

Line Integral:-

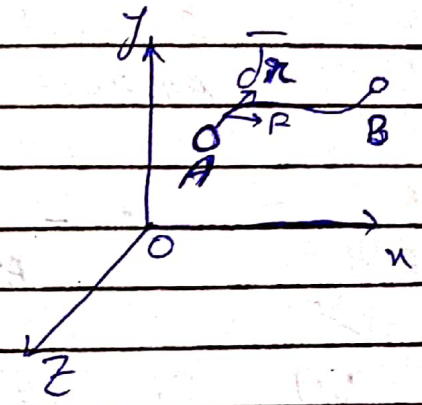
Let us consider a system of particles with finite masses under certain laws of Mechanics in an inertial frame then if the trajectory is evaluated along the curve that describes the evaluation of line integral if ϕ and A be the scalar and vector valued fn's then line integral describes for.

$$(a) \int_C \phi dr, \int_C A \cdot dr, (c) \int_C A \times dr$$

Note that $\oint_C A \cdot dr$ be taken for $C \in \mathbb{R}^n$ where (a) and (b) are vector while (b) gives a scalar quantity.

Also if the system of particles is located in 2-D or 3-D that gives for multiple integrals extended for double and triple integrals - representing the surface and volume integrals.

Let us consider a particle moving in inertial frame \mathcal{R} for infinitesimal changes or displacement, the total work done by applied force is calculated for:



$$dW = F \cdot dr \quad \text{--- (1)}$$

where for the sake of sol of Eq. (1) if we integrate on both sides then it will be referred to ordinary integral for the applied force F .

Example of line integral.

1) The line integral $\int_C F \cdot dr$ represents the work done by the force along the regular curve C for

different values of $m \in \mathbb{R}$

ii) The quantitative relationship b/w current i and magnetic field B is given by Ampere's Law:

$$\oint B \cdot d\vec{r} = \mu_0 i$$

ii) The Potential difference b/w two points A and B is related to the electric field by line integral

$$V_B - V_A = - \int_A^B E \cdot d\vec{r}$$

V_B

v) If $A = \text{grad } \phi$ then $\int_A^B A \cdot d\vec{r}$

depends only on the initial and boundary values of ϕ depending on the path also

$$\oint A \cdot d\vec{r} = 0$$

Conversely if $\oint A \cdot d\vec{r} = 0$

Then \exists a scalar fn ϕ s.t. $A = \text{grad } \phi$

$\oint A \cdot d\vec{r}$
is
said
conservative.

$$A = \text{grad } \phi$$

Then the vector field is said to be conservative

V) Magnetostatics is also very important example of line integral of a loop of wire in magnetic field B carrying I . Then

$$dF = I dr \times B$$

So the entire force is defined for

$$F = I \oint_C dr \times B$$

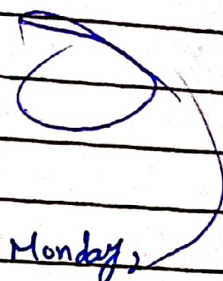


Heyl

~~about the wire~~

just

with

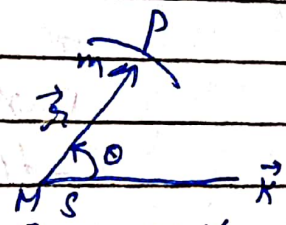


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Monday

Kepler's Problem:-

Let us consider a gravitational system in which particles of specified mass are moving in an inertial frame. Also if M be the mass of sun, k be the arbitrary fixed vector along an arbitrary reference line in Polar axis. Then the force acting on the planet by the sun is



$$F = - \left(\frac{GmM}{r^3} \right) \vec{r} \quad \text{--- (1)}$$

By Newton's II law for $F = ma$

$$\text{or } F = m \frac{dv}{dt} \quad \text{--- (2)}$$

From (1) and (2)

$$\frac{GMm}{r^3} \vec{r} = m \frac{dv}{dt}$$

$$\text{or } \frac{dv}{dt} = -\frac{GM}{r^3} \vec{r} \quad \text{--- (2)*}$$

Now,

$$\frac{d}{dt} (\vec{r} \times \vec{v}) = \vec{r} \times \frac{dv}{dt} + \frac{d\vec{r}}{dt} \times \vec{v} \quad \text{--- (3)}$$

$$\therefore \frac{d}{dt} (\vec{r} \times \vec{v}) = \vec{r} \times \frac{dv}{dt} \quad \text{--- (4)}$$

