3D Stress Components
From equilibrium principles:

$$
\tau_{\mathrm{xy}}=\tau_{\mathrm{yx}}, \tau_{\mathrm{xz}}=\tau_{\mathrm{zx}}, \tau_{\mathrm{zy}}=\tau_{\mathrm{yz}}
$$

The most general state of stress at a point may be represented by 6 components

Shear Stresses


Normal stress $(\sigma)$ : the subscript identifies the face on which the stress acts. Tension is positive and compression is negative. Shear stress ( $\tau$ ) : it has two subscripts. The first subscript denotes the face on which the stress acts. The second subscript denotes the direction on that face. A shear stress is positive if it acts on a positive face and positive direction or if it acts in a negative face and negative direction.


For static equilibrium $\tau_{x y}=\tau_{y x}, \tau_{x z}=\tau_{z x}, \tau_{z y}=\tau_{y z}$ resulting in ${ }_{s}$ ix independent scalar quantities. These six scalars can be arranged in a $3 \times 3$ matrix, giving us a stress tensor.
$\boldsymbol{\sigma}=\sigma_{i j}=\left[\begin{array}{ccc}\sigma_{x} & \tau_{y x} & \tau_{z x} \\ \tau_{x y} & \sigma_{y} & \tau_{z y} \\ \tau_{x z} & \tau_{y z} & \sigma_{z}\end{array}\right]$


The sign convention for the stress elements is that a positive force on a positive face or a negative force on a negative face is positive. All others are negative.

The stress state is a second order tensor since it is a quantity associated with two directions (two subscripts direction of the surface normal and direction of the stress).
Same state of stress is represented by a different set of components if axes are rotated. There is a special set of components (when axes are rotated) where all the shear components are zero (principal stresses).



A property of a symmetric tensor is that there exists an orthogonal set of axes 1,2 and 3 (called principal axes) with respect to which the tensor elements are all zero except for those in the diagonal.
$\boldsymbol{\sigma}=\sigma_{i j}=\left[\begin{array}{ccc}\sigma_{x} & \tau_{y x} & \tau_{z x} \\ \tau_{x y} & \sigma_{y} & \tau_{z y} \\ \tau_{x z} & \tau_{y z} & \sigma_{z}\end{array}\right] \quad \xrightarrow{\text { Eigen values }} \quad \boldsymbol{\sigma}^{\prime}=\sigma_{i j}^{\prime}=\left[\begin{array}{ccc}\sigma_{1} & 0 & 0 \\ 0 & \sigma_{2} & 0 \\ 0 & 0 & \sigma_{3}\end{array}\right]$
In matrix notation the transformation is known as the Eigen-values.
The principal stresses are the "new-axes" coordinate system. The angles between the "old-axes" and the "new-axes" are known as the Eigen-vectors.
"old" axes

|  | x1 | x 2 | x3 |
| :---: | :---: | :---: | :---: |
| x'1 | a11 | a12 | al3 |
| "new" x'2 | a21 | a22 | a23 |
| x'3 | a31 | a32 | a33 |
| principal stress | Cosine of angle between $X$ and the principal stress | Cosine of angle between Y and the principal stress | Cosine of angle between Z and the principal stress |
| $\sigma_{1}$ | k1 | 11 | m1 |
| $\sigma_{2}$ | k2 | 12 | m2 |
| $\sigma_{3}$ | k3 | 13 | m3 |

## Plane Stress

State of stress in which two faces of the cubic element are free of stress. For the illustrated example, the state of stress is defined by


$$
\sigma_{x}, \sigma_{y}, \tau_{\mathrm{xy}} \quad \text { and } \quad \sigma_{z}=\tau_{z x}=\tau_{z y}=0
$$

## Sign Conventions for Shear Stress and Strain



The Shear Stress will be considered positive when a pair of shear stress acting on opposite sides of the element produce a counterclockwise (ccw) torque (couple).


A shear strain in an element is positive when the angle between two positive faces (or two negative faces) is reduced, and is negative if the angle is increased.


$$
\left[\begin{array}{ccc}
\sigma_{x 1 x 1} & \tau_{y 1 x 1} & 0 \\
\tau_{x 1 y 1} & \sigma_{y 1 y 1} & 0 \\
0 & 0 & 0
\end{array}\right]
$$

$$
\left[\begin{array}{ccc}
\sigma_{1} & 0 & 0 \\
0 & \sigma_{2} & 0 \\
0 & 0 & 0
\end{array}\right]
$$

## Stresses on Inclined Sections

Knowing the normal and shear stresses acting in the element denoted by the $x y$ axes, we will calculate the normal and shear stresses acting in the element denoted by the axis $\boldsymbol{x}_{1} y_{1}$.



Equilibrium of forces:
Acting in $\boldsymbol{x}_{1}$

$$
\sigma_{X 1} \cdot \frac{A_{O}}{\cos \theta}=\sigma_{X} \cos \theta \cdot A_{O}+\tau_{X Y} \sin \theta \cdot A_{O}+\sigma_{Y} \operatorname{Sin} \theta \cdot A_{O} \frac{\operatorname{Sin} \theta}{\operatorname{Cos} \theta}+\tau_{Y X} \cos \theta \cdot A_{O} \frac{\operatorname{Sin} \theta}{\operatorname{Cos} \theta}
$$


$\tau_{x 1 y} A_{o} \sec \theta=-\sigma_{x} A_{o} \sin \theta+\tau_{x y} A_{o} \cos \theta+\sigma_{y} A_{o} \tan \theta \cos \theta-\tau_{y x} A_{o} \tan \theta \sin \theta$
Eliminating $A_{o}, \boldsymbol{\operatorname { s e c }} \theta=\mathbf{1} / \boldsymbol{\operatorname { c o s }} \theta$ and $\tau_{x y}=\tau_{y x}$
$\tau_{x \mid y 1}=-\sigma_{x} \cdot \sin \theta \cdot \cos \theta+\sigma_{y} \cdot \sin \theta \cdot \cos \theta+\tau_{x y}\left(\cos ^{2} \theta-\sin ^{2} \theta\right)$

## Transformation Equations for Plane Stress

Using the following trigonometric identities:
$\operatorname{Cos}^{2} \theta=1 / 2(1+\cos 2 \theta)$
$\sin ^{2} \theta=1 / 2(1-\cos 2 \theta)$
$\sin \theta \cos \theta=1 / 2 \sin 2 \theta$

$$
\begin{aligned}
& \sigma_{x 1}=\frac{\sigma_{x}+\sigma_{y}}{2}+\frac{\sigma_{x}-\sigma_{y}}{2} \cos 2 \theta+\tau_{x y} \sin 2 \theta \\
& \sigma_{y 1}=\frac{\sigma_{x}+\sigma_{y}}{2}-\frac{\sigma_{x}-\sigma_{y}}{2} \cos 2 \theta-\tau_{x y} \sin 2 \theta \\
& \tau_{x 1 y 1}=-\frac{\sigma_{x}-\sigma_{y}}{2} \sin 2 \theta+\tau_{x y} \cos 2 \theta
\end{aligned}
$$

These equations are known as the transformation equations for plane stress.

Special Cases $\quad \sigma_{x 1}=\frac{\sigma_{x}+\sigma_{y}}{2}+\frac{\sigma_{x}-\sigma_{y}}{2} \cos 2 \theta+\tau_{x y} \sin 2 \theta$
Case 1: Uniaxial stress $\tau_{x \mid y 1}=-\frac{\sigma_{x}-\sigma_{y}}{2} \sin 2 \theta+\tau_{x y} \cos 2 \theta$
$\sigma_{y}=0 \quad \tau_{\mathrm{xy}}=\tau_{y x}=0$
$\sigma_{x 1}=\sigma_{x} \cdot\left(\frac{1+\operatorname{Cos} 2 \theta}{2}\right)$
$\tau_{x 1, y 1}=-\sigma_{x} \cdot\left(\frac{\operatorname{Sin} 2 \theta}{2}\right)$


Case 2 : Pure Shear
$\sigma_{x}=\sigma_{y}=0$
$\sigma_{x 1}=\tau_{x y} \cdot \operatorname{Sin} 2 \theta$
$\tau_{x 1, y 1}=\tau_{x y} \cdot \operatorname{Cos} 2 \theta$


## Case 3: Biaxial stress

$$
\begin{aligned}
& \tau_{x y}=0 \\
& \sigma_{x 1}=\frac{\sigma_{x}+\sigma_{y}}{2}+\frac{\sigma_{x}-\sigma_{y}}{2} \cdot \operatorname{Cos} 2 \theta \\
& \tau_{x 1, y 1}=-\frac{\sigma_{x}-\sigma_{y}}{2} \cdot \operatorname{Sin} 2 \theta
\end{aligned}
$$



## Example

An element in plane stress is subjected to stresses $\sigma_{x}=16000 p s i, \sigma_{y}=6000 p s i$, and $\tau_{x y}=\tau_{y x}=4000$ psi (as shown in figure below). Determine the stresses acting on an element inclined at an angle $\theta=45^{\circ}$ (counterclockwise - ccw).


$$
\begin{aligned}
& \sigma_{x 1}=\frac{\sigma_{x}+\sigma_{y}}{2}+\frac{\sigma_{x}-\sigma_{y}}{2} \cos 2 \theta+\tau_{x y} \sin 2 \theta \\
& \sigma_{y 1}=\frac{\sigma_{x}+\sigma_{y}}{2}-\frac{\sigma_{x}-\sigma_{y}}{2} \cos 2 \theta-\tau_{x y} \sin 2 \theta \\
& \tau_{x 1 y 1}=-\frac{\sigma_{x}-\sigma_{y}}{2} \sin 2 \theta+\tau_{x y} \cos 2 \theta
\end{aligned}
$$

## Numerical substitution

For $x$-axis $\theta=+45^{\circ}$ (ccw)
$\operatorname{Sin}(2 \theta)=\operatorname{Sin}\left(90^{\circ}\right)=1$
$\operatorname{Cos}(2 \theta)=\operatorname{Cos}\left(90^{\circ}\right)=0$
For $y$-axis $\theta=+45^{\circ}+90^{\circ}$ (ccw)
$\operatorname{Sin}(2 \theta)=\operatorname{Sin}\left(270^{\circ}\right)=-1$
$\operatorname{Cos}(2 \theta)=\operatorname{Cos}\left(270^{\circ}\right)=0$


$$
\begin{aligned}
& \frac{\sigma_{x}+\sigma_{y}}{2}=\frac{(16000+6000)}{2}=11000 p s i \\
& \frac{\sigma_{x}-\sigma_{y}}{2}=\frac{(16000-6000)}{2}=5000 p s i \\
& \tau_{x y}=4000 p s i \\
& \sigma_{x 1}=11000+5000(0)+4000(1)=15000 p s i \\
& \sigma_{y 1}=11000-5000(0)-4000(1)=7000 p s i \\
& \tau_{x 1 y 1}=-5000(1)+4000(0)=-5000 p s i
\end{aligned}
$$

## Example

A plane stress condition exists at a point on the surface of a loaded structure such as shown below. Determine the stresses acting on an element that is oriented at a clockwise (cw) angle of $\mathbf{1 5}{ }^{\circ}$ with respect to the original element, the principal stresses, the maximum shear stress and the angle of inclination for the principal stresses


