

Principal axis make an angle of  $\pi/4$  with coordinate axis i.e. principal axis in this case are along the diagonals of the square.

Now using,  $I_{xx} = A \cos^2 \theta + B \sin^2 \theta + F \sin 2\theta$

$$= \frac{5ma^2}{2} \left(\frac{1}{2}\right) + \frac{5ma^2}{2} \left(\frac{1}{2}\right) + \frac{ma^2}{2} \times (1)$$

$$= \frac{12ma^2}{4} \quad (\text{by lem})$$

Similarly  $= 3ma^2$

$$I_{yy} = A \sin^2 \theta + B \cos^2 \theta - F \sin 2\theta$$

$$= 2ma^2$$

$$I_{zz} = I_{xx} + I_{yy} = 3ma^2 + 2ma^2 = 5ma^2$$

Equimomental Systems

Two systems are said to be equimomental systems if they have same (equal) moments of inertia about every line in space.

Theorems

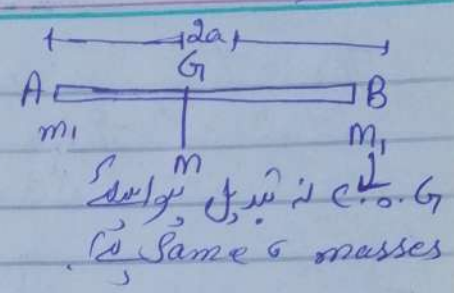
- Two systems are equimomental, if
- i) they have same total mass,
  - ii) They have same principal axis
  - iii) They have same centroid.

Q<sub>5</sub>

Show that a uniform rod of mass  $m$  is equimomental to three particles of masses  $\frac{m}{6}$ ,  $\frac{m}{6}$  &  $\frac{2m}{3}$  attached first & 2<sup>nd</sup> at its ends & 3<sup>rd</sup> at its centre?

Sol:

Consider a rod AB of length  $2a$  & mass  $M$ . Then if  $m$  is at a line  $\perp$  to the rod



$$AB = \frac{1}{3} Ma^2 \rightarrow \text{①}$$

Now, Consider three masses  $m_1, m_1, m_2$  attached at A, B & G resp.

$$\begin{aligned} \& m_1 + m_2 + m_1 &= M \\ m_2 &= M - 2m_1 \end{aligned}$$

Now, m.g of these three masses,

$$\begin{aligned} I &= \sum_{i=1}^3 m_i d_i^2 \\ &= m_1 d_1^2 + m_2 d_2^2 + m_3 d_3^2 \\ &= m_1 a^2 + m_1 a^2 + m_2 (0) \Rightarrow 2m_1 a^2 \end{aligned}$$

The system would be equimomental if  $2m_1 a^2 = \frac{1}{3} Ma^2$ .

$$\begin{aligned} m_1 &= \frac{1}{6} M \\ \text{for } m_2 &= M - 2m_1 \end{aligned}$$

i.e for equimomental system masses are,  $\frac{m}{6}, \frac{m}{6}$  &  $\frac{2m}{6}$ .

Note  $I = \frac{m}{6} \cdot a^2 + \frac{m}{6} a^2 + \frac{2m}{6} \times 0$

Theorems

The two given systems are equimomental if they have the same total mass ii) Same p. axis iii) same Centroid

Proof

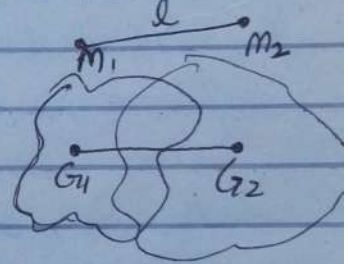
Consider  $m_1$  &  $m_2$  are two masses with  $G_1$  &  $G_2$  their C.O.M (centroid)

Join  $G_1$  &  $G_2$  as the two systems are equimomental. i.e They have same M.O about any line in space. we consider M.O of the two systems abt the line  $G_1, G_2$  as I.

Draw a line  $l \parallel$  to  $G_1, G_2$  then by using parallel axis theorem

M.O of 1st system abt  $l = I + m_1 h^2$

M.O of 2nd system abt  $l = I + m_2 h^2$



\* Systems are equimomental.

$$I + m_1 h^2 = I + m_2 h^2$$

$$m_1 h^2 = m_2 h^2$$

$$m_1 = m_2 = M \text{ (Say)}$$

Consider two  $\parallel$  lines  $G_1 H_1$  &  $G_2 H_2$  & each of the line,  $G_1 H_1$  &  $G_2 H_2$  is  $\perp$  to  $G_1 G_2$  <sup>also</sup> Then

