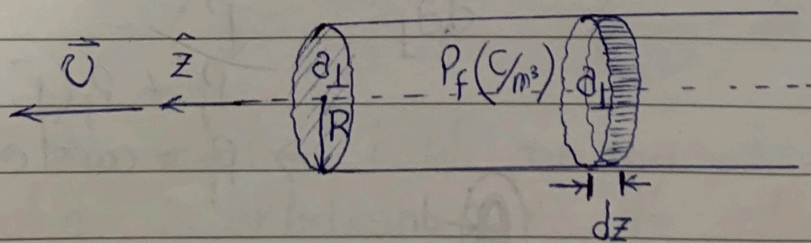


## Volume Currents:



$$d\vec{a}_\perp = r dr d\phi \hat{z}$$

$$d\vec{a}_\perp = r dr d\phi$$

$$dV = \int_0^R r dr \int_0^{2\pi} d\phi$$

$$dq_f = \rho_f dV = \rho_f a_\perp dz = \pi R^2$$

$$= \rho_f \pi R^2 dz$$

DALMATIAN

**Date:**

Consider an infinitely long circular conducting rod of radius  $R$ , with cross-sectional area  $a_{\perp} = \pi R^2$ . For simplicity, let us assume the rod is initially charged with constant free volume charge density  $\rho_f$ , and  $\rho_f \neq \rho_f(\vec{r}, t)$ . We again place a potential difference  $\Delta V$  across the ends of the rod, and a volume current  $I$  flows in the rod.

$$(I \text{ or}) I_f = \frac{dq_f}{dt} = \frac{\rho_f dz}{dt} = \rho_f (\pi R^2) \frac{dz}{dt} = \rho_f (\pi R^2) v$$

$$I(z, t) = \rho_f (\pi R^2) v(z, t)$$

In general,

$$\vec{I}(\vec{r}, t) = \rho_f (\pi R^2) \vec{v}(\vec{r}, t)$$

$$\vec{J}(\vec{r}, t) \equiv \frac{\vec{I}(\vec{r}, t)}{\pi R^2} = \rho_f \vec{v}(\vec{r}, t)$$

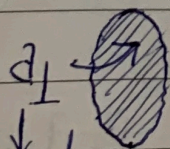
(Free) Volume current density ; SI Unit: Amperes/m<sup>2</sup>  
Note: not Amperes/m<sup>3</sup> !!!

$$\text{or } \vec{J}(\vec{r}, t) = \frac{\vec{I}(\vec{r}, t)}{a_{\perp}} = \rho_f \vec{v}(\vec{r}, t)$$

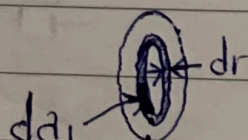
The differential form is

$$\vec{J}(\vec{r}, t) = \frac{d\vec{I}(\vec{r}, t)}{da_{\perp}} = \rho_f \vec{v}(\vec{r}, t)$$

$\rho_f \neq \rho_f(\vec{r}, t)$   
 $\rightarrow \rho_f = \text{constant}$



cross-sectional area of conductor



infinitesimal cross-sectional area element of  $a_{\perp}$

Date:

$$\vec{I}(\vec{r}, t) = \int_{S_1} \vec{J}(\vec{r}, t) \cdot d\vec{a}_1$$

In general,  $\vec{I}(\vec{r}, t) = \int_{S_1} \vec{J}(\vec{r}, t) \cdot d\vec{a}$

Current through any closed surface,  $S$  is

$$\vec{I}(\vec{r}, t) = \oint_S \vec{J}(\vec{r}, t) \cdot d\vec{a}$$

the rate at which the charge is decreasing in volume,  $V$ .

$$\oint_S \vec{J}(\vec{r}, t) \cdot d\vec{a} = - \frac{dq_f}{dt} = - \frac{d}{dt} \int_V \rho_f(\vec{r}, t) d\tau$$

current flowing out through

surface,  $S$  (bounding the volume,  $V$ )

decrease in the charge density in volume,  $V$ .