Solution: We will use the following transformation equations:

$$
\begin{aligned}
& \sigma_{x 1}=\frac{\sigma_{x}+\sigma_{y}}{2}+\frac{\sigma_{x}-\sigma_{y}}{2} \cos 2 \theta+\tau_{x y} \sin 2 \theta \\
& \tau_{x 1 y 1}=-\frac{\sigma_{x}-\sigma_{y}}{2} \sin 2 \theta+\tau_{x y} \cos 2 \theta
\end{aligned}
$$

$\frac{\sigma_{x}+\sigma_{y}}{2}=\frac{(-46+12)}{2}=-17 \mathrm{MPa}$
$\frac{\sigma_{x}-\sigma_{y}}{2}=\frac{(-46-12)}{2}=-29 M P a$
$\tau_{x y}=-19 M P a$
$\sigma_{x 1}=(-17)+(-29)(0.866)+(-19)(-0.5)=-32.6 M P a$
$\tau_{x 1 y 1}=(-29)(-0.5)+(-19)(0.866)=-31 M P a$

$$
\begin{aligned}
& \text { For } x \text {-axis } \theta=-15^{\circ}(c w) \\
& \operatorname{Sin}(2 \theta)=\operatorname{Sin}\left(-30^{\circ}\right)=-0.5 \\
& \operatorname{Cos}(2 \theta)=\operatorname{Cos}\left(-30^{\circ}\right)=0.866
\end{aligned}
$$

$$
\begin{aligned}
& \sigma_{x}+\sigma_{y}=\sigma_{x 1}+\sigma_{y 1} \\
& \sigma_{y 1}=-1.4 M P a
\end{aligned}
$$



## Example

A rectangular plate of dimensions 3.0 in $\times 5.0$ in is formed by welding two triangular plates (see figure). The plate is subjected to a tensile stress of 600 psi in the long direction and a compressive stress of 250psi in the short direction.
Determine the normal stress sw acting perpendicular to the line or the weld and the shear stress $t w$ acting parallel to the weld. (Assume sw is positive when it acts in tension and tw is positive when it acts counterclockwise against the weld).


$$
\sigma_{x 1}=375 p s i \quad \tau_{x 1 y 1}=-375 p s i
$$

$$
\sigma_{y 1}=-25 p s i
$$

For $x$-axis $\tan \theta=\frac{3}{5} \Rightarrow$
$\theta=30.96^{\circ} \Rightarrow 2 \theta=61.92^{0}$

## Stresses acting on the weld

$$
\begin{aligned}
& \sigma_{x}+\sigma_{y}=\sigma_{x 1}+\sigma_{y 1} \\
& \sigma_{y 1}=-25 p s i
\end{aligned}
$$

$$
\frac{\sigma_{x}+\sigma_{y}}{2}=175 p s i \quad \frac{\sigma_{x}-\sigma_{y}}{2}=425 p s i \quad \tau_{x y}=0 p s i
$$

600 psi




$$
\sigma_{w}=-25 p s i \text { and } \tau_{w}=375 p s i
$$

## Principal Stresses and Maximum Shear Stresses

The sum of the normal stresses acting on perpendicular faces of plane stress elements is constant and independent of the angle $\theta$.

$$
\begin{gathered}
\sigma_{X 1}=\sigma_{X} \cos ^{2} \theta+\sigma_{Y} \sin ^{2} \theta+2 \tau_{X Y} \sin \theta \cos \theta \\
\sigma_{Y 1}=\sigma_{X} \sin ^{2} \theta+\sigma_{Y} \cos ^{2} \theta-2 \tau_{X Y} \sin \theta \cos \theta \\
\sigma_{X 1}+\sigma_{Y 1}=\sigma_{X}+\sigma_{Y}
\end{gathered}
$$

As we change the angle $\theta$ there will be maximum and minimum normal and shear stresses that are needed for design purposes.

## The maximum and minimum normal

 stresses are known as the principal stresses. These stresses are found by taking the derivative of $\sigma_{x 1}$ with respect to $\theta$ and setting equal to zero.$$
\begin{aligned}
& \sigma_{x 1}=\frac{\sigma_{x}+\sigma_{y}}{2}+\frac{\sigma_{x}-\sigma_{y}}{2} \cos 2 \theta+\tau_{x y} \sin 2 \theta \\
& \frac{\delta \sigma_{x 1}}{\delta \theta}=-\left(\sigma_{x}-\sigma_{y}\right) \sin 2 \theta+2 \tau_{x y} \cos 2 \theta=0 \\
& \tan 2 \theta_{P}=\frac{\tau_{x y}}{\left(\frac{\sigma_{x}-\sigma_{y}}{2}\right)}
\end{aligned}
$$

The subscript $\boldsymbol{p}$ indicates that the angle $\theta_{p}$ defines the orientation of the principal planes. The angle $\theta_{p}$ has two values that differ by $90^{\circ}$. They are known as the principal angles. For one of these angles $\sigma_{x 1}$ is a maximum principal stress and for the other a minimum. The principal stresses occur in mutually perpendicular planes.


$$
\tan 2 \theta_{P}=\frac{\tau_{x y}}{\left(\sigma_{x}-\sigma_{y}\right) / 2}
$$

$$
\tau_{x y} \sin 2 \theta_{P}=\frac{\tau_{x y}}{R} \quad \cos 2 \theta_{P}=\frac{\left(\sigma_{x}-\sigma_{y}\right)}{2 R}
$$

$$
\frac{\sigma_{x}-\sigma_{y}}{2}
$$

$$
\sigma_{x 1}=\frac{\sigma_{x}+\sigma_{y}}{2}+\frac{\sigma_{x}-\sigma_{y}}{2} \cos 2 \theta+\tau_{x y} \sin 2 \theta
$$

for the maximum stress $\theta=\theta_{\mathrm{P}}$
$\sigma_{1}=\frac{\sigma_{x}+\sigma_{y}}{2}+\frac{\sigma_{x}-\sigma_{y}}{2}\left(\frac{\sigma_{x}-\sigma_{y}}{2 R}\right)+\tau_{x y}\left(\frac{\tau_{x y}}{R}\right)=\frac{\sigma_{x}+\sigma_{y}}{2}+\left(\frac{\sigma_{x}-\sigma_{y}}{2}\right)^{2}\left(\frac{1}{R}\right)+\left(\tau_{x y}\right)^{2}\left(\frac{1}{R}\right)$
$\sigma_{1}=\frac{\sigma_{x}+\sigma_{y}}{2}+\left(\frac{1}{R}\right)\left[\left(\frac{\sigma_{x}-\sigma_{y}}{2}\right)^{2}+\left(\tau_{x y}\right)^{2}\right]$
But $\quad R=\sqrt{\left(\frac{\sigma_{x}-\sigma_{y}}{2}\right)^{2}+\left(\tau_{x y}\right)^{2}}$

$$
\sigma_{1}=\frac{\sigma_{x}+\sigma_{y}}{2}+\sqrt{\left(\frac{\sigma_{x}-\sigma_{y}}{2}\right)^{2}+\left(\tau_{x y}\right)^{2}}
$$

## Principal Stresses

$\sigma_{1}=\left(\frac{\sigma_{x}+\sigma_{y}}{2}\right)+\sqrt{\left(\frac{\sigma_{x}-\sigma_{y}}{2}\right)^{2}+\left(\tau_{x y}\right)^{2}}=\sigma_{\text {Average }}+R$
$\sigma_{2}=\left(\frac{\sigma_{x}+\sigma_{y}}{2}\right)-\sqrt{\left(\frac{\sigma_{x}-\sigma_{y}}{2}\right)^{2}+\left(\tau_{x y}\right)^{2}}=\sigma_{\text {Average }}-R$
The plus sign gives the algebraically larger principal stress and the minus sign the algebraically smaller principal stress.

This are the in-plane principal stresses. The third stress is zero in plane stress conditions

## Maximum Shear Stresses

The location of the angle for the maximum shear stress is obtained by taking the derivative of $\tau_{x 1 y 1}$ with respect to $\theta$ and setting it equal to zero.

$$
\tau_{x 1 y 1}=-\frac{\sigma_{x}-\sigma_{y}}{2} \sin 2 \theta+\tau_{x y} \cos 2 \theta
$$

$$
\begin{aligned}
& \frac{\delta \tau_{x 1 y 1}}{\delta \theta}=-\left(\sigma_{x}-\sigma_{y}\right) \cos 2 \theta-2 \tau_{x y} \sin 2 \theta=0 \\
& \tan 2 \theta_{S}=-\frac{\left(\sigma_{x}-\sigma_{y}\right) / 2}{\tau_{x y}}
\end{aligned}
$$

$$
\cos 2 \theta_{s 1}=\frac{\tau_{x y}}{R} \quad \sin 2 \theta_{s 1}=-\frac{\left(\sigma_{x}-\sigma_{y}\right)}{2 R} \text { and } \theta_{\mathrm{s} 1}=\theta_{\mathrm{P} 1}-45^{0}
$$

Therefore, $2 \theta_{s}-2 \theta_{p}=-90^{\circ}$ or $\theta_{s}=\theta_{p}+/-45^{\circ}$
The planes for maximum shear stress occurs at $45^{\circ}$ to the principal $\tan 2 \theta_{S}=-\frac{\left(\sigma_{x}-\sigma_{y}\right)}{2 \tau_{x y}}=-\frac{1}{\tan 2 \theta_{P}}=-\cot 2 \theta_{P}$
planes. The plane of the maximum planes. The plane of the maximum positive shear stress $\tau_{\max }$ is defined by the angle $\theta_{S 1}$ for which the

$$
\frac{\sin 2 \theta_{S}}{\cos 2 \theta_{S}}=-\frac{\cos 2 \theta_{P}}{\sin 2 \theta_{P}}=\frac{-\sin \left(90^{\circ}-2 \theta_{P}\right)}{\cos \left(90^{\circ}-2 \theta_{P}\right)}=\frac{\sin \left(2 \theta_{P}-90^{\circ}\right)}{\cos \left(2 \theta_{P}-90^{\circ}\right)}
$$ following equations apply:

The corresponding maximum shear is given by the equation

$$
\tau_{M A X}=\sqrt{\left(\frac{\sigma_{x}-\sigma_{y}}{2}\right)^{2}+\left(\tau_{x y}\right)^{2}}=R
$$

Another expression for the maximum shear stress

$$
\tau_{M A X}=\frac{\left(\sigma_{1}-\sigma_{2}\right)}{2}
$$

The normal stresses associated with the maximum shear stress are equal to

$$
\sigma_{A V E R}=\frac{\left(\sigma_{x}+\sigma_{y}\right)}{2}
$$

## Equations of a Circle

General equation

Consider

Equation (1)

$$
\begin{gathered}
\left(\sigma_{x 1}-\sigma_{A V E R}\right)^{2}=\left[\frac{\sigma_{x}-\sigma_{y}}{2} \cos 2 \theta+\tau_{x y} \sin 2 \theta\right]^{2} \\
\tau_{x 1 y 1}=-\frac{\sigma_{x}-\sigma_{y}}{2} \sin 2 \theta+\tau_{x y} \cos 2 \theta
\end{gathered}
$$

$$
\left(\tau_{x 1 y 1}\right)^{2}=\left[-\frac{\sigma_{x}-\sigma_{y}}{2} \sin 2 \theta+\tau_{x y} \cos 2 \theta\right]^{2}
$$