M.O of 1st system abt line  $G_1 H_1 = I + M(G_1 G_2)^2$

" " 2nd " " = I + 0

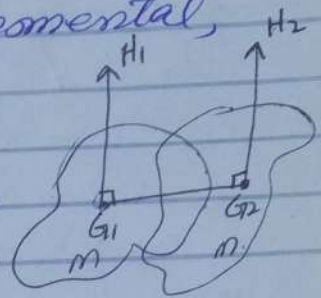
using Conditions of equimomental,

$$I + M |G_{11} G_{22}|^2 = I$$

$$M |G_{11} G_{22}|^2 = 0$$

$$M \neq 0; \Rightarrow |G_{11} G_{22}|^2 = 0$$

$$G_{11} = G_{22} = G \text{ (say)}$$

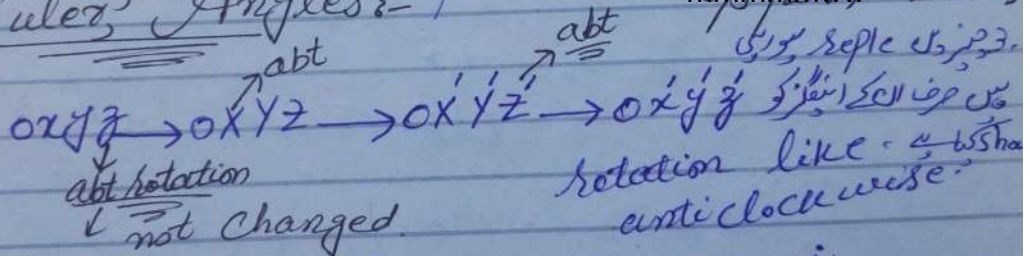


also, let the two systems have same m.o I about  $G_2 H$

∴ the two systems are equimom. i.e they have same m.o about any line, so they have same m.o about principal axis, so there. p. axis are same.

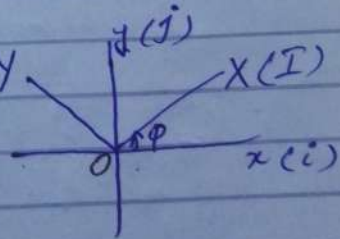
after mid Phewed

Euler's Angles:-

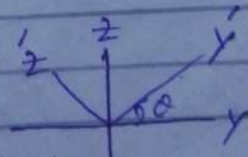


( $\phi, \theta, \psi$ )

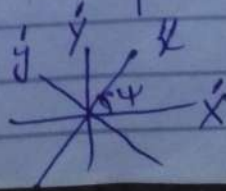
1st rotation:- ( $\phi$ )  
precession angle



2nd rotation:- ( $\theta$ )  
nutation angle



3rd rotation:- ( $\psi$ )  
body angle



Motion of a rigid body in space can be described by three angular displacements  $(\phi, \theta, \psi)$ , where  $\phi$  of the coordinate system  $OXYZ$  to  $O'X'Y'Z'$ ;  $O'X'Y'Z'$  to  $O''X''Y''Z''$  &  $O''X''Y''Z''$  to  $O'X'Y'Z'$  respectively. Where, 1st rotation is about  $Z$ -axis, 2nd rotation is about  $X$ -axis, & 3rd rotation is about  $Z'$ -axis.

Let,  $i, j, k; I, J, K; I', J', K'$  are the unit vectors along the axis in the respective fixed coordinate system.

for first rotation about  $Z$ -axis through an angle  $\phi$ .

$$i = (i \cdot I)I + (i \cdot J)J + (i \cdot K)K$$

This relation is from,

$$\vec{A} = A_1 \hat{i} + A_2 \hat{j} + A_3 \hat{k}$$

$$\vec{A} = (\vec{A} \cdot \hat{i})\hat{i} + (\vec{A} \cdot \hat{j})\hat{j} + (\vec{A} \cdot \hat{k})\hat{k}$$

$$\begin{cases} \because A \cdot i = A_1 \\ \because A \cdot j = A_2 \\ \because A \cdot k = A_3 \end{cases}$$

$\rightarrow$  Anticlockwise  $\rightarrow$   $i, j, k$

$$\begin{aligned} i &= \cos \phi I + \cos(90 + \phi)J + \cos 90 K \\ &= \cos \phi I - \sin \phi J \rightarrow (1) \end{aligned}$$

$$\begin{aligned} j &= (j \cdot I)I + (j \cdot J)J + (j \cdot K)K \\ &= \cos(90 - \phi)I + \cos \phi J + \cos 90 K \\ &= \sin \phi I + \cos \phi J \rightarrow (2) \end{aligned}$$

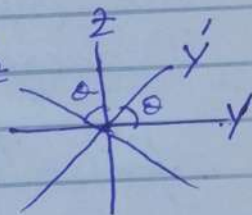
$$\begin{aligned} k &= (k \cdot I)I + (k \cdot J)J + (k \cdot K)K \\ &= \cos 90 I + \cos 90 J + \cos 0 K \\ &= K \rightarrow (3) \end{aligned}$$

So, first transformation eqs are given by (1), (2) & (3)

Now

ii) for 2<sup>nd</sup> rotation abt x-axis, through an angle  $\theta$ , we have,

$$\begin{aligned} \hat{i} &= (\hat{i} \cdot \hat{i}')\hat{i}' + (\hat{i} \cdot \hat{j}')\hat{j}' + (\hat{i} \cdot \hat{k}')\hat{k}' \\ &= \cos(0)\hat{i}' + \cos 90\hat{j}' + \cos 90\hat{k}' \\ &= \hat{i}' \rightarrow (4) \end{aligned}$$



