}mv^2)} dv = 1 \quad (19)$$

The integral on L.H.S of the form

$$\int_0^\infty x^n e^{-ax^2} dx$$

when $n=2$ $a = \frac{1}{2}m\beta$ & its value is

$$\begin{aligned} \frac{1}{4} \sqrt{\frac{\pi}{a^3}} &= \frac{1}{4} \sqrt{\frac{\pi}{(\frac{1}{2}m\beta)^3}} \\ &= \frac{1}{4} \frac{\sqrt{\pi}}{(\frac{1}{2}m\beta)^{3/2}} \end{aligned}$$

using this value in above relation (19) we