

QNO1:

Define Separation of variables,
Separate the variables in Schrodinger's
wave equation for hydrogen atom?

In mathematics separation of variables
is any of several methods for
solving ordinary and partial differential
equations in which algebra allows
one to re-write an equation
so that each of two variables
occurs on different side of
equation.

Schrodinger's wave equation for
Hydrogen atom:

For hydrogen atom one proton,
one electron and the electrostatic
potential that holds them together
The potential energy in this
case is.

$$V = \frac{e^2}{4\pi\epsilon_0 r}$$

Which is the attractive potential between charges of $+e$ and $-e$ separated by a distance r .

Now this potential looks quite simple. But notice that it is a function of $a \cdot r$, not x or (xyz) .

One approach would be to express v in terms of (xyz) , where

$$x^2 + y^2 + z^2 = r^2$$

In some cases you might be able to get away with that but here you would make the problem quite difficult (the square root of sum of squares is a pain to deal with).

The appropriate approach is to let the symmetric potential tell us to use spherical polar coordinates.

In spherical polar coordinates r , is length of radius vector from the origin to a point (xyz)

$$r = \sqrt{x^2 + y^2 + z^2}$$

Another coordinate is θ , the angle between the radius vector and +z axis

$$\theta = \cos^{-1} \left(\frac{z}{\sqrt{x^2 + y^2 + z^2}} \right)$$

The third coordinate (because we are in three dimensions, we need three coordinates) is ϕ the angle between projection of the radius vector into the xy plane and the +x axis.

$$\phi = \tan^{-1} \left[\frac{y}{x} \right]$$

Note that we have three coordinates (r, θ, ϕ) which we have relate to (x, y, z) above. we can also express (x, y, z) in terms of (r, θ, ϕ) :

$$x = r \sin \theta \cos \phi \quad y = r \sin \theta \sin \phi \quad z = r \cos \theta$$

Now, we can re-write Schrodinger's equation

$$\frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} + \frac{\partial^2 \psi}{\partial z^2} + \frac{2m}{\hbar^2} (E - V) \psi = 0$$

In spherical polar coordinate as

$$\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial \psi}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial \psi}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 \psi}{\partial \phi^2}$$

$$+ \frac{2m}{\hbar^2} (E - V) = 0.$$

if we plug in our potential V and multiply both sides by $r^2 \sin^2 \theta$, we get

$$\sin^2 \theta \frac{\partial}{\partial r} \left(r^2 \frac{\partial \psi}{\partial r} \right) + \sin \theta \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial \psi}{\partial \theta} \right) + \frac{\partial^2 \psi}{\partial \phi^2} + \frac{2m r^2 \sin^2 \theta}{\hbar^2} (E - V) \psi = 0$$

$$\left(\frac{e^2}{4\pi\epsilon_0 r} + E \right) = 0$$

Let's talk about this equation
it looks nasty, but we will
see that it's not quite bad
as it looks

This equation gives us the wave function ψ for electron in hydrogen atom. if we can solve for ψ in principal we know every thing that is to know about the hydrogen atom. When we solved equation in one dimensions we found that one Quantum number was necessary to describe our systems. For example, in bohr's atom, the electron moves in an orbit but we need only one quantum number. Here in three dimensions and with three boundary conditions we will find that we need three quantum numbers to describe our electron.

Beiser at the end of this section tells what Quantum numbers for hydrogen atom are and give their possible values, but untill we see where they come from and what they mean and they are not of much use to us.

Another comment, we are really solving Schrodinger's equation for electron in a hydrogen atom aren't we. Nevertheless we talk about doing the "hydrogen atom" because our solution will provide us with much of what we need to know about Hydrogen.

Separation of Variables:

We now discuss the technique for solving our equation for the electron in hydrogen atom

Here are some things we like in decreasing order of liking

- Linear algebraic equations
- Coupled algebraic equations
(i.e. xy multiplied together)
- Linear differential equations
- Coupled differential equations

When we have the equation like one above, we like to see if we can "separate" the variables "split" the equations into different parts

with only one variable in each part
Our problem will then be much
simplified if we can write

$$\psi(r, \theta, \varphi) = R(r) \Theta(\theta) \Phi(\varphi) = R \Theta \Phi$$

Let assume we can write ψ like
that and see where it leads us
if our assumption works, in the
orderly work of mathematics we
know it must have been right.