Equation (1) + Equation (2)

$$
\left(\sigma_{x 1}-\sigma_{\text {AVER }}\right)^{2}+\left(\tau_{x \mid y 1}\right)^{2}=\left[\frac{\sigma_{x}-\sigma_{y}}{2} \cos 2 \theta+\tau_{x y} \sin 2 \theta\right]^{2}+\left[-\frac{\sigma_{x}-\sigma_{y}}{2} \sin 2 \theta+\tau_{x y} \cos 2 \theta\right]^{2}
$$

$$
\left[\frac{\sigma_{x}-\sigma_{y}}{2} \cos 2 \theta+\tau_{x y} \sin 2 \theta\right]^{2}=\left(\frac{\sigma_{x}-\sigma_{y}}{2}\right)^{2} \cos ^{2} 2 \theta+\left(\tau_{x y}\right)^{2} \sin ^{2} 2 \theta+2\left(\frac{\sigma_{x}-\sigma_{y}}{2}\right)\left(\tau_{x y}\right) \sin 2 \theta \cos 2 \theta
$$

$$
\left[-\frac{\sigma_{x}-\sigma_{y}}{2} \sin 2 \theta+\tau_{x y} \cos 2 \theta\right]^{2}=\left(-\frac{\sigma_{x}-\sigma_{y}}{2}\right)^{2} \sin ^{2} 2 \theta+\left(\tau_{x y}\right)^{2} \cos ^{2} 2 \theta-2\left(\frac{\sigma_{x}-\sigma_{y}}{2}\right)\left(\tau_{x y}\right) \sin 2 \theta \cos 2 \theta
$$

$$
\begin{gathered}
=\left(\frac{\sigma_{x}-\sigma_{y}}{2}\right)^{2}+\left(\tau_{x y}\right)^{2}=R^{2} \\
\left(\sigma_{x 1}-\sigma_{A V E R}\right)^{2}+\left(\tau_{x 1 y 1}\right)^{2}=R^{2}
\end{gathered}
$$

## Mohr Circle

The radius of the Mohr circle is the magnitude $\boldsymbol{R}$.
$\tau$

$$
R=\sqrt{\left(\frac{\sigma_{x}-\sigma_{y}}{2}\right)^{2}+\left(\tau_{x y}\right)^{2}}
$$

The center of the Mohr circle is the magnitude

$$
\sigma_{A V E R}=\frac{\left(\sigma_{x}+\sigma_{y}\right)}{2}
$$

State of Stresses

$$
\left(\sigma_{x 1}-\sigma_{A V E R}\right)^{2}+\left(\tau_{x 1 y 1}\right)^{2}=\left(\frac{\sigma_{x}-\sigma_{y}}{2}\right)^{2}+\left(\tau_{x y}\right)^{2}
$$


(a)

(b)

(c)

Alternative sign conversion for shear stresses:
(a)clockwise shear stress,
(b)counterclockwise shear stress, and
(c) axes for Mohr's circle.

Note that clockwise shear stresses are plotted upward and counterclockwise shear stresses are plotted downward.

## Forms of Mohr's Circle

a) We can plot the normal stress $\sigma_{x 1}$ positive to the right and the shear stress $\tau_{x 1 y 1}$ positive downwards, i.e. the angle $2 \theta$ will be positive when counterclockwise or
b) We can plot the normal stress $\sigma_{x 1}$ positive to the right and the shear stress $\tau_{x 1 y 1}$ positive upwards, i.e. the angle $2 \theta$ will be positive when clockwise.
Both forms are mathematically correct. We use (a)

Two forms of Mohr's circle:
$\tau_{\mathrm{x} 1 \mathrm{y} 1}$ is positive downward and the angle $2 \theta$ is positive counterclockwise, and
$\tau_{\mathbf{x 1 y 1}}$ is positive upward and the angle $2 \theta$ is positive clockwise. (Note: The first form is used here)

(a)

(b)

(a)

(b)

Construction of Mohr's circle for plane stress.

(c)

## Example

At a point on the surface of a pressurized cylinder, the material is subjected to biaxial stresses $\sigma_{x}=90 M P a$ and $\sigma_{y}=$ 20MPa as shown in the element below.
Using the Mohr circle, determine the stresses acting on an element inclined at an angle $\boldsymbol{\theta}=\mathbf{3 0 ^ { \circ }}$ (Sketch a properly oriented element). ( $\sigma_{x}=90 \mathrm{MPa}, \sigma_{y}=20 \mathrm{MPa}$ and $\left.\tau_{x y}=0 \mathrm{MPa}\right)$


Because the shear stress is zero, these are the principal stresses.
Construction of the Mohr's circle The center of the circle is

$$
\sigma_{\text {Average }}=\frac{\left(\sigma_{x}+\sigma_{y}\right)}{2}=\frac{(90+20)}{2}=55 \mathrm{MPa}
$$



Stresses on an element inclined at $\boldsymbol{\theta}=30^{\circ}$
By inspection of the circle, the coordinates of point $\boldsymbol{D}$ are

$$
\begin{aligned}
& \sigma_{x 1}=\sigma_{\text {Average }}+R \cdot \operatorname{Cos} 60^{\circ}=55+35 \cdot \operatorname{Cos} 60=72.5 \mathrm{MPa} \\
& \tau_{x 1 y 1}=-R \cdot \operatorname{Sin} 60^{\circ}=-35\left(\operatorname{Sin} 60^{\circ}\right)=-30.3 \mathrm{MPa} \\
& \sigma_{y 1}=\sigma_{\text {Average }}-R \cdot \operatorname{Cos} 60^{\circ}=55-35 \cdot \operatorname{Cos} 60=37.5 \mathrm{MPa}
\end{aligned}
$$




Example
An element in plane stress at the surface of a lat machine is subjected to stresses $\sigma_{x}=15000$ psi, 5000psi and $\tau_{x y}=4000 p s i$.
Using the Mohr's circle determine the following:
a) The stresses acting on an element inclined at an angle $\theta=40^{\circ}$
b) The principal stresses and
c) The maximum shear stresses.


Construction of Mohr's circle:
Center of the circle (Point C):

$$
\sigma_{\text {Average }}=\frac{\left(\sigma_{x}+\sigma_{y}\right)}{2}=\frac{(15000+5000)}{2}=10000 \mathrm{psi}
$$

Radius of the circle:

$$
R=\sqrt{\left(\frac{\sigma_{x}-\sigma_{y}}{2}\right)^{2}+\left(\tau_{x y}\right)^{2}}=\sqrt{\left(\frac{15000-5000}{2}\right)^{2}+(4000)^{2}}=6403 p s i
$$

Point A, representing the stresses on the $x$ face of the element $\left(\theta=0^{\circ}\right)$ has the coordinates $\sigma_{\mathrm{x} 1}=15000 \mathrm{psi}$ and $\tau_{\mathrm{x} 1 \mathrm{y} 1}=4000 \mathrm{psi}$

Point B, representing the stresses on the $y$ face of the element $\left(\theta=90^{\circ}\right)$ has the coordinates $\sigma_{y 1}=5000 \mathrm{psi}$ and $\tau_{\mathrm{y} 1 \times 1}=-4000 \mathrm{psi}$
The circle is now drawn through points A and B with center C and radius R
$S_{2}\left(\theta_{s_{2}}=64.3^{\circ}\right)$


Stresses on an element inclined at $\theta$ $=40^{\circ}$

These are given by the coordinates of point $D$ which is at an angle $2 \boldsymbol{\theta}=80^{\circ}$ from point A

By inspection the angle $\mathrm{ACP}_{1}$ for the principal stresses (point $\mathrm{P}_{1}$ ) is :

$$
\tan \left(\overline{A C P_{1}}\right)=\frac{4000}{5000} \Rightarrow \overline{A C P_{1}}=38.66^{\circ}
$$

Then, the angle $\mathrm{P}_{1} \mathrm{CD}$ is $80^{\circ}-$ $38.66^{\circ}=41.34^{\circ}$

Knowing this angle, we can calculate the coordinates of point D (by inspection)

$$
\begin{aligned}
& \sigma_{x 1}=\sigma_{\text {Average }}+R \cdot \operatorname{Cos} 41.34^{\circ}=10000+6403 \cdot \operatorname{Cos} 41.34^{\circ}=14810 p s i \\
& \tau_{x 1 y 1}=-R \cdot \operatorname{Sin} 41.34^{o}=-6403\left(\operatorname{Sin} 41.34^{\circ}\right)=-4230 p s i \\
& \sigma_{y 1}=\sigma_{\text {Average }}-R \cdot \operatorname{Cos} 41.34^{\circ}=10000-6403 \cdot \operatorname{Cos} 41.34^{\circ}=5190 p s i
\end{aligned}
$$

And of course, the sum of the normal stresses is 14810psi $+5190 p s i=15000 p s i+5000 p s i$

## Principal Stresses

The principal stresses are represented by points $\mathrm{P}_{1}$ and $\mathrm{P}_{2}$ on Mohr's circle.
The angle it was found to be $2 \theta=38.66^{\circ}$ or $\theta=19.3^{\circ}$

$\tau_{x_{1} y_{1}}$

## Maximum Shear Stresses

These are represented by point $\mathrm{S}_{1}$ and $\mathrm{S}_{2}$ in Mohr's circle. Algebraically the maximum shear stress is given by the radius of the circle.
The angle $\mathrm{ACS}_{1}$ from point A to point $\mathrm{S}_{1}$ is $2 \theta_{\mathrm{S} 1}=51.34^{\circ}$. This angle is negative because is measured clockwise on the circle. Then the corresponding $\theta_{\mathrm{S} 1}$ value is $25.7^{\circ}$.

3-D stress state $\left[\begin{array}{ccc}15000 & 4000 & 0 \\ 4000 & 5000 & 0 \\ 0 & 0 & 0\end{array}\right]$ Psi Transform to $\left[\begin{array}{ccc}16403 & 0 & 0 \\ 0 & 3597 & 0 \\ 0 & 0 & 0\end{array}\right]$ Psi

In matrix notation the transformation is known as the Eigen-values.
The principal stresses are the "new-axes" coordinate system. The angles between the "old-axes" and the "new-axes" are known as the Eigen-vectors.

| principal stress | Cosine of angle <br> between $X$ and the <br> principal stress | Cosine of angle <br> between $Y$ and the <br> principal stress | Cosine of angle <br> between $Z$ and the <br> principal stress |
| :---: | :---: | :---: | :---: |
| $\mathbf{1 6 4 0 3 . 1 2 4 2}$ | 0.94362832 | 0.331006939 | 0 |
| 3596.876 | -0.33101 | 0.943628 | 0 |
| $\mathbf{0}$ | 0 | 0 | 1 |

## Example

At a point on the surface of a generator shaft the stresses are $\sigma_{\mathrm{x}}=-50 \mathrm{MPa}, \sigma_{\mathrm{y}}=10 \mathrm{MPa}$ and $\tau_{\mathrm{xy}}=-40 \mathrm{MPa}$ as shown in the figure. Using Mohr's circle determine the following:
(a)Stresses acting on an element inclined at an angle $\theta=45^{\circ}$,
(b)The principal stresses and
(c)The maximum shear stresses


Construction of Mohr's circle Center of the circle (Point C):

$$
\begin{gathered}
\sigma_{\text {Average }}=\frac{\left(\sigma_{x}+\sigma_{y}\right)}{2}=\frac{(-50)+(10)}{2}=-20 \mathrm{MPa} \\
R=\sqrt{\left(\frac{\sigma_{x}-\sigma_{y}}{2}\right)^{2}+\left(\tau_{x y}\right)^{2}}=\sqrt{\left(\frac{(-50)-(10)}{2}\right)^{2}+(-40)^{2}}=50 \mathrm{MPa}
\end{gathered}
$$

Point A, representing the stresses on the $x$ face of the element $\left(\theta=0^{\circ}\right)$ has the coordinates $\sigma_{x 1}=-50 \mathrm{MPa}$ and $\tau_{x 1 y 1}=-40 \mathrm{MPa}$
Point B , representing the stresses on the $y$ face of the element $\left(\theta=90^{\circ}\right)$ has the coordinates $\sigma_{\mathrm{y} 1}=10 \mathrm{MPa}$ and $\tau_{\mathrm{y} 1 \times 1}=40 \mathrm{MPa}$
The circle is now drawn through points A and B with center C and radius R .