Now

$$\begin{aligned} \hat{j} &= (\hat{j} \cdot \hat{i}')\hat{i}' + (\hat{j} \cdot \hat{j}')\hat{j}' + (\hat{j} \cdot \hat{k}')\hat{k}' \\ &= \cos(90)\hat{i}' + \cos \theta \hat{j}' + \cos(90+\theta)\hat{k}' \\ &= \cos \theta \hat{j}' - \sin \theta \hat{k}' \rightarrow (5) \end{aligned}$$

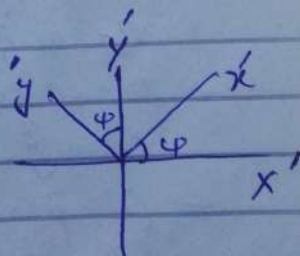
$$\begin{aligned} \hat{k} &= (\hat{k} \cdot \hat{i}')\hat{i}' + (\hat{k} \cdot \hat{j}')\hat{j}' + (\hat{k} \cdot \hat{k}')\hat{k}' \\ &= \cos(90)\hat{i}' + \cos(90-\theta)\hat{j}' + \cos \theta \hat{k}' \\ &= \sin \theta \hat{j}' + \cos \theta \hat{k}' \rightarrow (6) \end{aligned}$$

So, 2<sup>nd</sup> transformation eqs are given by (4), (5) & (6)

Now

iii) for 3<sup>rd</sup> rotation abt z-axis through an angle  $\psi$ ,

$$\begin{aligned} \hat{i}' &= (\hat{i}' \cdot \hat{i})\hat{i} + (\hat{i}' \cdot \hat{j})\hat{j} + (\hat{i}' \cdot \hat{k})\hat{k} \\ &= \cos \psi \hat{i} + \cos(90+\psi)\hat{j} + \cos 90\hat{k} \\ &= \cos \psi \hat{i} - \sin \psi \hat{j} \rightarrow (7) \end{aligned}$$



$$\begin{aligned} \hat{j}' &= (\hat{j}' \cdot \hat{i})\hat{i} + (\hat{j}' \cdot \hat{j})\hat{j} + (\hat{j}' \cdot \hat{k})\hat{k} \\ &= \cos(90-\psi)\hat{i} + \cos \psi \hat{j} + \cos 90\hat{k} \\ &= \sin \psi \hat{i} + \cos \psi \hat{j} \rightarrow (8) \end{aligned}$$

$$\begin{aligned} \hat{k}' &= (\hat{k}' \cdot \hat{i}')\hat{i}' + (\hat{k}' \cdot \hat{j}')\hat{j}' + (\hat{k}' \cdot \hat{k}')\hat{k}' \\ &= \cos 90^\circ \hat{i}' + \cos 90^\circ \hat{j}' + \cos(0)\hat{k}' \\ &= \hat{k}' \rightarrow \textcircled{9} \end{aligned}$$

So,

Transformation eqn<sup>s</sup> for 3<sup>rd</sup> rotation are given by  $\textcircled{7}$ ,  $\textcircled{8}$  &  $\textcircled{9}$ .

Deductions:

find Transformation eqn<sup>s</sup> from  $\hat{i}, \hat{j}, \hat{k}$  into  $\hat{i}', \hat{j}', \hat{k}'$

$$\begin{aligned} \hat{i} &= \cos \phi \hat{I} - \sin \phi \hat{J} \\ &= \cos \phi \hat{I}' - \sin \phi (\cos \theta \hat{J}' - \sin \theta \hat{K}') \\ &= \cos \phi \hat{I}' - \sin \phi \cos \theta \hat{J}' + \sin \phi \sin \theta \hat{K}' \\ &= \cos \phi (\cos \psi \hat{i} - \sin \psi \hat{j}) - \sin \phi \cos \theta (\sin \psi \hat{i} + \cos \psi \hat{j}) + \sin \phi \sin \theta \hat{k} \\ &= (\cos \phi \cos \psi - \sin \phi \sin \psi \cos \theta) \hat{i} - (\cos \phi \sin \psi + \sin \phi \cos \psi \cos \theta) \hat{j} + \sin \phi \sin \theta \hat{k} \\ &\rightarrow \textcircled{A} \end{aligned}$$

Now

$$\begin{aligned} \hat{j} &= \sin \phi \hat{I} + \cos \phi \hat{J} \\ &= \sin \phi \hat{I}' + \cos \phi (\cos \theta \hat{J}' - \sin \theta \hat{K}') \\ &= \sin \phi (\cos \psi \hat{i} - \sin \psi \hat{j}) + \cos \phi (\cos \theta \sin \psi \hat{i} + \cos \theta \cos \psi \hat{j} - \sin \theta \hat{k}) \\ &= \sin \phi \cos \psi \hat{i} - \sin \phi \sin \psi \hat{j} + \cos \phi \cos \theta \sin \psi \hat{i} + \cos \phi \cos \theta \cos \psi \hat{j} - \cos \phi \sin \theta \hat{k} \\ &= (\sin \phi \cos \psi + \cos \phi \sin \psi \cos \theta) \hat{i} + (\cos \phi \cos \theta \cos \psi - \sin \phi \sin \psi) \hat{j} - \cos \phi \sin \theta \hat{k} \rightarrow \textcircled{B} \end{aligned}$$

