

## The Biot-Savart Law:

Stationary charges produce electric fields that are constant in time; hence the term electrostatics. Steady currents produce magnetic fields that are constant in time. That's why the theory of steady currents is called magnetostatics.

Stationary charges  $\Rightarrow$  constant electric fields : electrostatics

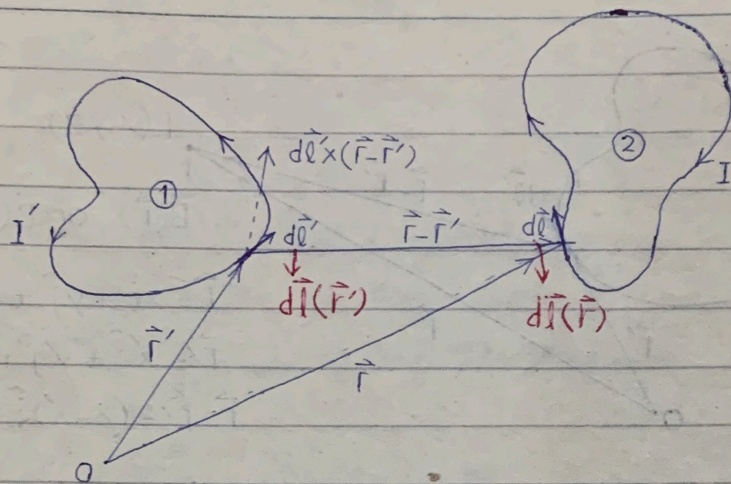
Steady currents  $\Rightarrow$  constant magnetic fields : magnetostatics

In 1820, Oersted first discovered that currents produce magnetic fields. Just after a few weeks of this discovery, Ampere presented the results of a series of experiments that may be generalized and expressed in modern mathematical language as

$$\vec{B} = \frac{\mu_0}{4\pi} I I' \oint \oint \frac{d\vec{l} \times [d\vec{l}' \times (\vec{r} - \vec{r}')] }{|\vec{r} - \vec{r}'|^3}$$

This expression can be understood with reference to the Figure.





The force  $\vec{F}$  is the force exerted on circuit ② due to the influence of circuit ①. By definition,

$$\frac{\mu_0}{4\pi} = 10^{-7} \text{ N/A}^2 \text{ (in mks units)}$$

It plays the same role as that of  $1$  in electrostatics, and is a constant. ' $\mu_0$ ' is called permeability of free space.

Now from our knowledge about current carrying conductors, we know that the force experienced by a current carrying element  $I d\vec{l}$  of a closed circuit when placed in a magnetic field of magnetic induction  $\vec{B}$  is given by