Electric charge conservation: The electric charge can neither be created, nor destroyed.

The electric charge is conserved, whatever electric charge flows out of or flows into the surface  $S$  must come from or go into the volume  $V$  respectively.

↓ Applying Divergence theorem,

$$\int_V (\vec{\nabla} \cdot \vec{J}) d\tau = \int_V - \frac{\partial \rho_f(\vec{r}, t)}{\partial t} d\tau$$

Since, this relation holds for any volume ' $V$ ' (i.e., the choice of volume is arbitrary) so the integrands on the LHS and RHS must be equal.

Date:

$$\nabla \cdot \vec{J}(\vec{r}, t) = - \frac{\partial \rho(\vec{r}, t)}{\partial t}$$

= Equation of continuity  
(Electric charge conservation)

Under static conditions, (For steady currents)

No time dependence.

If  $\rho \neq \rho(\vec{r}, t)$   $\nabla \cdot \vec{J}(\vec{r}) + \frac{\partial \rho(\vec{r})}{\partial t} = 0$   
If  $\rho = \text{constant} \Rightarrow$  either  $\rho \neq \rho(\vec{r}, t)$  or  $\rho = \rho(\vec{r})$

$$\therefore \frac{\partial \rho}{\partial t} = 0$$

$$\Rightarrow \nabla \cdot \vec{J}(\vec{r}) = 0$$

$$\nabla \cdot \vec{J}(\vec{r}) = 0$$

$$\int_V \nabla \cdot \vec{J}(\vec{r}, t) d\tau = 0$$

Applying DT,  
 $\oint_S \vec{J}(\vec{r}, t) \cdot d\vec{a} = 0$

Summary:

volume current flux through any closed surface, S

line charge density

Surface charge density

Volume charge density

$$Q = \int_C \lambda(\vec{r}) dl$$

$$Q = \int_S \sigma(\vec{r}) da$$

$$Q = \int_V \rho(\vec{r}) d\tau$$

line current  
 $I$  (A)

Surface current density  
 $\vec{K}(\vec{r})$  (A/m)

Volume current density  
 $\vec{J}(\vec{r})$  (A/m<sup>2</sup>)

$$I = \int_C \vec{K}(\vec{r}) d\vec{l}$$

$$I = \int_S \vec{J}(\vec{r}) \cdot d\vec{a}$$

Bending the motion of charges along a curve is typical behavior of magnetic forces.

Date:

### Example 5.4:

(a)

Since,  $I$  is uniformly distributed over cross-sectional area  $A$  of the wire  $\Rightarrow \vec{J}$  must also be uniform.

$$J = \frac{I}{A} = \frac{I}{\pi a^2}$$

(b)

$$J(s) = ks$$

$s$  = radial distance from cylindrical symmetry axis

$k$  = constant (SI Unit:  $A/m^2$ )

$$\vec{J}(s) = ks \hat{z}$$

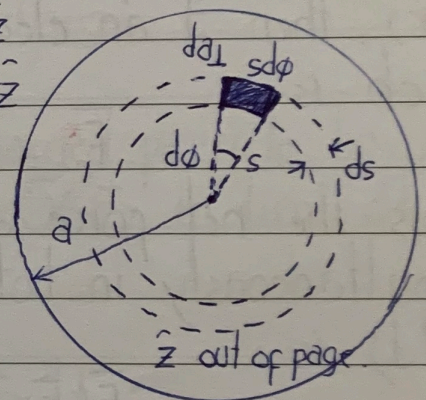
$$d\vec{a} = da_{\perp} \hat{z} \\ = s ds d\phi \hat{z}$$

$$I = \int_{S_{\perp}} \vec{J}(s) \cdot d\vec{a}$$

$$= \int_{S_{\perp}} ks (s ds d\phi) \underbrace{\hat{z} \cdot \hat{z}}_{=1}$$

$$= k \int_0^a s^2 ds \int_0^{2\pi} d\phi$$

$$= k \frac{a^3}{3} 2\pi = \frac{2\pi}{3} k a^3$$



cross-sectional view  
of the wire.