Similarly,

$$\hat{k} = K$$

$$= \sin\theta \hat{j} + \cos\theta \hat{k}'$$

$$= \sin\theta (\sin\psi \hat{i} + \cos\psi \hat{j}) + \cos\theta \hat{k}'$$

$$= \sin\theta \sin\psi \hat{i} + \sin\theta \cos\psi \hat{j} + \cos\theta \hat{k}' \rightarrow \textcircled{C}$$

(A), (B), (C) represents required Transformation

Alternatively,

first rotation can be written as,

$$\begin{bmatrix} \hat{i} \\ \hat{j} \\ \hat{k}' \end{bmatrix} = \begin{bmatrix} \cos\phi & -\sin\phi & 0 \\ \sin\phi & \cos\phi & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} I \\ J \\ K \end{bmatrix}$$

$$\bar{x} = R_\phi \bar{X} \rightarrow \textcircled{A}$$

2<sup>nd</sup> rotation.

$$I = I'$$

$$J = \cos\theta J' - \sin\theta K'$$

$$K = \sin\theta J' + \cos\theta K'$$

$$\begin{bmatrix} I \\ J \\ K \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos\theta & -\sin\theta \\ 0 & \sin\theta & \cos\theta \end{bmatrix} \begin{bmatrix} I' \\ J' \\ K' \end{bmatrix}$$

$$\bar{X} = R_\theta \bar{X}' \rightarrow \textcircled{B}$$

3<sup>rd</sup> rotation

$$\begin{bmatrix} I' \\ J' \\ K' \end{bmatrix} = \begin{bmatrix} \cos\psi & -\sin\psi & 0 \\ \sin\psi & \cos\psi & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} \hat{i} \\ \hat{j}' \\ \hat{k}' \end{bmatrix}$$

$$\bar{X}' = R_\psi \bar{x} \rightarrow \textcircled{C}$$



from (1)

$$\begin{aligned} \bar{x} &= R_p \bar{x}_p \\ &= R_p R_o \bar{x}' \\ &= R_p R_o R_y \bar{x}' \end{aligned}$$

$$\bar{x} = R \bar{x}'$$

Where,  $R = R_p R_o R_y$ .

Derived Euler's Geometrical eqs

If  $\bar{\omega}'$  is angular velocity after the three rotations.  $\hat{i}, \hat{j}, \hat{k}$  are unit vectors in the original frame, and  $\hat{i}', \hat{j}', \hat{k}'$  are unit vectors in the rotated frame.

$$\begin{aligned} \bar{\omega}' &= \omega_p \hat{k} + \omega_o \hat{i}' + \omega_y \hat{k}' \\ &= \omega_p (\sin \theta \hat{j} + \cos \theta \hat{k}) + \omega_o (\cos \psi \hat{i} - \sin \psi \hat{j}) + \omega_y \hat{k}' \\ &= \omega_p [\sin \theta (\sin \psi \hat{i} + \cos \psi \hat{j}) + \cos \theta \hat{k}] + \omega_o \cos \psi \hat{i} \\ &\quad - \omega_o \sin \psi \hat{j} + \omega_y \hat{k}' \\ &= \omega_p \sin \theta \sin \psi \hat{i} + \omega_p \sin \theta \cos \psi \hat{j} + \omega_p \cos \theta \hat{k} + \omega_o \cos \psi \hat{i} \\ &\quad - \omega_o \sin \psi \hat{j} + \omega_y \hat{k}' \\ \bar{\omega}' &= (\omega_p \sin \theta \sin \psi + \omega_o \cos \psi) \hat{i} + (\omega_p \sin \theta \cos \psi - \omega_o \sin \psi) \hat{j} \\ &\quad + (\omega_p \cos \theta + \omega_y) \hat{k}' \end{aligned}$$

Where,

$$\begin{cases} \dot{\omega}_x = \omega_p \sin \theta \sin \psi + \omega_o \cos \psi = \dot{\phi} \sin \theta \sin \psi + \dot{\alpha} \cos \psi \\ \dot{\omega}_y = \omega_p \sin \theta \cos \psi - \omega_o \sin \psi = \dot{\phi} \sin \theta \cos \psi - \dot{\alpha} \sin \psi \\ \dot{\omega}_z = \omega_p \cos \theta + \omega_y = \dot{\phi} \cos \theta + \dot{\psi} \end{cases}$$

are known as Euler's geometrical eqs.

Rules, Dynamical eqs of motion

Let a rigid body is rotating with angular velocity  $\vec{\omega}$  abt a fixed point O, let  $OX, OY$  &  $OZ$  are principal axis, if  $L$  &  $[I_{ij}]$  are angular momentum & inertia matrix then,

$$[L] = [I][\vec{\omega}]$$

$$\begin{bmatrix} L_x \\ L_y \\ L_z \end{bmatrix} = \begin{bmatrix} I_{xx} & I_{xy} & I_{xz} \\ I_{yx} & I_{yy} & I_{yz} \\ I_{zx} & I_{zy} & I_{zz} \end{bmatrix} \begin{bmatrix} \omega_x \\ \omega_y \\ \omega_z \end{bmatrix}$$