Stresses on an element inclined at $\theta=45^{\circ}$ These stresses are given by the coordinates of point $D$ ( $2 \theta=90^{\circ}$ of point A). To calculate its magnitude we need to determine the angles ACP 2 and P2CD.

$$
\tan \mathrm{ACP}_{2}=40 / 30=4 / 3
$$

$$
\mathrm{ACP}_{2}=53.13^{\circ}
$$

$$
\mathrm{P}_{2} \mathrm{CD}=90^{\circ}-53.13^{\circ}=36.87^{\circ}
$$

Then, the coordinates of point D are

$$
\begin{aligned}
& \sigma_{x 1}=\sigma_{\text {Average }}+R \cdot \operatorname{Cos} 36.87^{\circ}=(-20)+50 \cdot \operatorname{Cos} 36.87^{\circ}=-60 M P a \\
& \tau_{x 1 y 1}=R \cdot \operatorname{Sin} 36.87^{\circ}=50\left(\operatorname{Sin} 36.87^{\circ}\right)=30 M P a \\
& \sigma_{y 1}=\sigma_{\text {Average }}-R \cdot \operatorname{Cos} 36.87^{\circ}=(-20)+50 \cdot \operatorname{Cos} 36.87^{\circ}=20 \mathrm{MPa}
\end{aligned}
$$

And of course, the sum of the normal stresses is $-50 \mathrm{MPa}+10 \mathrm{MPa}=-60 \mathrm{MPa}+20 \mathrm{MPa}$

## Principal Stresses

They are represented by points $\mathrm{P}_{1}$ and $\mathrm{P}_{2}$ on Mohr's circle.

$$
\begin{aligned}
& \sigma_{1}=\sigma_{\text {Average }}+R=-20+50=30 M P a \\
& \sigma_{2}=\sigma_{\text {Average }}-R=-20-50=-70 M P a
\end{aligned}
$$

The angle $\mathrm{ACP}_{1}$ is $2 \theta_{\mathrm{P} 1}=180^{\circ}+53.13^{\circ}$ $=233.13^{\circ}$ or $\theta_{\mathrm{P} 1}=116.6^{\circ}$
The angle $\mathrm{ACP}_{2}$ is $2 \theta_{\mathrm{P} 2}=53.13^{\circ}$ or $\theta_{\mathrm{P} 2}$ $=26.6^{\circ}$

## Maximum Shear Stresses

These are represented by point $S_{1}$ and $S_{2}$ in Mohr's circle.
The angle $\mathrm{ACS}_{1}$ is $2 \theta_{\mathrm{S} 1}=90^{\circ}+53.13^{\circ}=$ $143.13^{\circ}$ or $\theta=71.6^{\circ}$.
The magnitude of the maximum shear stress is 50 MPa and the normal stresses corresponding to point S 1 is the average stress -20MPa.

3-D stress state $\left[\begin{array}{ccc}-50 & -40 & 0 \\ -40 & 10 & 0 \\ 0 & 0 & 0\end{array}\right] M P a$ Transform to $\left[\begin{array}{ccc}30 & 0 & 0 \\ 0 & -70 & 0 \\ 0 & 0 & 0\end{array}\right] M P a$

In matrix notation the transformation is known as the Eigen-values.
The principal stresses are the "new-axes" coordinate system. The angles between the "old-axes" and the "new-axes" are known as the Eigen-vectors.

| principal stress | Cosine of angle <br> between $X$ and the <br> principal stress | Cosine of angle <br> between $Y$ and the <br> principal stress | Cosine of angle <br> between $Z$ and the <br> principal stress |
| :---: | :---: | :---: | :---: |
| 30 | -0.44721359 | 0.894427193 | 0 |
| -70 | 0.894427 | 0.447214 | 0 |
| 0 | 0 | 0 | 1 |

## Hooke's Law for Plane Stress

The stress transformations equations were derived solely from equilibrium conditions and they are material independent.
Here the material properties will be considered (strain) taking into account the following:
a)The material is uniform throughout the body (homogeneous)
b)The material has the same properties in all directions (isotropic)
c) The material follows Hooke's law (linearly elastic material)

Hooke's law: Linear relationship between stress and strain For uniaxial stress:
( $\boldsymbol{E}=$ modulus of elasticity or Young's modulus)

$$
\sigma=E \varepsilon
$$

Poisson's ratio: $\quad v=-\frac{\text { lateral strain }}{a \text { xial strain }}=-\frac{\varepsilon_{\text {transverse }}}{\varepsilon_{\text {longitudinal }}}$
For pure shear : $(\boldsymbol{G}=$ Shear modulus of elasticity $)$

$$
\tau=G \gamma
$$



Consider the normal strains $\varepsilon_{x}, \varepsilon_{y}, \varepsilon_{z}$ in plane stress.
All are shown with positive elongation.
The strains can be expressed in terms of the stresses by superimposing the effect of the individual stresses.
For instance the strain $\varepsilon_{x}$ in the $x$ direction:
a)Due to the stress $\sigma_{x}$ is equal to $\sigma_{x} / E$.
b)Due to the stress $\sigma_{y}$ is equal to $-v \sigma_{y} / E$. The resulting strain is:


$$
\begin{aligned}
& \varepsilon_{x}=\frac{\sigma_{x}}{E}-v \frac{\sigma_{y}}{E} \\
& \varepsilon_{y}=-v \frac{\sigma_{x}}{E}+\frac{\sigma_{y}}{E} \\
& \varepsilon_{z}=-v \frac{\sigma_{x}}{E}-v \frac{\sigma_{y}}{E}
\end{aligned}
$$

The shear stress causes a distortion of the element such that each $\boldsymbol{z}$ face becomes a rhombus.

$$
\gamma_{X Y}=\frac{\tau_{X Y}}{G}
$$

The normal stresses $\sigma_{x}$ and $\sigma_{y}$ have no effect on the shear strain $\gamma_{x y}$

$\varepsilon_{x}=\frac{\sigma_{x}}{E}-v \frac{\sigma_{y}}{E} \quad \varepsilon_{y}=-v \frac{\sigma_{x}}{E}+\frac{\sigma_{y}}{E} \quad \varepsilon_{z}=-\nu \frac{\sigma_{x}}{E}-v \frac{\sigma_{y}}{E}$
or rearranging the equations:

$$
\sigma_{x}=\frac{E \varepsilon_{x}}{\left(1-v^{2}\right)}+\frac{v E \varepsilon_{y}}{\left(1-v^{2}\right)} \quad \sigma_{Y}=\frac{v E \varepsilon_{x}}{\left(1-v^{2}\right)}+\frac{E \varepsilon_{y}}{\left(1-v^{2}\right)}+\quad \tau_{\mathrm{XY}}=G \gamma_{X Y}
$$

These equations are known collectively as Hooke's Law for plane stress
These equations contain three material constants ( $E, G$ and $v$ ) but only two are independent because of the relationship:

$$
G=\frac{E}{2(1+v)}
$$

## Special cases of Hooke's law ( $\sigma_{z}=0$ )

## Uniaxial stress :

$\sigma_{y}=0 \quad \tau_{x y}=0$
$\varepsilon_{x}=\frac{\sigma_{x}}{E} \quad \varepsilon_{y}=\varepsilon_{z}=-v \frac{\sigma_{x}}{E}$
Pure Shear :

$$
\begin{aligned}
& \sigma_{x}=\sigma_{y}=0 \\
& \varepsilon_{x}=\varepsilon_{y}=\varepsilon_{z}=0 \quad \gamma_{x y}=\frac{\tau_{x y}}{G}
\end{aligned}
$$

Biaxial stress :

$$
\tau_{x y}=0
$$

$$
\varepsilon_{x}=\frac{\sigma_{x}}{E}-v \frac{\sigma_{y}}{E} \quad \varepsilon_{y}=-v \frac{\sigma_{x}}{E}+\frac{\sigma_{y}}{E} \quad \varepsilon_{z}=-v \frac{\sigma_{x}}{E}-v \frac{\sigma_{y}}{E}
$$

## Volume Change

When a solid object undergoes strains, both its dimensions and its volume will change.

Consider an object of dimensions $a, b, c$. The original volume is $V_{o}=a b c$ and its final volume is
$V_{1}=\left(a+a \varepsilon_{\chi}\right)\left(b+b \varepsilon_{y}\right)\left(c+c \varepsilon_{z}\right)$
$V_{1}=a b c\left(1+\varepsilon_{\chi}\right)\left(1+\varepsilon_{y}\right)\left(1+\varepsilon_{\nu}\right)$


Upon expanding the terms:

$$
V_{1}=V_{o}\left(1+\varepsilon_{x}+\varepsilon_{y}+\varepsilon_{z}+\varepsilon_{x} \varepsilon_{y}+\varepsilon_{x} \varepsilon_{z}+\varepsilon_{y} \varepsilon_{z}+\varepsilon_{x} \varepsilon_{y} \varepsilon_{z}\right)
$$

For small strains:

$$
V_{1}=V_{o}\left(1+\varepsilon_{x}+\varepsilon_{y}+\varepsilon_{z}\right)
$$

The volume change is

$$
\Delta V=V_{1}-V_{o}=V_{o}\left(\varepsilon_{x}+\varepsilon_{y}+\varepsilon_{z}\right)
$$

The unit volume change $e$, also known as dilatation is defines as:

$$
e=\Delta V / V_{o}=\varepsilon_{x}+\varepsilon_{y}+\varepsilon_{z}
$$

Positive strains are considered for elongations and negative strains for shortening, i.e. positive values of $e$ for an increase in volume.

$$
\varepsilon_{x}=\frac{\sigma_{x}}{E}-v \frac{\sigma_{y}}{E} \quad \varepsilon_{y}=-v \frac{\sigma_{x}}{E}+\frac{\sigma_{y}}{E} \quad \varepsilon_{z}=-v \frac{\sigma_{x}}{E}-v \frac{\sigma_{y}}{E}
$$

$$
e=\frac{\Delta V}{V}=\left(\sigma_{x}+\sigma_{y}\right) \frac{(1-2 v)}{E}
$$

For uniaxial stress $\sigma_{y}=0$

$$
e=\sigma_{x} \frac{(1-2 v)}{E}
$$

We can notice that the maximum possible value of Poisson's ratio is 0.5 , because a larger value means that the volume decreases when the material is in tension (contrary to physical behavior).

