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Free Convection Flow along a vertical plate:

This article investigates the free convection flow of viscous incompressible fluid along a heated vertical plate with uniform surface temperature. With the appropriate transformations, the boundary layer equations governing the flow are reduced to local nonsimilarity equations valid in free convection regime. A group of transformation is also introduced to reduce the boundary layer equations to set of ~~local~~ nonsimilar equations.

The problem of free convection flows past a vertical surfaces has been studied extensively because of its wide application in industry.

Mathematical formulation:

Consider a steady two dimensional flow of a viscous incompressible fluid along a semi-infinite vertical plate. Under the Boussinesq approximation the flow is governed by the following boundary layer equations:

Equation of continuity (Law of conservation of mass).

$$u \frac{\partial u}{\partial x} + v \frac{\partial v}{\partial y} = 0 \quad \longrightarrow \quad (1)$$

Momentum Equation (Navier-Stokes Equation).

$$u \frac{\partial u}{\partial x} + v \frac{\partial v}{\partial y} = \nu \frac{\partial^2 u}{\partial y^2} + g \beta_E (T - T_\infty) + g \beta_C (C - C_\infty) \quad \longrightarrow \quad (2)$$

(2)

Energy Equation:

$$u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} = \alpha \frac{\partial^2 T}{\partial y^2} \longrightarrow (3)$$

Mass concentration Equation

$$u \frac{\partial c}{\partial x} + v \frac{\partial c}{\partial y} = D_m \frac{\partial^2 c}{\partial y^2} \longrightarrow (4)$$

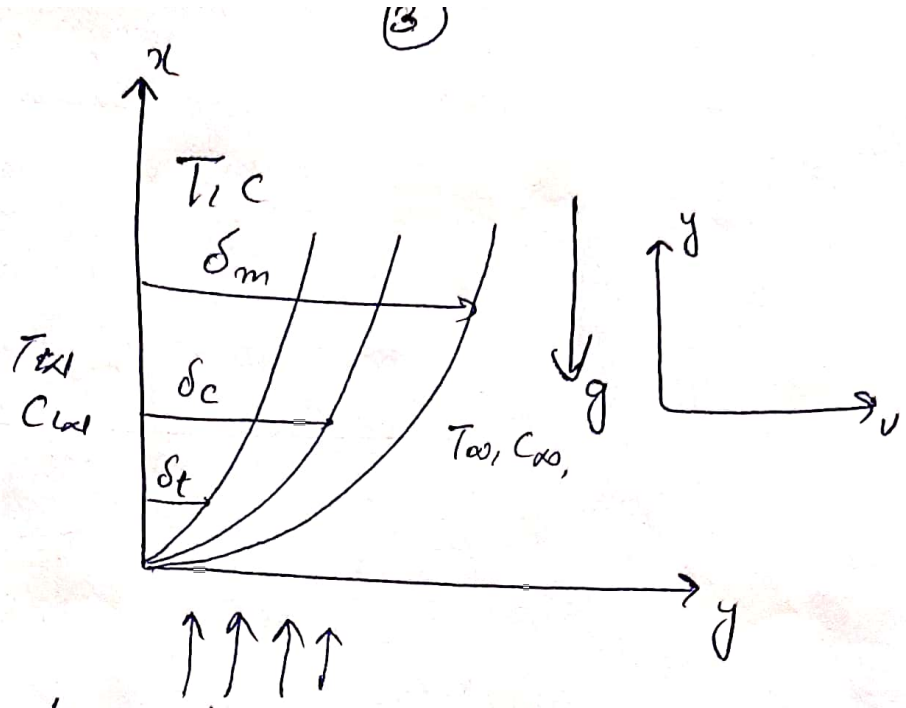
Subject to the boundary Conditions:

$$u=0, v=0, T=T_w, C=C_w \text{ at } y=0 \longrightarrow (5)$$
$$u \rightarrow 0, T \rightarrow T_\infty, C \rightarrow C_\infty, \text{ as } y \rightarrow \infty.$$

Here, u , and v are the velocity components along x and y directions respectively. The symbols ν , g , β_t and β_c are known as kinematic viscosity, gravitation acceleration (gravitational force), thermal expansion coefficient, & concentration expansion coefficient respectively.

Here, α and D_m denote thermal diffusion coefficient, and mass diffusion coefficient. where $\alpha = \frac{k}{\rho c_p}$, $\nu = \frac{\mu}{\rho}$; ρ is the density of fluid, c_p is specific heat at constant pressure, μ is dynamic viscosity.