$$d\vec{F}(\vec{r}) = I d\vec{l} \times \vec{B}(\vec{r})$$

$$\vec{F} = \oint I d\vec{l} \times \vec{B}(\vec{r})$$

Force experienced by whole closed loop

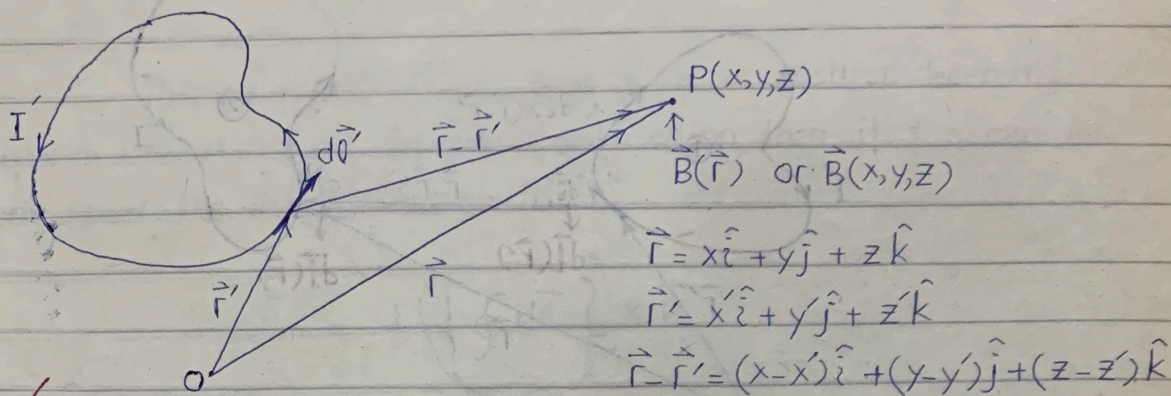
Compare it with

$$\vec{F} = \frac{\mu_0}{4\pi} \oint \oint \frac{(I d\vec{l}) \times I' [d\vec{l}' \times (\vec{r} - \vec{r}')] }{|\vec{r} - \vec{r}'|^3}, \text{ it is clear that}$$

$$\vec{B}(\vec{r}) = \frac{\mu_0}{4\pi} \oint \frac{I' d\vec{l}' \times (\vec{r} - \vec{r}')}{|\vec{r} - \vec{r}'|^3} = \vec{B}(P)$$

This is the general form of Biot-Savart Law, which gives us the value of magnetic induction  $\vec{B}$  at any point 'P', a distance ' $\vec{r}$ ' away from the origin and ' $\vec{r} - \vec{r}'$ ' away from a conducting loop through which current ' $I$ ' is flowing and  $I' d\vec{l}'$  is its current carrying element.





? Check it! ✓

The direction of  $\vec{B}(\vec{r})$  can be found by right hand rule, i.e., rotate  $I d\vec{l}'$  to coincide with  $(\vec{r} - \vec{r}')$  through smallest angle, curl right hand fingers along the direction of rotation and the right hand thumb will point in the direction of  $\vec{B}(\vec{r})$ .

For surface currents, the Biot-Savart law becomes

$$I d\vec{l}' \rightarrow \vec{K}(\vec{r}') da'$$

$$\therefore \vec{B}(\vec{r}) = \frac{\mu_0}{4\pi} \int_S \frac{\vec{K}(\vec{r}') \times (\vec{r} - \vec{r}')}{|\vec{r} - \vec{r}'|^3} da'$$

$$\vec{B}(\vec{r}) = \frac{\mu_0}{4\pi} \int \frac{I d\vec{l}' \times (\vec{r} - \vec{r}')}{|\vec{r} - \vec{r}'|^3}$$

and For the volume currents

$$I d\vec{l}' \rightarrow \vec{J}(\vec{r}') dv' \text{ or } \vec{J}(\vec{r}') d\vec{r}'$$

For open circuit.

$$\vec{B}(\vec{r}) = \frac{\mu_0}{4\pi} \int_V \frac{\vec{J}(\vec{r}') \times (\vec{r} - \vec{r}')}{|\vec{r} - \vec{r}'|^3} d\vec{r}'$$

$$I = I(\vec{r}')$$

$$d\vec{l}' = d\vec{l}(\vec{r}')$$

Differential form of Equations are

$$d\vec{B} = \frac{\mu_0}{4\pi} \frac{I d\vec{l}' \times (\vec{r} - \vec{r}')}{|\vec{r} - \vec{r}'|^3}$$

$$d\vec{B} = \frac{\mu_0}{4\pi} \frac{\vec{K}(\vec{r}') \times (\vec{r} - \vec{r}')}{|\vec{r} - \vec{r}'|^3} da'$$

$$d\vec{B} = \frac{\mu_0}{4\pi} \frac{\vec{J}(\vec{r}') \times (\vec{r} - \vec{r}')}{|\vec{r} - \vec{r}'|^3} d\vec{r}'$$