## STRAIN ENERGY DENSITY IN PLANE STRESS

The strain energy density $\boldsymbol{u}$ is the strain energy stored in a unit volume of the material.Because the normal and shear strains occur independently, we can add the strain energy of these two elements to obtain the total energy.

Work done $=$ Force $\times$ distance
Work done in the x - direction $=\frac{\left(\sigma_{\mathrm{x}}\right)(b c)}{2}\left(a \varepsilon_{x}\right)$
Work done in the y - direction $=\frac{\left(\sigma_{\mathrm{y}}\right)(a c)}{2}\left(b \varepsilon_{y}\right)$
The sum of the energies due to normal stresses:

$$
U=\frac{a b c}{2}\left(\sigma_{x} \varepsilon_{x}+\sigma_{y} \varepsilon_{y}\right)
$$

Then the strain energy density (strain per unit volume) $\quad u_{1}=\frac{1}{2}\left(\sigma_{x} \varepsilon_{x}+\sigma_{y} \varepsilon_{y}\right)$
Similarly, the strain energy density associated with the shear strain:

$$
u_{2}=\frac{1}{2} \tau_{x y} \gamma_{x y}
$$

By combining the strain energy densities for the normal and shear strains:

$$
u=\frac{1}{2}\left(\sigma_{x} \varepsilon_{x}+\sigma_{y} \varepsilon_{y}+\tau_{x y} \gamma_{x y}\right)
$$

The strain energy density in terms of stress alone:

$$
u=\frac{\sigma_{X}^{2}}{2 E}+\frac{\sigma_{Y}^{2}}{2 E}-v \frac{\sigma_{X} \sigma_{Y}}{E}+\frac{\tau_{X Y}^{2}}{2 G}
$$

The strain energy density in terms of strain alone:

$$
u=\frac{E}{2\left(1-v^{2}\right)}\left(\varepsilon_{X}^{2}+\varepsilon_{Y}^{2}-2 v \varepsilon_{X} \varepsilon_{Y}\right)+\frac{G}{2} \gamma_{X Y}^{2}
$$

For the special case of uniaxial stress:

$$
\begin{aligned}
& \sigma_{y}=0 \quad \tau_{x y}=0 \quad \varepsilon_{y}=-v \varepsilon_{x} \quad \gamma_{x y}=0 \\
& u=\frac{\sigma_{x}^{2}}{2 E} \quad \text { or } \quad u=\frac{E \varepsilon_{x}^{2}}{2}
\end{aligned}
$$

## TRIAXIAL STRESS

An element of the material subjected to normal stresses $\sigma_{x} \sigma_{y}$ and $\sigma_{z}$ acting in three mutually perpendicular directions is said to be in a state of triaxial stress. Since there is no shear in $x, y$ or $z$ faces then the stresses $\sigma_{x} \sigma_{y}$ and $\sigma_{z}$ are the principal stresses in the material.

For the special case of pure shear:

$$
\begin{array}{lll}
\sigma_{x}=0 & \sigma_{y}=0 & \varepsilon_{y}=\varepsilon_{x}=0 \\
u=\frac{\tau_{x y}^{2}}{2 G} & \text { or } & u=\frac{G \gamma_{x y}^{2}}{2}
\end{array}
$$



If an inclined plane parallel to the $\mathbf{z}$-axis is cut through the element, the only stress of the inclined face are the normal stress $\sigma$ and the shear stress $\tau$, both of which act parallel to the $x y$ plane.


Because these stresses are independent of the $\sigma_{z}$, we can use the transformation equations of plane stress, as well as the Mohr's circle for plane stress, when determining the stresses $\sigma$ and $\tau$ in triaxial stress.

The same general conclusion hold for normal and shear stresses acting on inclined planes cut through the element parallel to the $x$ and $y$ axes.
Maximum Shear Stress
For a material in triaxial stress, the maximum shear stresses occur on elements oriented at angles of $45^{\circ}$ to the $\boldsymbol{x}, \boldsymbol{y}$ and $\boldsymbol{z}$ axes.
for the inclined plane $/ / \mathrm{z}$-axis
for the inclined plane // x -axis
for the inclined plane // y-axis

$$
\begin{aligned}
& \left(\tau_{M A X}\right)_{Z}= \pm \frac{\sigma_{X}-\sigma_{Y}}{2} \\
& \left(\tau_{M A X}\right)_{X}= \pm \frac{\sigma_{Y}-\sigma_{Z}}{2} \\
& \left(\tau_{M A X}\right)_{Y}= \pm \frac{\sigma_{X}-\sigma_{Z}}{2}
\end{aligned}
$$

The absolute maximum of the shear stress is the numerically largest of the above.

The stresses acting on elements oriented at various angles to the $\boldsymbol{x}, \boldsymbol{y}$ and $\boldsymbol{z}$ axes can be visualized with the aid of the Mohr's circle.

In this case $\sigma_{x}>\sigma_{y}>\sigma_{z}$


If Hooke's law is obeyed, it is possible to obtain the relationship between normal stresses and normal strains

$$
\begin{aligned}
& \varepsilon_{x}=\frac{\sigma_{x}}{E}-v \frac{\sigma_{y}}{E}-v \frac{\sigma_{z}}{E} \\
& \varepsilon_{y}=-v \frac{\sigma_{x}}{E}+\frac{\sigma_{y}}{E}-v \frac{\sigma_{z}}{E}
\end{aligned}
$$ using the same procedure as for plane stress.

Or

$$
\sigma_{x}=\frac{E}{(1+v)(1-2 v)}\left[(1-v) \varepsilon_{x}+v \varepsilon_{y}+v \varepsilon_{z}\right]
$$

$$
\varepsilon_{z}=-v \frac{\sigma_{x}}{E}-v \frac{\sigma_{y}}{E}+\frac{\sigma_{z}}{E}
$$

$$
\sigma_{y}=\frac{E}{(1+v)(1-2 v)}\left[v \varepsilon_{x}+(1-v) \varepsilon_{y}+v \varepsilon_{z}\right]
$$

They are known as the Hooke's law for triaxial

$$
\sigma_{z}=\frac{E}{(1+v)(1-2 v)}\left[v \varepsilon_{x}+v \varepsilon_{y}+(1-v) \varepsilon_{z}\right]
$$ stress.

## Unit Volume Change

The unit volume change (or dilatation) for an element in triaxial stress is obtained in the same manner as for plane stress.

$$
e=\frac{\Delta V}{V_{O}}=\varepsilon_{X}+\varepsilon_{Y}+\varepsilon_{Z}
$$

If Hooke's laws apply, then

$$
e=\frac{\Delta V}{V_{O}}=\frac{(1-2 v)}{E}\left(\sigma_{X}+\sigma_{Y}+\sigma_{Z}\right)
$$

## Strain Energy Density

The strain energy density for an element in triaxial stress is obtained by the same method used for plane stress.

In terms of the strains:

$$
u=\frac{\sigma_{x} \varepsilon_{x}}{2}+\frac{\sigma_{y} \varepsilon_{y}}{2}+\frac{\sigma_{z} \varepsilon_{z}}{2}
$$

In terms of stresses :

$$
u=\frac{E}{2(1+v)(1-2 v)}\left[(1-v)\left(\varepsilon_{x}^{2}+\varepsilon_{y}^{2}+\varepsilon_{z}^{2}\right)+2 v\left(\varepsilon_{x} \varepsilon_{y}+\varepsilon_{x} \varepsilon_{z}+\varepsilon_{y} \varepsilon_{z}\right)\right]
$$

## Spherical Stress

A special case of triaxial stress, called spherical stress, occurs

$$
\sigma_{x}=\sigma_{y}=\sigma_{z}=\sigma_{0}
$$ whenever all three normal stresses are equal:

The Mohr's circle is reduced to a single point. Any plane cut through the element will be free of shear stress and will be subjected to the same normal stress so and it is a principal plane.
The normal strains in spherical stress are also the same in all directions, provided the material is isotropic and if Hooke's law applies:

$$
\begin{aligned}
& \varepsilon_{O}=\frac{\sigma_{O}}{E}(1-2 v) \text { The volume change } \\
& \mathrm{e}=3 \varepsilon_{0}=3 \sigma_{O} \frac{(1-2 v)}{E}
\end{aligned}
$$



Element in spherical stress.
$\boldsymbol{K}=$ bulk or volume modulus of elasticity

$$
K=\frac{E}{3(1-2 v)} \quad K=\frac{\sigma_{0}}{e}
$$

If $v=1 / 3$ then $K=E$
If $v=\mathbf{0}$ then $\boldsymbol{K}=\boldsymbol{E} / \mathbf{3}$
If $v=\mathbf{1} / \mathbf{2}$ then $K=$ infinite (rigid material having no change in volume)
These formulas also apply to and element in uniform compression, for example rock deep within the earth or material submerged in water (hydrostatic stress).

## PLAIN STRAIN

Strains are measured by strain gages.
A material is said to be in a state of plain strain if the only deformations are those in the $\boldsymbol{x y}$ plane, i.e. it has only three strain components $\varepsilon_{x}, \varepsilon_{y}$ and $\gamma_{x y}$.
Plain stress is analogous to plane stress, but under ordinary conditions they do not occur simultaneously
Exception when $\sigma_{x}=-\sigma_{y}$ and when $v=0$
Strain components $\varepsilon_{x}, \varepsilon_{y}$, and $\gamma_{\mathrm{xy}}$ in the $x y$ plane (plane strain).


|  | Plane stress | Plane strain |
| :---: | :---: | :---: |
|  |  |  |
| Stresses | $\begin{aligned} & \sigma_{z}=0 \quad \tau_{x z}=0 \quad \tau_{y z}=0 \\ & \sigma_{x}, \sigma_{y}, \text { and } \tau_{x y} \text { may have } \\ & \text { nonzero values } \end{aligned}$ | $\tau_{x z}=0 \quad \tau_{y z}=0$ <br> $\sigma_{x}, \sigma_{y}, \sigma_{z}$, and $\tau_{x y}$ may have nonzero values |
| Strains | $\gamma_{x z}=0 \quad \gamma_{y z}=0$ <br> $\varepsilon_{x}, \varepsilon_{y}, \varepsilon_{z}$, and $\gamma_{x}$ may have nonzero values | $\begin{aligned} & \varepsilon_{z}=0 \quad \gamma_{x z}=0 \quad \gamma_{y z}=0 \\ & \varepsilon_{x}, \varepsilon_{y}, \text { and } \quad \gamma_{x y} \text { may have } \\ & \text { nonzero values } \end{aligned}$ |

Comparison of plane stress and plane strain.