It is noted that 2nd term on the R.H.S of (3) vanishes because \parallel vectors so after using (2)* in (4)

$$\frac{d}{dt} (\vec{r} \times \vec{v}) = \vec{r} \times \left(-\frac{GM}{r^3} \vec{r} \right) = 0$$

and $\vec{r} \times \vec{v} = \vec{h} = \text{const. Vector}$
or,

$$\vec{r} \times \frac{d\vec{r}}{dt} = \vec{h} \quad \text{for } \vec{v} = \frac{d\vec{r}}{dt}$$

Since,

$|\vec{r} \times \frac{d\vec{r}}{dt}| = 2 \times \text{twice the area of sector.}$

That we get

$$2 \left| \frac{dA}{dt} \right| = |h| \quad (5)$$

This eq (5) is the Kepler's
2nd law of planetary motion.
If two vectors in scalar
triple product are identical
then the product is zero.

$$\mathbf{r} \cdot (\mathbf{r} \times \frac{d\mathbf{r}}{dt}) = \mathbf{r} \cdot \mathbf{h} = 0 \quad (6)$$

where (6) shows that $\mathbf{r} \perp \mathbf{h}$ and
the motion of planet takes
cross product of (6)

$$\frac{d\mathbf{r}}{dt} \times \mathbf{h} = - \frac{GM\mathbf{r}}{r^3} \times \mathbf{h}$$

$$= - \frac{GM\mathbf{r}}{r^3} \times (\mathbf{r} \times \mathbf{v}) \quad (7)$$

Now,

$$\mathbf{r} = r \hat{e}_r \quad \text{for } \hat{e}_r \text{ be unit vector}$$

so,

$$\mathbf{v} = \frac{d\mathbf{r}}{dt} = \frac{dr}{dt} \hat{e}_r + r \frac{d\hat{e}_r}{dt}$$

So (7) reduces to,

$$\begin{aligned} \frac{d}{dt} (\mathbf{v} \times \mathbf{h}) &= - \frac{GM}{r^3} \mathbf{r} \times \left(\mathbf{r} \times r \frac{d\hat{e}_r}{dt} \right) \\ &= - GM \hat{e}_r \times \left(\hat{e}_r \times r \frac{d\hat{e}_r}{dt} \right) \end{aligned}$$

or,

$$\frac{d}{dt} (\mathbf{v} \times \mathbf{h}) = - GM \left[\left(\hat{e}_r \cdot \frac{d\hat{e}_r}{dt} \right) \mathbf{r} - \left(\hat{e}_r \cdot \hat{e}_r \right) r \frac{d\hat{e}_r}{dt} \right]$$

Using the vector identities, we have,

$$\frac{d}{dt}(v \times h) = GM \frac{de_r}{dt}, \text{ or integrating gives}$$

$$\int \frac{d}{dt}(v \times h) = \int GM \frac{de_r}{dt}$$

$$\Rightarrow v \times h = GM e_r$$

if

$$r \cdot (v \times h) = r \cdot (GM e_r + K)$$

Then implies for

$$r \cdot (v \times h) = GM(r \cdot e_r) + r \cdot K$$

from the polar curves,

$$r \cdot (v \times h) = GM r + r K \cos \theta$$

where K is arbitrary const with specified magnitude, θ is angle b/w K and e_r . so using relation for 3-vectors in Ex. 9 gives,

$$r \cdot (v \times h) = (r \times v) \cdot h = h \cdot h = h^2$$

for $r \times v = h$.

$$\Rightarrow \boxed{r = \frac{h^2}{GM + K \cos \theta}} \rightarrow \textcircled{8}$$

This is taken as the Kepler's first law. It is responsible for the motion of the planets around sun.

to show that curl is elliptic in nature for a closed curve C .

Surface Integrals:-

Under the classical forces is taken for the double of integrated area about a closed curve path gives for,

$$(a) \int_S \phi ds \quad (b) \int_S A \cdot ds \quad (c) \int_S A \times ds$$

OR,

$$\oint_S \phi ds \quad \text{or} \quad \oint_S A \cdot ds \quad \text{or} \quad \oint_S A \times ds.$$