∴ Abt principal axis, we have,  
 $I_{xy} = I_{yx} = I_{yz} = I_{zy} = I_{zx} = I_{xz} = 0$

(i.e product of inertia vanishes abt principal axis")

So,

$$\begin{bmatrix} L_x \\ L_y \\ L_z \end{bmatrix} = \begin{bmatrix} I_{xx} & 0 & 0 \\ 0 & I_{yy} & 0 \\ 0 & 0 & I_{zz} \end{bmatrix} \begin{bmatrix} \omega_x \\ \omega_y \\ \omega_z \end{bmatrix}$$

$$L_x = I_{xx} \omega_x \quad \text{or} \quad L_1 = I_1 \omega_1$$

$$L_y = I_{yy} \omega_y \quad \text{or} \quad L_2 = I_2 \omega_2$$

$$L_z = I_{zz} \omega_z \quad \text{or} \quad L_3 = I_3 \omega_3$$

So,

$$\vec{L} = L_1 \hat{i} + L_2 \hat{j} + L_3 \hat{k}$$

$$\vec{L} = I_1 \omega_1 \hat{i} + I_2 \omega_2 \hat{j} + I_3 \omega_3 \hat{k} \rightarrow \textcircled{1}$$

using rotating axis theorem

$$\left(\frac{d\bar{A}}{dt}\right)_f = \left(\frac{d\bar{A}}{dt}\right)_r + \bar{\omega} \times \bar{A}$$

So,

$$\left(\frac{d\bar{L}}{dt}\right)_f = \left(\frac{d\bar{L}}{dt}\right)_r + \bar{\omega} \times \bar{L} \rightarrow (2)$$

as,  $\underline{L} = \underline{r} \times \underline{p} \rightarrow$  linear Mom.

$$\left(\frac{d\bar{L}}{dt}\right)_f = \frac{d\bar{r}}{dt} \times \bar{p} + \bar{r} \times \frac{d\bar{p}}{dt}$$

$$= \bar{r} \times \bar{p} + \bar{r} \times m\bar{a}$$

$$= 0 + \bar{r} \times \bar{F}$$

$$= \underline{\tau} \text{ (say)} \rightarrow (3)$$

$$\left. \begin{array}{l} \text{if } \underline{p} = m\underline{v} \\ \frac{d}{dt}(m\underline{v}) \\ m \frac{d\underline{v}}{dt} \\ m\bar{a} \end{array} \right\}$$

$$\left(\frac{d\bar{L}}{dt}\right)_r = \frac{d}{dt} (I_1 \omega_1 \hat{i} + I_2 \omega_2 \hat{j} + I_3 \omega_3 \hat{k})$$

$$= I_1 \dot{\omega}_1 \hat{i} + I_2 \dot{\omega}_2 \hat{j} + I_3 \dot{\omega}_3 \hat{k} \rightarrow (4)$$

Now

$$\bar{\omega} \times \bar{L} = \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ \omega_1 & \omega_2 & \omega_3 \\ I_1 \omega_1 & I_2 \omega_2 & I_3 \omega_3 \end{vmatrix}$$

$$= \omega_2 \omega_3 (I_3 - I_2) \hat{i} + \omega_1 \omega_3 (I_1 - I_3) \hat{j} + \omega_1 \omega_2 (I_2 - I_1) \hat{k}$$

$\rightarrow (5)$

using values from,

(3), (4) & (5) into eq (2)

$$\bar{L} = I_1 \omega_1 \hat{i} + I_2 \omega_2 \hat{j} + I_3 \omega_3 \hat{k} + \omega_2 \omega_3 (I_3 - I_2) \hat{i} \\ + \omega_1 \omega_3 (I_1 - I_3) \hat{j} + \omega_1 \omega_2 (I_2 - I_1) \hat{k}$$

Now simplifying:

$$\left. \begin{aligned} \tau_1 &= I_1 \dot{\omega}_1 + \omega_2 \omega_3 (I_3 - I_2) \\ \tau_2 &= I_2 \dot{\omega}_2 + \omega_1 \omega_3 (I_1 - I_3) \\ \tau_3 &= I_3 \dot{\omega}_3 + \omega_1 \omega_2 (I_2 - I_1) \end{aligned} \right\} \rightarrow (6)$$

These are known as Euler dynamical eqns.

Deductions:

for torque free motion  $\Rightarrow \tau = 0$   
then Euler's eqns becomes,

$I_1 \dot{\omega}_1 + \omega_2 \omega_3 (I_3 - I_2) = 0 \rightarrow (7)$

$I_2 \dot{\omega}_2 + \omega_1 \omega_3 (I_1 - I_3) = 0 \rightarrow (8)$

$I_3 \dot{\omega}_3 + \omega_1 \omega_2 (I_2 - I_1) = 0 \rightarrow (9)$

xply eq (7) by  $\omega_1$ , (8) by  $\omega_2$ , (9) by  $\omega_3$

$I_1 \omega_1 \dot{\omega}_1 + \omega_1 \omega_2 \omega_3 (I_3 - I_2) = 0$

$I_2 \omega_2 \dot{\omega}_2 + \omega_1 \omega_2 \omega_3 (I_1 - I_3) = 0$

$I_3 \omega_3 \dot{\omega}_3 + \omega_1 \omega_2 \omega_3 (I_2 - I_1) = 0$

addif these, we get,

$I_1 \omega_1 \dot{\omega}_1 + I_2 \omega_2 \dot{\omega}_2 + I_3 \omega_3 \dot{\omega}_3 = 0.$

or  $\frac{1}{2} \frac{d}{dt} [I_1 \omega_1^2 + I_2 \omega_2^2 + I_3 \omega_3^2] = 0$

intef.