## APPLICATION OF THE TRANSFORMATION EQUATIONS

The transformation equations for plane stress are valid even when $\sigma_{z}$ is not zero, because $\sigma_{z}$ does not enter the equations of equilibrium. Therefore, the transformations equations for plane stress can also be used for stresses in plane strain.
Similarly, the strain transformation equations that will be derived for the case of plain strain in the $x y$ plane are valid even when $\varepsilon_{z}$ is not zero, because the strain $\varepsilon_{z}$ does not affect the geometric relationship used for the derivation. Therefore, the transformations equations for plane strain can also be used for strains in plane stress.

## Transformation Equations for Plain Strain

We will assume that the strain $\varepsilon_{x}, \varepsilon_{y}$ and $\gamma_{x y}$ associated with the $x y$ plane are known.
We need to determine the normal and shear strains $\left(\varepsilon_{x 1}\right.$ and $\left.\gamma_{x y y}\right)$ associated with the $x_{1} y_{1}$ axis. $\varepsilon_{y 1}$ can be obtained from the equation of $\varepsilon_{x 1}$ by substituting $\theta+90$ for $\theta$.


For an element of size $\boldsymbol{d x}, \boldsymbol{d y}$
In the $x$ direction:
the strain $\varepsilon_{x}$ produces an elongation $\varepsilon_{x} d x$.
The diagonal increases in length by $\varepsilon_{x} d x \cos \theta$.


In the $y$ direction:
the strain $\varepsilon_{y}$ produces an elongation $\varepsilon_{y} d y$.
The diagonal increases in length by $\varepsilon_{y} d y \sin \theta$.

The shear strain $\gamma_{x y}$ in the plane $x y$ produces a distortion of the element such that the angle at the lower left corner decreases by an amount equal to the shear strain. Consequently, the upper face moves to the right by an amount $\gamma_{x y} d y$. This deformation results in an increase in the length of the diagonal equal to: $\gamma_{x y} d y \cos \theta$


The total increase $\boldsymbol{\Delta} \boldsymbol{\delta}$ of the diagonal is the sum of the preceding three expressions, thus:

$$
\Delta d=\varepsilon_{x} d x(\operatorname{Cos} \theta)+\varepsilon_{y} d y(\operatorname{Sin} \theta)+\gamma_{x y} d y(\operatorname{Cos} \theta)
$$

But

$$
\begin{aligned}
& \varepsilon_{x 1}=\frac{\Delta d}{d s}=\varepsilon_{x}\left(\frac{d x}{d s}\right)(\operatorname{Cos} \theta)+\varepsilon_{y}\left(\frac{d y}{d s}\right)(\operatorname{Sin} \theta)+\gamma_{x y}\left(\frac{d y}{d s}\right)(\operatorname{Cos} \theta) \\
& \frac{d x}{d s}=\operatorname{Cos} \theta \quad \frac{d y}{d s}=\operatorname{Sin} \theta \\
& \quad \varepsilon_{x 1}=\varepsilon_{x} \operatorname{Cos}^{2} \theta+\varepsilon_{y} \operatorname{Sin}^{2} \theta+\gamma_{x y} \operatorname{Sin} \theta \operatorname{Cos} \theta
\end{aligned}
$$

Shear Strain $\gamma_{x 1 y 1}$ associated with $x_{1} y_{1}$ axes.
This strain is equal to the decrease in angle between lines in the material that were initially along the $\boldsymbol{x} \mathbf{1}$ and $\boldsymbol{y} \mathbf{1}$ axes.
$\boldsymbol{O a}$ and $\boldsymbol{O b}$ were the lines initially along the $\boldsymbol{x} \mathbf{1}$ and $y 1$ axis respectively. The deformation caused by the strains $\varepsilon_{x}, \varepsilon_{y}$ and $\gamma_{x y}$ caused the $\boldsymbol{O} \boldsymbol{a}$ and $\boldsymbol{O b}$ lines to rotate and angle $\alpha$ and $\beta$ from the $\boldsymbol{x}_{1}$ and $\boldsymbol{y}_{1}$ axis respectively. The shear strain $\gamma_{x 1 y 1}$ is the decrease in angle between the two lines that originally were at right angles, therefore, $\gamma_{x 1 y 1}=\alpha+\beta$.


The angle $\alpha$ can be found from the deformations produced by the strains $\varepsilon_{x}, \varepsilon_{y}$ and $\gamma_{x y}$. The strains $\varepsilon_{x}$ and $\gamma_{x y}$ produce a clockwise rotation, while the strain $\varepsilon_{y}$ produces a counterclockwise rotation.



Let us denote the angle of rotation produced by $\varepsilon_{x}, \varepsilon_{y}$ and $\gamma_{x y}$ as $\alpha_{1}, \alpha_{2}$ and $\alpha_{3}$ respectively.

$$
\begin{aligned}
& \alpha_{1}=\varepsilon_{x} \operatorname{Sin} \theta \frac{d x}{d s} \quad \frac{d x}{d s}=\operatorname{Cos} \theta \quad \frac{d y}{d s}=\operatorname{Sin} \theta \\
& \alpha_{2}=\varepsilon_{y} \operatorname{Cos} \theta \frac{d y}{d s} \\
& \alpha_{1}=\gamma_{x y} \operatorname{Sin} \theta \frac{d y}{d s} \\
& \alpha=-\alpha_{1}+\alpha_{2}-\alpha_{3}=-\left(\varepsilon_{x}-\varepsilon_{y}\right) \operatorname{Sin} \theta \operatorname{Cos} \theta-\gamma_{x y} \operatorname{Sin}^{2} \theta
\end{aligned}
$$

The rotation of line $O \boldsymbol{b}$ which initially was at $90^{\circ}$ to the line $O a$ can be found

$$
\beta=\left(\varepsilon_{x}-\varepsilon_{y}\right) \operatorname{Sin}(\theta+90) \operatorname{Cos}(\theta+90)+\gamma_{x y} \operatorname{Sin}^{2}(\theta+90)
$$ by substituting $\theta+90$ for $\theta$ in the expression for $\alpha$. Because $\beta$ is positive $\beta=-\left(\varepsilon_{x}-\varepsilon_{y}\right) \operatorname{Sin} \theta \operatorname{Cos} \theta+\gamma_{x y} \operatorname{Cos}^{2} \theta$ when clockwise. Thus:

$$
\gamma_{x 1 y 1}=\alpha+\beta=-\left(\varepsilon_{x}-\varepsilon_{y}\right) \operatorname{Sin} \theta \operatorname{Cos} \theta+\frac{\gamma_{x y}}{2}\left[\operatorname{Cos}^{2} \theta-\operatorname{Sin}^{2} \theta\right]
$$

## Transformation Equations for Plain Strain

Using the following trigonometric identities:
$\operatorname{Cos}^{2} \theta=1 / 2(1+\operatorname{Cos} 2 \theta)$ $\operatorname{Sin}^{2} \theta=1 / 2(1-\operatorname{Cos} 2 \theta)$ $\operatorname{Sin} \theta \cos \theta=1 / 2 \operatorname{Sin} 2 \theta$

$$
\begin{aligned}
& \varepsilon_{x 1}=\frac{\left(\varepsilon_{x}+\varepsilon_{y}\right)}{2}+\frac{\left(\varepsilon_{x}-\varepsilon_{y}\right)}{2} \cos 2 \theta+\frac{\gamma_{x y}}{2} \sin 2 \theta \\
& \frac{\gamma_{x 1 y 1}}{2}=-\frac{\left(\varepsilon_{x}-\varepsilon_{y}\right)}{2} \sin 2 \theta+\frac{\gamma_{x y}}{2} \cos 2 \theta
\end{aligned}
$$

$$
\text { In variant }=\varepsilon_{x}+\varepsilon_{y}=\varepsilon_{x 1}+\varepsilon_{y 1}
$$

$$
\varepsilon_{\text {Average }}=\frac{\left(\varepsilon_{x}+\varepsilon_{y}\right)}{2}
$$

## PRINCIPAL STRAINS

The angle for the principal strains is :
The value for the principal strains are

$$
\tan 2 \theta_{P}==\frac{\gamma_{x y} / 2}{\left(\varepsilon_{x}-\varepsilon_{y}\right) / 2}=\frac{\gamma_{x y}}{\varepsilon_{x}-\varepsilon_{y}}
$$

$$
\begin{aligned}
& \varepsilon_{1}=\frac{\left(\varepsilon_{x}+\varepsilon_{y}\right)}{2}+\sqrt{\left(\frac{\varepsilon_{x}-\varepsilon_{y}}{2}\right)^{2}+\left(\frac{\gamma_{x y}}{2}\right)^{2}} \\
& \varepsilon_{2}=\frac{\left(\varepsilon_{x}+\varepsilon_{y}\right)}{2}-\sqrt{\left(\frac{\varepsilon_{x}-\varepsilon_{y}}{2}\right)^{2}+\left(\frac{\gamma_{x y}}{2}\right)^{2}}
\end{aligned}
$$

## Maximum Shear

The maximum shear strains in the $x y$ plane are associated with axes at $45^{\circ}$ to the directions of

$$
\frac{\gamma_{M a x}}{2}=+\sqrt{\left(\frac{\varepsilon_{x}-\varepsilon_{y}}{2}\right)^{2}+\left(\frac{\gamma_{x y}}{2}\right)^{2}} \text { or } \gamma_{M a x}=\left(\varepsilon_{1}-\varepsilon_{2}\right)
$$ the principal strains

For isotropic materials, at a given point in an stressed body, the principal strains and principal stresses occur in the same directions.

## MOHR'S CIRCLE FOR PLANE STRAIN

It is constructed in the same manner as the Mohr's circle for plane stress with the following similarities:


$$
\begin{array}{c||c}
\sigma_{x} & \varepsilon_{x} \\
\sigma_{y} & \varepsilon_{y} \\
\tau_{x y} & \gamma_{x y} / 2 \\
\sigma_{x 1} & \varepsilon_{x 1} \\
\tau_{x \mid y 1} & \gamma_{x \mid y 1} / 2
\end{array}
$$

## Strain Measurements

An electrical-resistance strain gage is a device for measuring normal strains $(\varepsilon)$ on the surface of a stressed object.
The gages are small (less than $1 / 2$ inch) made of wires that are bonded to the surface of the object. Each gage that is stretched or shortened when the object is strained at the point, changes its electrical resistance. This change in resistance is converted into a measurement of strain.
From three measurements it is possible to calculate the strains in any direction. A group of three gages arranged in a particular pattern is called a strain rosette. Because the rosette is mounted in the surface of the body, where the material is in plane stress, we can use the transformation equations for plane strain to calculate the strains in various directions.

$45^{\circ}$ strain rosette, and element oriented at an angle $\boldsymbol{\theta}$ to the $\boldsymbol{x y}$ axes.