Here, How we take the partial
derivatives in Schrodinger's equation:

$$\frac{\partial \psi}{\partial r} = \Theta \Phi \frac{\partial R}{\partial r} = \Theta \Phi \frac{dR}{dr}$$

$$\frac{\partial \psi}{\partial \theta} = R \Phi \frac{\partial \Theta}{\partial \theta} = R \Phi \frac{d\Theta}{d\theta}$$

$$\frac{\partial^2 \psi}{\partial \varphi^2} = R \Theta \frac{\partial^2 \Phi}{\partial \varphi^2} = R \Theta \frac{d^2 \Phi}{d\varphi^2}$$

Where the partial derivatives become
full derivatives because R , Θ and Φ
depends on r , θ and φ only.

To separate the variables, plug $\psi = R(\theta)\Phi$ into Schrodinger's equation and divide by $R(\theta)\Phi$. The result is

$$\frac{\sin^2\theta}{R} \frac{d}{dr} \left[r^2 \frac{dR}{dr} \right] + \frac{\sin\theta}{\Theta} \frac{d}{d\theta} \left[\sin\theta \frac{d\Theta}{d\theta} \right] + \frac{1}{\Phi} \frac{d^2\Phi}{d\phi^2} +$$

$$\frac{2mr^2 \sin^2\theta}{h^2} \left(\frac{e^2}{4\pi\epsilon_0 r} + E \right) = 0$$

Now we have separated out the ϕ variable, the term

$$\frac{1}{\Phi} \frac{d^2\Phi}{d\phi^2}$$

is a function of ϕ only. Let us put it over on the right hand side of equation. This gives us.

$$\frac{\sin^2\theta}{R} \frac{d}{dr} \left[r^2 \frac{dR}{dr} \right] + \frac{\sin\theta}{\Theta} \frac{d}{d\theta} \left[\sin\theta \frac{d\Theta}{d\theta} \right] +$$

$$\frac{2mr^2 \sin^2\theta}{h^2} \left(\frac{e^2}{4\pi\epsilon_0 r} + E \right) = -\frac{1}{\Phi} \frac{d^2\Phi}{d\phi^2}$$

Notice the form of equation it is in the form $f(r, \theta) = g(\phi)$.

Where f is function of r and θ only and g is function of ϕ only

Under what condition, this can be possible?

If an r or θ shows up on the LHS, it can't be satisfied because r and θ never show up on the RHS.

Similarly, if a φ shows up on RHS, the equation cannot be satisfied because φ never shows up on the LHS. The only way for the equation to be satisfied is for

$$f(r, \theta) = \text{a constant, independent of } r, \theta$$
$$\text{and } \varphi = g(\varphi)$$

Notice what we have done,

We have taken the one nasty equation for r, θ and φ , and separated into two equations one is r and θ and other in φ only.

At this point, we will begin quoting some results from differential equations

As we have seen several times before in this course some differential equations can be solved only if certain conditions are satisfied. These conditions have led to quantum numbers. They will do so in this case also.

Let me follow Beiser's development and anticipate some of requirements.

It turns out that the constant in the last equation on previous page is square of an integer. Thus we can write the RHS of Beiser's eq as

$$-\frac{1}{\Phi} \frac{d^2 \Phi}{d\varphi^2} = m_l^2$$

We will hear more about this m_l later.

If we set LHS equal to m_l^2 , divide by $\sin^2 \theta$ and rearrange, we get.

$$r \frac{1}{R} \frac{d}{dr} \left[r^2 \frac{dR}{dr} \right] + \frac{2mr^2}{h^2} \left(\frac{e^2}{4\pi\epsilon_0 r} + E \right) = \frac{m_l^2}{\sin^2 \theta} \frac{1}{\sin \theta} \frac{d}{d\theta} \left[\sin \theta \frac{d\Theta}{d\theta} \right]$$

Once again we have separated variables

The LHS is the function of r only
and the RHS is a function of θ only

Again the only way to satisfy
this equation is for LHS = a
Constant = RHS.

Conditions on the constant will
arrive from the solution of
differential equations. In this
case, we call the constant $l(l+1)$

We conclude this section by taking
our three differential equations
for r , θ and ϕ rearranging them
slightly and writing them in
the form

$$\frac{d^2 \phi}{d\phi^2} + m_l^2 \phi = 0$$

$$\frac{1}{\sin \theta} \frac{d}{d\theta} \left(\sin \theta \frac{d\Theta}{d\theta} \right) + \left[l(l+1) - \frac{m_l^2}{\sin^2 \theta} \right] \Theta = 0$$

$$\frac{1}{r^2} \frac{d}{dr} \left(r^2 \frac{dR}{dr} \right) + \left[\frac{2m}{\hbar^2} \left(\frac{e^2}{4\pi\epsilon_0 r} \right) + E - \frac{l(l+1)}{r^2} \right] R = 0$$

All of above seems like a lot of work but we have separated our big partial differential equation into three simple ones each of which is a function of a single variable and each of which is much easier to handle.