Proof:  $\vec{\nabla} \cdot \vec{B} = 0$

$$\text{As, } \vec{B}(\vec{r}) = \frac{\mu_0}{4\pi} \int \frac{\vec{J}(\vec{r}') \times (\vec{r} - \vec{r}')}{|\vec{r} - \vec{r}'|^3} d\tau'$$

$$\text{As, } \vec{r} - \vec{r}' = \vec{r}$$

$$\therefore \vec{B}(\vec{r}) = \frac{\mu_0}{4\pi} \int \frac{\vec{J}(\vec{r}') \times \vec{r}}{r^3} d\tau'$$

$$\Rightarrow \vec{B}(\vec{r}) = \frac{\mu_0}{4\pi} \int \frac{\vec{J}(\vec{r}') \times \hat{r}}{r^2} d\tau'$$

$$\text{As, } \vec{r} = x\hat{i} + y\hat{j} + z\hat{k}$$

$$\vec{r}' = x'\hat{i} + y'\hat{j} + z'\hat{k}$$

$$\vec{r} = \vec{r} - \vec{r}' = (x-x')\hat{i} + (y-y')\hat{j} + (z-z')\hat{k}$$

So,

$\vec{B}$  is a function of  $(x, y, z)$ ,

$\vec{J}$  is a function of  $(x', y', z')$ ,

$$d\tau' = dx' dy' dz'$$

$$\vec{\nabla} \cdot \vec{B} = \frac{\mu_0}{4\pi} \vec{\nabla} \cdot \int \frac{\vec{J}(\vec{r}') \times \hat{r}}{r^2} d\tau'$$

$$\text{As, } \vec{\nabla} = \frac{\partial}{\partial x} \hat{i} + \frac{\partial}{\partial y} \hat{j} + \frac{\partial}{\partial z} \hat{k}$$

$$\therefore \vec{\nabla} \cdot \vec{B} = \frac{\mu_0}{4\pi} \int \vec{\nabla} \cdot \left[ \frac{\vec{J}(\vec{r}') \times \hat{r}}{r^2} \right] d\tau'$$

Using vector identity,

$$\vec{\nabla} \cdot (\vec{A} \times \vec{B}) = \vec{B} \cdot (\vec{\nabla} \times \vec{A}) - \vec{A} \cdot (\vec{\nabla} \times \vec{B})$$

$$\therefore \vec{\nabla} \cdot \vec{B} = \frac{\mu_0}{4\pi} \int \frac{\hat{r}}{r^2} \cdot (\vec{\nabla} \times \vec{J}(\vec{r}')) d\tau' - \frac{\mu_0}{4\pi} \int \vec{J}(\vec{r}') \cdot (\vec{\nabla} \times \frac{\hat{r}}{r^2}) d\tau'$$

Because  $\vec{J}(\vec{r}')$  doesn't depend on the unprimed variables.

$$\text{As, } \frac{\hat{r}}{r^2} = -\vec{\nabla} \left( \frac{1}{r} \right)$$

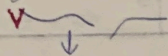
$$\therefore \vec{\nabla} \times \frac{\hat{r}}{r^2} = -\vec{\nabla} \times \vec{\nabla} \left( \frac{1}{r} \right) = 0 \quad (\text{since, curl of a gradient is zero})$$

$$\cancel{\vec{\nabla} \cdot \frac{(\vec{r} - \vec{r}')}{|\vec{r} - \vec{r}'|^3}} = -\vec{\nabla} \cdot \left( \frac{1}{|\vec{r} - \vec{r}'|} \right)$$



$\therefore \vec{\nabla} \cdot \vec{B} = 0 \equiv$  Maxwell's second Equation

$$\Rightarrow \int (\vec{\nabla} \cdot \vec{B}) d\tau = 0$$



Applying Divergence theorem,

$$\oint \vec{B} \cdot d\vec{a} = 0$$

i.e, the Magnetic<sup>s</sup> flux through any closed surface is always equal to zero.

$\Rightarrow$  Magnetic monopoles do not exist.