$\frac{1}{2} [I_1 \omega_1^2 + I_2 \omega_2^2 + I_3 \omega_3^2] = \text{const.}$

$\left( \begin{aligned} &= \frac{1}{2} \frac{d}{dt} I \omega^2 \\ &= \frac{1}{2} \frac{d}{dt} I \omega^2 \end{aligned} \right)$

$\frac{1}{2} [I_1 \omega_1 \hat{i} + I_2 \omega_2 \hat{j} + I_3 \omega_3 \hat{k}] \cdot [\omega_1 \hat{i} + \omega_2 \hat{j} + \omega_3 \hat{k}] = C$

$\frac{1}{2} \underline{L} \cdot \underline{\omega} = C \rightarrow (10)$

This is known as energy theorem.

Q1: Show that for a torque free motion, magnitude of angular mom. is constant.

Sol:

for a torque free motion, Euler's dynamical eq<sup>s</sup> are,

$$I_1 \dot{\omega}_1 - \omega_2 \omega_3 (I_2 - I_3) = 0 \rightarrow (1)$$

$$I_2 \dot{\omega}_2 - \omega_1 \omega_3 (I_3 - I_1) = 0 \rightarrow (2)$$

$$I_3 \dot{\omega}_3 - \omega_1 \omega_2 (I_1 - I_2) = 0 \rightarrow (3)$$

xply (1)  $I_1 \omega_1$ , (2)  $I_2 \omega_2$  & (3)  $I_3 \omega_3$

$$I_1^2 \omega_1 \dot{\omega}_1 - \omega_1 \omega_2 \omega_3 (I_1 I_2 - I_1 I_3) = 0$$

$$I_2^2 \omega_2 \dot{\omega}_2 - \omega_1 \omega_2 \omega_3 (I_2 I_3 - I_2 I_1) = 0$$

$$I_3^2 \omega_3 \dot{\omega}_3 - \omega_1 \omega_2 \omega_3 (I_3 I_1 - I_3 I_2) = 0$$

adding all these,

$$I_1^2 \omega_1 \dot{\omega}_1 + I_2^2 \omega_2 \dot{\omega}_2 + I_3^2 \omega_3 \dot{\omega}_3 = 0$$

$$\frac{1}{2} \frac{d}{dt} (I_1^2 \omega_1^2 + I_2^2 \omega_2^2 + I_3^2 \omega_3^2) = 0$$

Integ.

$$\frac{1}{2} (I_1^2 \omega_1^2 + I_2^2 \omega_2^2 + I_3^2 \omega_3^2) = C^x$$

$$(I_1 \omega_1 \hat{i} + I_2 \omega_2 \hat{j} + I_3 \omega_3 \hat{k}) \cdot (I_1 \omega_1 \hat{i} + I_2 \omega_2 \hat{j} + I_3 \omega_3 \hat{k}) = C$$

$$L \cdot L = C$$

$$|L|^2 = C$$

$$|L| = C$$

$\Rightarrow$  mag. of A.M is const.

Q32. In case of Spherical Torque free motion, angular velocity is constant.

Soln

• for a torque free motion, Euler's dynamical eqs are,

$$I_1 \dot{\omega}_1 - \omega_2 \omega_3 (I_2 - I_3) = 0$$

$$I_2 \dot{\omega}_2 - \omega_1 \omega_3 (I_3 - I_1) = 0$$

$$I_3 \dot{\omega}_3 - \omega_1 \omega_2 (I_1 - I_2) = 0$$

for any Spherical Top.

$$I_1 = I_2 = I_3 = I \neq 0$$

$$\Rightarrow I_1 \dot{\omega}_1 = 0 ; I_2 \dot{\omega}_2 = 0, I_3 \dot{\omega}_3 = 0$$

integ. all,

$$\Rightarrow \omega_1 = C_1$$

$$\Rightarrow \omega_2 = C_2$$

$$\Rightarrow \omega_3 = C_3$$

Therefore:

$$\vec{\omega} = \omega_1 \hat{i} + \omega_2 \hat{j} + \omega_3 \hat{k} = \text{const.}$$

$\Rightarrow$  Angular vel is const.