## General Equations




## Example

An element of material in plane strain undergoes the following strains: $\varepsilon_{x}=340 \times 10^{-6}$ $; \varepsilon_{y}=110 \times 10^{-6} \quad ; \gamma_{x y}=180 \times 10^{-6}$. Determine the following quantities:
(a) the strains of an element oriented at an angle $\theta=30^{\circ}$;
(b) the principal strains and
(c) the maximum shear strains.
(a) Element oriented at an angle

$$
\theta=30^{\circ}\left(2 \theta=60^{\circ}\right)
$$

$\varepsilon_{x 1}=\frac{\left(\varepsilon_{x}+\varepsilon_{y}\right)}{2}+\frac{\left(\varepsilon_{x}-\varepsilon_{y}\right)}{2} \cos 2 \theta+\frac{\gamma_{x y}}{2} \sin 2 \theta$
$\varepsilon_{x 1}=\left[\frac{340+110}{2}+\frac{340-110}{2} \operatorname{Cos} 60+\frac{180}{2} \operatorname{Sin} 60\right] \times 10^{-6}$

$$
\varepsilon_{x 1}=360 \times 10^{-6}
$$



$$
\begin{aligned}
& \frac{\gamma_{x 1 y 1}}{2}=-\frac{\left(\varepsilon_{x}-\varepsilon_{y}\right)}{2} \sin 2 \theta+\frac{\gamma_{x y}}{2} \cos 2 \theta \\
& \frac{\gamma_{x 1 y 1}}{2}=\left[-\frac{340-110}{2} \operatorname{Sin} 60+\frac{180}{2} \operatorname{Cos} 60\right] \times 10^{-6}=-55 \times 10^{-6}
\end{aligned}
$$

$$
\begin{aligned}
& \varepsilon_{x}+\varepsilon_{y}=\varepsilon_{x 1}+\varepsilon_{y 1} \\
& 340 \mu+110 \mu=360 \mu+\varepsilon_{y 1} \\
& \varepsilon_{y 1}=90 \mu
\end{aligned}
$$

$$
\varepsilon_{\text {Average }}=225 \mu
$$

## (b) Principal Strains and Angle of Rotation

$$
\begin{aligned}
& \varepsilon_{1,2}=\frac{\left(\varepsilon_{x}+\varepsilon_{y}\right)}{2} \pm \sqrt{\left(\frac{\varepsilon_{x}-\varepsilon_{y}}{2}\right)^{2}+\left(\frac{\gamma_{x y}}{2}\right)^{2}} \\
& \varepsilon_{1}=225 \mu+\sqrt{\left(\frac{340 \mu-110 \mu}{2}\right)^{2}+\left(\frac{180}{2}\right)^{2}}=370 \mu \\
& \varepsilon_{2}=225 \mu-\sqrt{\left(\frac{340 \mu-110 \mu}{2}\right)^{2}+\left(\frac{180}{2}\right)^{2}}=80 \mu
\end{aligned}
$$


$\tan 2 \theta_{P}=\frac{\gamma_{x y}}{\varepsilon_{x}-\varepsilon_{y}}=\frac{-55 \mu}{340 \mu-110 \mu}=0.7826$
$\theta_{P}=19^{0}$

(c) In-Plane Maximum Shear Strain

$$
\begin{aligned}
& \frac{\gamma_{M a x}}{2}=+\sqrt{\left(\frac{\varepsilon_{x}-\varepsilon_{y}}{2}\right)^{2}+\left(\frac{\gamma_{x y}}{2}\right)^{2}}=\sqrt{\left(\frac{340 \mu-110 \mu}{2}\right)^{2}+\left(\frac{180 \mu}{2}\right)^{2}}=145 \mu \\
& \gamma_{\operatorname{Max}}=\left(\varepsilon_{1}-\varepsilon_{2}\right)=370 \mu-80 \mu=290 \mu
\end{aligned}
$$

$$
\varepsilon_{\text {Average }}=225 \mu
$$


(d) Out-of-Plane Maximum Shear Strain

$$
\gamma_{M a x}=\left(\varepsilon_{1}-\varepsilon_{3}\right)=370 \mu-0 \mu=370 \mu
$$

## Transformation Equations

$$
\begin{aligned}
& \varepsilon_{X 1}=\varepsilon_{X} \cos ^{2} \theta+\varepsilon_{Y} \sin ^{2} \theta+\left(\frac{\gamma_{X Y}}{2}\right) 2 \sin \theta \cos \theta \\
& \varepsilon_{Y 1}=\varepsilon_{X} \sin ^{2} \theta+\varepsilon_{Y} \cos ^{2} \theta-\left(\frac{\gamma_{X Y}}{2}\right) 2 \sin \theta \cos \theta
\end{aligned}
$$

$$
\frac{\gamma_{X|Y|}}{2}=-\varepsilon_{X} \sin \theta \cos \theta+\varepsilon_{Y} \sin \theta \cos \theta+\frac{\gamma_{X Y}}{2}\left(\cos ^{2} \theta-\sin ^{2} \theta\right)
$$

$$
\left[\begin{array}{c}
\varepsilon_{X 1} \\
\varepsilon_{Y 1} \\
\frac{\gamma_{X 1 Y}}{2}
\end{array}\right]=[T] \times\left[\begin{array}{c}
\varepsilon_{X} \\
\varepsilon_{Y} \\
\frac{\gamma_{X Y}}{2}
\end{array}\right] \quad[T]=\left[\begin{array}{ccc}
\cos ^{2} \theta & \sin ^{2} \theta & 2 \sin \theta \cos \theta \\
\sin ^{2} \theta & \cos ^{2} \theta & -2 \sin \theta \cos \theta \\
-\sin \theta \cos \theta & \sin \theta \cos \theta & \left(\cos ^{2} \theta-\sin ^{2} \theta\right)
\end{array}\right]
$$

$$
\left[\begin{array}{c}
\varepsilon_{X} \\
\varepsilon_{Y} \\
\frac{\gamma_{X I}}{2}
\end{array}\right]=[T]^{-1} \times\left[\begin{array}{c}
\varepsilon_{X 1} \\
\frac{\varepsilon_{X 1}}{2}
\end{array}\right]\left[\begin{array}{c}
\varepsilon_{X 1} \\
\frac{\gamma_{X 1}}{2}
\end{array}\right]=\left[\begin{array}{ccc}
\cos ^{2} \theta & \sin ^{2} \theta & 2 \sin \theta \cos \theta \\
\varepsilon_{Y 1} \operatorname{rin}^{2} \theta & \cos ^{2} \theta & -2 \sin \theta \cos \theta \\
-\sin \theta \cos \theta & \sin \theta \cos \theta & \left(\cos ^{2} \theta-\sin ^{2} \theta\right)
\end{array}\right]\left[\begin{array}{c}
\varepsilon_{X} \\
\varepsilon_{Y} \\
\frac{\gamma_{X Y}}{2}
\end{array}\right]
$$

## Example For $\Theta=30$ degrees

$$
\begin{aligned}
& {\left[\begin{array}{c}
\varepsilon_{X 1} \\
\varepsilon_{Y 1} \\
\frac{\gamma_{X 1 Y 1}}{2}
\end{array}\right]=\left[\begin{array}{ccc}
\cos ^{2} 30 & \sin ^{2} 30 & 2 \sin 30 \cos 30 \\
\sin ^{2} 30 & \cos ^{2} 30 & -2 \sin 30 \cos 30 \\
-\sin 30 \cos 30 & \sin 30 \cos 30 & \left(\cos ^{2} 30-\sin ^{2} 30\right)
\end{array}\right]\left[\begin{array}{c}
\varepsilon_{X} \\
\varepsilon_{Y} \\
\frac{\gamma_{X Y}}{2}
\end{array}\right]} \\
& {\left[\begin{array}{c}
\varepsilon_{X 1} \\
\varepsilon_{Y 1} \\
\frac{\gamma_{X 1 Y 1}}{2}
\end{array}\right]=\left[\begin{array}{ccc}
0.75 & 0.25 & 0.876 \\
0.25 & 0.75 & -0.876 \\
-0.438 & 0.438 & 0.5
\end{array}\right]\left[\begin{array}{c}
340 \\
110 \\
90
\end{array}\right] \mu=\left[\begin{array}{c}
361.3 \\
88.6 \\
-55.8
\end{array}\right] \mu}
\end{aligned}
$$

F

$$
]
$$

$$
\checkmark
$$

$$
[\varepsilon]=\left[\begin{array}{ccc}
\varepsilon_{x x} & \frac{1}{2} \gamma_{y x} & \frac{1}{2} \gamma_{z x} \\
\frac{1}{2} \gamma_{x y} & \varepsilon_{y y} & \frac{1}{2} \gamma_{z y} \\
\frac{1}{2} \gamma_{x z} & \frac{1}{2} \gamma_{y z} & \varepsilon_{z z}
\end{array}\right]=\text { Strain_Tensor }
$$

## Example

$$
[\varepsilon]=\left[\begin{array}{ccc}
\varepsilon_{x x} & \frac{1}{2} \gamma_{y x} & \frac{1}{2} \gamma_{z x} \\
\frac{1}{2} \gamma_{x y} & \varepsilon_{y y} & \frac{1}{2} \gamma_{z y} \\
\frac{1}{2} \gamma_{x z} & \frac{1}{2} \gamma_{y z} & \varepsilon_{z z}
\end{array}\right]=\left[\begin{array}{ccc}
340 & 180 / 2 & 0 \\
180 / 2 & 110 & 0 \\
0 & 0 & 0
\end{array}\right] \mu
$$

Eigen_Values $\Rightarrow\left[\begin{array}{ccc}371 & 0 & 0 \\ 0 & 79 & 0 \\ 0 & 0 & 0\end{array}\right] \mu$


## Example

A $45^{\circ}$ strain rosette (rectangular rosette) consists of three electrical-resistance strain gages, arranged to measure strains in two perpendicular directions and also at a $45^{\circ}$ angle (as shown below). The rosette is bonded to the surface of the structure before it is loaded. Gages $\boldsymbol{A}, \boldsymbol{B}$ and $\boldsymbol{C}$ measure the normal strains $\boldsymbol{\varepsilon}_{\boldsymbol{a}}, \boldsymbol{\varepsilon}_{\boldsymbol{b}}$ and $\boldsymbol{\varepsilon}_{\boldsymbol{c}}$ in the directions of the lines $\boldsymbol{O a}, \boldsymbol{O b}$ and $\boldsymbol{O c}$, respectively.
Explain how to obtain the strains $\varepsilon_{x 1}, \varepsilon_{y 1}$ and $\gamma_{x 1 y 1}$, associated with an element oriented at an angle $\boldsymbol{\theta}$ to the $\boldsymbol{x y}$ axes.



$$
\varepsilon_{a}=\varepsilon_{x} \cos ^{2} \theta+\varepsilon_{y} \sin ^{2} \theta+\left(\frac{\gamma_{x y}}{2}\right) 2 \sin \theta \cos \theta
$$

Angles with respect to x -axis: (a) is zero ; (b) is 45 degrees CCW ; (c) 90 degrees CCW

$$
\begin{aligned}
& \varepsilon_{a}=\varepsilon_{x} \cos ^{2} 0+\varepsilon_{y} \sin ^{2} 0+\left(\frac{\gamma_{x y}}{2}\right) 2 \sin 0 \cos 0 \\
& \varepsilon_{b}=\varepsilon_{x} \cos ^{2} 45+\varepsilon_{y} \sin ^{2} 45+\left(\frac{\gamma_{x y}}{2}\right)(2 \sin 45 \cos 45) \\
& \varepsilon_{b}=\varepsilon_{x} \cos ^{2} 90+\varepsilon_{y} \sin ^{2} 90+\left(\frac{\gamma_{x y}}{2}\right)(2 \sin 90 \cos 90)
\end{aligned}
$$

$$
\begin{aligned}
& \varepsilon_{x}=\varepsilon_{a} \\
& \varepsilon_{y}=\varepsilon_{b} \\
& \gamma_{x y}=2 \varepsilon_{b}-\varepsilon_{a}-\varepsilon_{c}
\end{aligned}
$$