~~and~~



T_w, C_w wall temperature and mass concentration. T_{∞}, C_{∞} are air free stream temperature and mass concentration. $\delta_t, \delta_c, \delta_m$ are known as thermal boundary layer thickness, concentration boundary layer thickness and momentum boundary layer thickness. T and C denote the temperature and concentration of the fluid flow domain.

The temperature of the surface is held uniform at T_w which is higher than the ambient temperature T_{∞} . The species concentration at the surface is maintained uniform at C_w , which is taken to be C_{∞} and that considered to be C_{∞} away from the wall.

Non-similar transformations

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Far from the leading edge, i.e. in the free convection regime, the boundary layer is formed due to buoyancy forces. This suggests the following transformations:

$$\left. \begin{aligned} \Psi &= \nu Gr_x^{1/4} F(\xi, \eta) \quad \eta = \frac{y}{x} Gr_x^{1/4}, \quad \xi = x \\ \theta &= \frac{T - T_\infty}{T_w - T_\infty}, \quad \phi = \frac{C - C_\infty}{C_w - C_\infty}, \quad Gr_x = \frac{g \beta_c \Delta T x^3}{\nu^2} \\ &Gr_{Cx} = \frac{g \beta_c \Delta C x^3}{\nu^2} \end{aligned} \right\} \rightarrow (6)$$

As we know that

$$u = \frac{\partial \Psi}{\partial y}, \quad v = - \frac{\partial \Psi}{\partial x}$$

$$\text{So, } u = \frac{\partial \Psi}{\partial y} = \frac{\partial}{\partial y} \left(\nu Gr_x^{1/4} F(\xi, \eta) \right) = \nu Gr_x^{1/4} \frac{\partial}{\partial y} F(\xi, \eta)$$

$$u = \nu Gr_x^{1/4} \left(\frac{\partial F}{\partial \xi} \cdot \frac{\partial \xi}{\partial y} + \frac{\partial F}{\partial \eta} \cdot \frac{\partial \eta}{\partial y} \right)$$

$$u = \nu Gr_x^{1/4} F' \left(\frac{Gr_x^{1/4}}{x} \right)$$

$$u = \frac{\nu Gr_x^{1/2}}{x} F'$$

$$\frac{\partial \Psi}{\partial x} = \nu \frac{\partial}{\partial x} \left[\frac{1}{x} Gr_x^{1/2} F' \right]$$

$$\begin{aligned} \eta &= \frac{y}{x} Gr_x^{1/4} \\ \frac{\partial \eta}{\partial x} &= y \frac{\partial}{\partial x} \left(x^{-1} Gr_x^{1/4} \right) \\ &= y \left[-\frac{1}{x^2} Gr_x^{1/4} + \frac{1}{4} \frac{Gr_x^{1/4}}{x} \cdot \frac{\partial (Gr_x)}{\partial x} \right] \\ &= y \left[-\frac{Gr_x^{1/4}}{x^2} + \frac{Gr_x^{1/4}}{4 Gr_x^{3/4}} \cdot \frac{\partial (g \beta_c \Delta T x^3)}{\partial x} \right] \\ &= y \left[-\frac{Gr_x^{1/4}}{x^2} + \frac{Gr_x^{1/4} \cdot x \cdot g \beta_c \Delta T \cdot 3x^2}{4 Gr_x^{3/4} \cdot x} \right] \\ &= y \left[-\frac{Gr_x^{1/4}}{x^2} + \frac{x \cdot Gr_x^{1/4} \cdot g \beta_c \Delta T \cdot 3x^2}{4 Gr_x^{3/4} \cdot x} \right] \\ &= y \left[-\frac{Gr_x^{1/4}}{x^2} + \frac{3 Gr_x^{1/4}}{4 x^2 Gr_x} \cdot \frac{\partial (g \beta_c \Delta T x^3)}{\partial x} \right] \end{aligned}$$

$$\frac{\partial u}{\partial x} = v \left[-\frac{1}{x^2} G_{rx} \frac{1}{2} F' + \frac{1}{x} \cdot \frac{1}{2} \frac{G_{rx} \frac{1}{2}}{G_{rx}} \cdot \frac{g \rho \Delta T}{2x^2} F' + \frac{1}{x} G_{rx} \frac{1}{2} \left(\frac{\partial F'}{\partial \beta} \cdot \frac{\partial \beta}{\partial x} + \frac{\partial F'}{\partial \eta} \cdot \frac{\partial \eta}{\partial x} \right) \right] \quad (5)$$

$$\frac{\partial u}{\partial x} = v \left[-\frac{G_{rx} \frac{1}{2}}{x^2} F' + \frac{3 G_{rx} \frac{1}{2}}{2x^2} \cdot \frac{1}{G_{rx}} \cdot \frac{g \rho \Delta T}{2x^2} F' + \frac{1}{x} G_{rx} \frac{1}{2} \left(\frac{\partial F'}{\partial \beta} + \frac{\eta F''}{2x} \right) \right]$$