Q33

An ellipsoid torque free motion about its centre, set into motion at  $t=0$  &  $(\omega_1, \omega_2, \omega_3)$  as components of its angular velocity  $\vec{\omega}$  along the principal axis,  $6A > 3A$  &  $A$  are M.I about the axis,

find angular velocity at any time  $t$  & show that when  $t \rightarrow \infty$  the mag. of the angular velocity is  $n\sqrt{5}$

Sol:

for torque free motion, Euler's eq of motion are,  $I_1 \dot{\omega}_1 - \omega_2 \omega_3 (I_2 - I_3) = 0 \rightarrow \textcircled{1}$

$$I_2 \dot{\omega}_2 - (I_3 - I_1) \omega_1 \omega_3 = 0 \rightarrow \textcircled{2}$$

$$I_3 \dot{\omega}_3 - (I_1 - I_2) \omega_1 \omega_2 = 0 \rightarrow \textcircled{3}$$

using  $I_1 = 6A$ ,  $I_2 = 3A$ ,  $I_3 = A$

We have,

$$6A \dot{\omega}_1 - 2A \omega_2 \omega_3 = 0 \quad \text{by } \frac{1}{A}$$

$$3A \dot{\omega}_1 - A \omega_2 \omega_3 = 0 \rightarrow \textcircled{4}$$

$$3A \dot{\omega}_2 + 5A \omega_1 \omega_3 = 0 \rightarrow \textcircled{5}$$

$$A \dot{\omega}_3 - 3A \omega_1 \omega_2 = 0 \rightarrow \textcircled{6}$$

$$\textcircled{4} \text{ } \omega_1 \Rightarrow 15A \omega_1 \dot{\omega}_1 - 5A \omega_1 \omega_2 \omega_3 = 0$$

$$\textcircled{5} \text{ } \omega_2 \Rightarrow 3A \omega_2 \dot{\omega}_2 + 5A \omega_1 \omega_2 \omega_3 = 0$$

$$15A \omega_1 \dot{\omega}_1 + 3A \omega_2 \dot{\omega}_2 = 0$$

$$3A (5\omega_1 \dot{\omega}_1 + \omega_2 \dot{\omega}_2) = 0$$

$$5\omega_1 \dot{\omega}_1 + \omega_2 \dot{\omega}_2 = 0$$

Integ

$$\frac{5\omega_1^2}{2} + \frac{\omega_2^2}{2} = C_1$$

$$5\omega_1^2 + \omega_2^2 = 2C_1$$

$$(\because 2C_1 = e_1)$$

$$5\omega_1^2 + \omega_2^2 = C_1 \rightarrow \textcircled{7}$$

at  $t=0$ ,  $\omega_1 = n$ ,  $\omega_2 = 0$

$$5n^2 = C_1$$

$$5\omega_1^2 + \omega_2^2 = 5n^2$$

$$5\omega_1^2 = 5n^2 - \omega_2^2$$

$$\omega_1 = \sqrt{\frac{5n^2 - \omega_2^2}{5}} \rightarrow \textcircled{8}$$

from (4) & (6)

(4)  $3\omega_1$ , (6)  $\omega_3$ .

$$9A \dot{\omega}_1 \omega_1 - 3A \omega_1 \omega_2 \omega_3 = 0$$

$$-A \dot{\omega}_3 \omega_3 + 3A \omega_1 \omega_2 \omega_3 = 0$$

$$9A \dot{\omega}_1 \omega_1 - A \dot{\omega}_3 \omega_3 = 0$$

$$9 \dot{\omega}_1 \omega_1 - \dot{\omega}_3 \omega_3 = 0$$

integ.

$$9 \frac{\omega_1^2}{2} - \frac{\omega_3^2}{2} = C_2$$

$$9\omega_1^2 - \omega_3^2 = C_2$$

at  $t=0$ ,  $\omega_1 = n$ ,  $\omega_3 = 3n$ .

$$9n^2 - 9n^2 = C_2$$

$$\Rightarrow C_2 = 0$$

$$\Rightarrow 9\omega_1^2 = \omega_3^2$$

$$3\omega_1 = \omega_3$$

$$\omega_3 = 3 \sqrt{\frac{5n^2 - \omega_2^2}{5}} \rightarrow (10)$$

using these results, (8), (10) into (5)

$$3A \dot{\omega}_2 + 5A \omega_1 \omega_3 = 0$$

$$3A \dot{\omega}_2 + 5A \left( \sqrt{\frac{5n^2 - \omega_2^2}{5}} \right) \left( 3 \sqrt{\frac{5n^2 - \omega_2^2}{5}} \right) = 0$$

$$3A \dot{\omega}_2 + 5A \left( 3 \left( \frac{5n^2 - \omega_2^2}{5} \right) \right) = 0$$

$$3A \dot{\omega}_2 + 15A \left( \frac{5n^2 - \omega_2^2}{5} \right) = 0$$

$$15A \dot{\omega}_2 + 15A \left( 5n^2 - \omega_2^2 \right) = 0$$

5

$$15A \dot{\omega}_2 + 15A \left( 5n^2 - \omega_2^2 \right) = 0$$

$$\dot{\omega}_2 + \left( 5n^2 - \omega_2^2 \right) = 0$$



$$\dot{w}_2 - \dot{w}_2^2 + \sin^2 = 0$$

$$\frac{dw_2}{dt} = w_2^2 - \sin^2$$

$$dw_2 = (w_2^2 - \sin^2) dt$$

$$\frac{dw_2}{w_2^2 - \sin^2} = dt$$

$$\int \frac{dw_2}{w_2^2 - \sin^2} = \int dt$$

$$\int \frac{dw_2}{w_2^2 - \sin^2} = \int dt$$

$$-\int \frac{dw_2}{\sin^2 - w_2^2} = \int dt$$

$$-\frac{1}{2\sqrt{\sin^2}} \ln \frac{\sqrt{\sin^2} + w_2}{\sqrt{\sin^2} - w_2} = t + C_1$$