## Example

The following results are obtained from a $60^{\circ}$ strain gauge rosette:
Strain in direction of strain gauge $\mathrm{A}=750 \mu$;
Strain in direction of SG B, $60^{\circ}$ to A $=350 \mu$;
Strain in direction of SG C, $120^{\circ}$ to $\mathrm{A}=100 \mu$.
Determine the principal strains and their directions.

$$
\varepsilon_{a}=\varepsilon_{x}=750 \mu
$$

$$
\begin{aligned}
& \varepsilon_{b}=\varepsilon_{x} \cos ^{2} 60+\varepsilon_{y} \sin ^{2} 60+\left(\frac{\gamma_{x y}}{2}\right)(2 \sin 60 \cos 60) \\
& \varepsilon_{b}=\varepsilon_{x}(0.25)+\varepsilon_{y}(0.75)+\left(\frac{\gamma_{x y}}{2}\right)(0.433)
\end{aligned}
$$

$$
\varepsilon_{b}=\varepsilon_{x} \cos ^{2} 120+\varepsilon_{y} \sin ^{2} 120+\left(\frac{\gamma_{x y}}{2}\right)(2 \sin 120 \cos 120)
$$

$$
\varepsilon_{b}=\varepsilon_{x}(0.25)+\varepsilon_{y}(0.75)+\left(\frac{\gamma_{x y}}{2}\right)(-0.433)
$$

$$
\varepsilon_{\mathrm{x}}=750, \varepsilon_{\mathrm{y}}=50 \text { and } \gamma_{\mathrm{xy}}=289
$$



$$
\Rightarrow \left\lvert\, \begin{gathered}
\varepsilon_{y}=50 \mu \\
\gamma_{x y}=289 \mu
\end{gathered}\right.
$$

$$
\begin{gathered}
\theta=11.2^{\circ} \text { and } 101.2^{\circ} \\
\varepsilon_{11.2}=779 \mu \varepsilon \\
\varepsilon_{101.2}=21 \mu \varepsilon
\end{gathered}
$$

$$
\begin{aligned}
& {[\varepsilon]=\left[\begin{array}{ccc}
\varepsilon_{x x} & \frac{1}{2} \gamma_{y x} & \frac{1}{2} \gamma_{z x} \\
\frac{1}{2} \gamma_{x y} & \varepsilon_{y y} & \frac{1}{2} \gamma_{z y} \\
\frac{1}{2} \gamma_{x z} & \frac{1}{2} \gamma_{y z} & \varepsilon_{z z}
\end{array}\right]=\left[\begin{array}{ccc}
750 & 289 / 2 & 0 \\
289 / 2 & 50 & 0 \\
0 & 0 & 0
\end{array}\right] \mu} \\
& \text { Eigen_Values } \Rightarrow\left[\begin{array}{ccc}
779 & 0 & 0 \\
0 & 21 & 0 \\
0 & 0 & 0
\end{array}\right] \\
& \text { Eigen_Vectors } \Rightarrow\left[\begin{array}{ccc}
0.981 & 0.1945 & 0 \\
-0.1945 & 0.981 & 0 \\
0 & 0 & 1
\end{array}\right]
\end{aligned}
$$

## Example

(A) Using the transformation equations define the maximum and minimum principal strains, maximum shearing strain and principal angles given $\varepsilon_{X}=3500 \mu ; \varepsilon_{Y}=700 \mu$ and $\gamma_{X Y}=-1050 \mu$
(B) Repeat using the Mohr's circle.
$[\varepsilon]=\left[\begin{array}{ccc}3500 & -1050 / 2 & 0 \\ -1050 / 2 & 700 & 0 \\ 0 & 0 & 0\end{array}\right] \mu$
Eigen_Values $\Rightarrow\left[\begin{array}{ccc}3595.2 & 0 & 0 \\ 0 & 604.8 & 0 \\ 0 & 0 & 0\end{array}\right] \mu$
Eigen_Vectors $\Rightarrow\left[\begin{array}{ccc}0.984 & -0.178 & 0 \\ 0.178 & 0.984 & 0 \\ 0 & 0 & 1\end{array}\right]$
(d) Out-of-Plane Maximum Shear Strain

$$
\gamma_{\operatorname{Max}}=\left(\varepsilon_{1}-\varepsilon_{3}\right)=3595.2 \mu-0 \mu=3595.2 \mu
$$

$\operatorname{ArcCos}($ angle $)=0.984$
Angle $=10.28$ degress
(c) In-Plane Maximum Shear Strain

$$
\gamma_{M a x}=\left(\varepsilon_{1}-\varepsilon_{2}\right)=3595.2 \mu-604.8 \mu=2990.4 \mu
$$



## Example

The state of stress at a point in a structural member is determined to be as shown. Knowing that for this material $E=210 \mathrm{GPa}$ and $v=0.3$, use the Mohr's circle to determine: (1) the principal stresses; (2) the in-plane maximum shear stress; (3) the absolute maximum shear stress; (4) principal angles; (5) the strains and the principal strains; (6) the maximum shear strain; (7) the principal angles.


$$
[\sigma]=\text { Stress_Tensor }=\left[\begin{array}{ccc}
-56 & 11.2 & 0 \\
11.2 & -14 & 0 \\
0 & 0 & 0
\end{array}\right] M P a \quad \text { Eigen_Values } \Rightarrow\left[\begin{array}{ccc}
0 & 0 & 0 \\
0 & -11.2 & 0 \\
0 & 0 & -58.8
\end{array}\right] M P a
$$

2. In-Plane Maximum Shear Stress

$$
\tau_{M a x}=R=23.8 M P a
$$

## 3. The Absolute Maximum Shear Stress (Out of plane)


4. Angle between the $x$-axis and the Principal Stresses

$$
(\text { In_Plane })=\tan \left(2 \theta_{P}\right)=\frac{\tau_{x y}}{\left(\sigma_{x}-\sigma_{y} y / 2\right.}=\frac{11.2}{(-56-(-14)) / 2}
$$

$$
\tan \left(2 \theta_{P}\right)=\frac{11.2}{-21}=-0.533
$$

$$
2 \theta_{P}=-28.07 \mathrm{deg}
$$

For $\sigma_{1}=0$
ArcCos(angle) $=0.0$
Angle $=90$ degress

For $\sigma_{2}=-11.2 \mathrm{MPa}$
ArcCos(angle) $=0.2425$
Angle $=76$ degress

For $\sigma_{3}=-58.8 \mathrm{MPa}$
ArcCos(angle) $=0.9701$
Angle $=14$ degress
5. Strains and Principal Strains

$$
\text { Eigen_Values } \Rightarrow\left[\begin{array}{ccc}
100 & 0 & 0 \\
0 & 30.8 & 0 \\
0 & 0 & -264.0
\end{array}\right] \mu
$$

$$
\begin{aligned}
& \varepsilon_{x}=\frac{\sigma_{x}}{E}-v \frac{\sigma_{y}}{E}-v \frac{\sigma_{z}}{E} \quad \varepsilon_{x}=\frac{-56}{210000}-(0.3) \frac{-14}{210000}=(-266.6+20) \mu=-246.6 \mu \\
& \varepsilon_{y}=-v \frac{\sigma_{x}}{E}+\frac{\sigma_{y}}{E}-v \frac{\sigma_{z}}{E} \quad \varepsilon_{y}=-(0.3) \frac{-56}{210000}+\frac{-14}{210000}=(80-66.6) \mu=13.4 \mu \\
& \varepsilon_{z}=-v \frac{\sigma_{x}}{E}-v \frac{\sigma_{y}}{E}+\frac{\sigma_{z}}{E} \quad \varepsilon_{z}=-(0.3) \frac{-56}{210000}-(0.3) \frac{-14}{210000}=(80+20) \mu=100 \mu \\
& \gamma_{X Y}=\frac{\tau_{X Y}}{G} \gamma_{x y}=\frac{\tau_{x y}}{G}=\frac{11.2}{80770}=139 \mu \\
& G=\frac{E}{2(1+v)}=\frac{210}{2(1+0.3)}=80.77 \mathrm{GPa} \\
& {[\varepsilon]=\left[\begin{array}{ccc}
\varepsilon_{x x} & \frac{1}{2} \gamma_{y x} & \frac{1}{2} \gamma_{z x} \\
\frac{1}{2} \gamma_{x y} & \varepsilon_{y y} & \frac{1}{2} \gamma_{z y} \\
\frac{1}{2} \gamma_{x z} & \frac{1}{2} \gamma_{y z} & \varepsilon_{z z}
\end{array}\right]=\left[\begin{array}{ccc}
-246.6 & 139 / 2 & 0 \\
139 / 2 & 13.4 & 0 \\
0 & 0 & 100
\end{array}\right] \mu} \\
& \varepsilon_{\text {Average }}=\frac{(-246.6+13.4)}{2}=-116.6 \mu \\
& \begin{array}{c}
R=\sqrt{\left(\frac{-246.6-13.4}{2}\right)^{2}+\left(\frac{139}{2}\right)^{2}}=147.4 \mu \\
\varepsilon_{1}=30.8 \mu \\
\varepsilon_{2}=-264 \mu
\end{array} \\
& \varepsilon_{1,2}=\frac{\left(\varepsilon_{x}+\varepsilon_{y}\right)}{2} \pm \sqrt{\left(\frac{\varepsilon_{x}-\varepsilon_{y}}{2}\right)^{2}+\left(\frac{\gamma_{x y}}{2}\right)^{2}} \\
& \varepsilon_{1,2}=\varepsilon_{\text {Average }} \pm R
\end{aligned}
$$

6. Maximum Shear Strain and Absolute Maximum Shear Strain

$$
\begin{aligned}
& \frac{\gamma_{\text {Max }}}{2}(\text { in_plane })=R=\sqrt{\left(\frac{-246.6-13.4}{2}\right)^{2}+\left(\frac{139}{2}\right)^{2}}=147.4 \mu \\
& \gamma_{\text {Max }}=294.8 \mu
\end{aligned}
$$

$$
\frac{\gamma_{\operatorname{Max}}}{2}(\text { out_of_plane })=\frac{\varepsilon_{1}-\varepsilon_{3}}{2}=\frac{100-(-264)}{2}=182 \mu
$$

$$
\gamma_{\operatorname{Max}}=364 \mu
$$

## 7. In-plane and Out-of-plane angles

$$
\tan 2 \theta_{P}=\frac{\gamma_{x y} / 2}{\left(\varepsilon_{x}-\varepsilon_{y}\right) / 2}=\frac{\gamma_{x y}}{\varepsilon_{x}-\varepsilon_{y}}=\frac{139}{(-246.6-13.4)}=-0.534
$$



$$
\begin{aligned}
& \tan \left(2 \theta_{P}\right)=-0.533 \\
& 2 \theta_{P}=-28.07 \mathrm{deg}
\end{aligned}
$$

For $\varepsilon_{3}=-264 \mu$ ArcCos(angle) $=0.9701$ Angle $=14$ degress