$$= -\frac{1}{2\sqrt{\sin^2}} \ln \frac{\sqrt{\sin^2} + w_2}{\sqrt{\sin^2} - w_2} = t + C_1$$

put  $t=0, w_2=0$

$$C_1 = -\frac{1}{2\sqrt{\sin^2}} \ln \frac{\sqrt{\sin^2}}{\sqrt{\sin^2}}$$

$$= -\frac{1}{2\sqrt{\sin^2}} \ln(1) \rightarrow 0$$

$$C_1 = 0$$

$$t = -\frac{1}{2\sqrt{\sin^2}} \ln \frac{\sqrt{\sin^2} + w_2}{\sqrt{\sin^2} - w_2}$$

$$\frac{1}{2\sqrt{\sin^2}} \ln \frac{\sqrt{\sin^2} + w_2}{\sqrt{\sin^2} - w_2} = -t$$

$$\frac{\ln \frac{\sqrt{\sin^2} + w_2}{\sqrt{\sin^2} - w_2}}{\sqrt{\sin^2} - w_2} = -2\sqrt{\sin^2} t$$

$$\frac{\sqrt{\sin^2} + w_2}{\sqrt{\sin^2} - w_2} = e^{-2\sqrt{\sin^2} t}$$

$$\frac{\sqrt{\sin^2} + w_2 + \sqrt{\sin^2} - w_2}{\sqrt{\sin^2} + w_2 - \sqrt{\sin^2} - w_2}$$

$$= \frac{e^{-2\sqrt{\sin^2} t} + 1}{e^{-2\sqrt{\sin^2} t} - 1}$$

$$\frac{2\sqrt{\sin^2}}{2w_2} = \frac{e^{-2\sqrt{\sin^2} t} + 1}{e^{-2\sqrt{\sin^2} t} - 1}$$

$$w_2 = \sqrt{\sin^2} \left[ \frac{e^{-2\sqrt{\sin^2} t} - 1}{e^{-2\sqrt{\sin^2} t} + 1} \right]$$

also

$$\int \frac{dw_2}{w_2^2 - \sin^2} = \int dt$$

$$-\int \frac{dw_2}{\sin^2 - w_2^2} = \int dt$$

$$-\int \frac{dw_2}{(\sin^2 - w_2^2)} = \int dt$$

$$-\frac{1}{\sqrt{\sin^2}} \tanh^{-1} \frac{w_2}{\sqrt{\sin^2}} = t$$

$$\tanh^{-1} \frac{w_2}{\sqrt{\sin^2}} = -\sqrt{\sin^2} t$$

$$\frac{w_2}{\sqrt{\sin^2}} = -\tanh \sqrt{\sin^2} t$$

$$w_2 = -\sqrt{\sin^2} \tanh \sqrt{\sin^2} t$$

$$\omega_2 = n\sqrt{5} \left[ \frac{e^{-\sqrt{5}nt}}{e^{\sqrt{5}nt}} \left( \frac{e^{-\sqrt{5}nt} - e^{\sqrt{5}nt}}{e^{-\sqrt{5}nt} + e^{\sqrt{5}nt}} \right) \right]$$

$$\omega_2 = n\sqrt{5} \tanh \sqrt{5}nt$$

$$\text{Now } \omega_1 = \sqrt{\frac{5n^2 - (-\sqrt{5}n \tanh \sqrt{5}nt)^2}{5}}$$

$$\rightarrow = \sqrt{\frac{5n^2 - 5n^2 \tanh^2 \sqrt{5}nt}{5}}$$

$$= n \sqrt{1 - \tanh^2 \sqrt{5}nt}$$

$$= n \operatorname{sech} \sqrt{5}nt$$

$$\omega_1 = n \operatorname{sech} \sqrt{5}nt$$

$$\Rightarrow \omega_3 = 3n \operatorname{sech} \sqrt{5}nt$$

$$\bar{\omega} = \omega_1 i + \omega_2 j + \omega_3 k$$

$$= n \operatorname{sech} \sqrt{5}nt i - \sqrt{5}n \tanh \sqrt{5}nt j + 3n \operatorname{sech} \sqrt{5}nt k$$

$$t \rightarrow \infty$$

$$\bar{\omega} = 0 - \sqrt{5}n(1) + 0$$

$$= -n\sqrt{5}$$

$$\bar{\omega} = \sqrt{5}n \text{ as required.}$$

$$\left( \begin{array}{l} \tanh \infty = 1 \\ \operatorname{sech} \infty = 0 \end{array} \right)$$

formulas,

Ans.

$$\int \frac{1}{a^2 - x^2} dx = \frac{1}{2a} \ln \frac{a+x}{a-x}$$

$$\int \frac{1}{x^2 - a^2} dx = \frac{1}{2a} \ln \frac{x-a}{x+a}$$

$$\text{or } \int \frac{1}{\sqrt{a^2 - x^2}} dx = \frac{1}{a} \tanh^{-1} \frac{x}{a} \text{ (using)}$$