$$\frac{\partial u}{\partial x} = v \left[-\frac{G_{rx} \frac{1}{2}}{x^2} F' + \frac{3 G_{rx} \frac{1}{2}}{2x^2} F' + \frac{G_{rx} \frac{1}{2}}{x} \left(\frac{\partial F'}{\partial \beta} - \frac{\eta F''}{2x} \right) \right]$$

$$\frac{\partial u}{\partial x} = \frac{v}{x^2} G_{rx} \frac{1}{2} \left[-F' + \frac{3}{2} F' + x \frac{\partial F'}{\partial \beta} - \frac{\eta F''}{2} \right]$$

$$\frac{\partial u}{\partial x} = \frac{v}{x^2} G_{rx} \frac{1}{2} \left[\frac{1}{2} F' + \frac{3}{2} \frac{\partial F'}{\partial \beta} + \frac{\eta F''}{2} \right]$$

Now $u \frac{\partial u}{\partial x} = \frac{v \cdot G_{rx} \frac{1}{2}}{x} \cdot \frac{v}{x^2} G_{rx} \frac{1}{2} \left(\frac{1}{2} F' + \frac{3}{2} \frac{\partial F'}{\partial \beta} + \frac{\eta F''}{2} \right)$

$$u \frac{\partial u}{\partial x} = \frac{v^2 \cdot G_{rx}}{x^3} \left(\frac{1}{2} F'^2 + \frac{3}{2} \frac{\partial F'}{\partial \beta} F' + \frac{\eta F' F''}{2} \right)$$

$$u \frac{\partial u}{\partial x} = \frac{v^2 G_{rx}}{x^3} \left(\frac{1}{2} F'^2 + \frac{3}{2} \frac{\partial F'}{\partial \beta} F' + \frac{\eta F' F''}{2} \right) \quad (7)$$

$$y = \frac{1}{x} \left[-\frac{G_{rx} \frac{1}{4}}{x^2} + \frac{3}{4x^2} \frac{G_{rx} \frac{1}{4}}{G_{rx}} \cdot G_{rx} \right]$$

$$\frac{\partial \eta}{\partial x} = \frac{1}{x} \left[-\frac{G_{rx} \frac{1}{4}}{x^2} + \frac{3}{4x^2} \frac{G_{rx} \frac{1}{4}}{G_{rx}} \right]$$

$$\frac{\partial \eta}{\partial x} = -\frac{1}{x} \frac{G_{rx} \frac{1}{4}}{x^2} + \frac{3}{4x^2} \frac{G_{rx} \frac{1}{4}}{G_{rx}}$$

$$\frac{\partial \eta}{\partial x} = \left[-\frac{1}{x} \cdot \frac{1}{x} \frac{G_{rx} \frac{1}{4}}{x} + \frac{3}{4x} \cdot \frac{1}{x} \frac{G_{rx} \frac{1}{4}}{x} \right]$$

$$\frac{\partial \eta}{\partial x} = \left[-\frac{\eta}{x} + \frac{3\eta}{4x} \right]$$

$$\frac{\partial \eta}{\partial x} = \frac{-4\eta + 3\eta}{4x}$$

$$\frac{\partial \eta}{\partial x} = +\frac{\eta}{2x} = \frac{\eta}{2x}$$

$$\frac{\partial \eta}{\partial x} = \frac{G_{rx} \frac{1}{4}}{x}$$

$$\frac{\partial u}{\partial y} = \frac{2G\mu x^{\frac{1}{2}}}{x} \left(\frac{\partial F'}{\partial \xi} \cdot \frac{\partial \xi}{\partial y} + \frac{\partial F'}{\partial \eta} \cdot \frac{\partial \eta}{\partial y} \right) \quad (6)$$

$$\frac{\partial u}{\partial y} = \frac{2G\mu x^{\frac{1}{2}}}{x} F'' \cdot \left(\frac{G\mu x^{\frac{1}{4}}}{x} \right) = \frac{2G\mu x^{\frac{3}{4}}}{x^2} F''$$

$$\boxed{\frac{\partial u}{\partial y} = \frac{2G\mu x^{\frac{3}{4}}}{x^2} F''}$$

$$\frac{\partial^2 u}{\partial y^2} = \frac{2 \cdot G\mu x^{\frac{3}{4}}}{x^2} F''' \cdot \frac{G\mu x^{\frac{1}{4}}}{x} = \frac{2 \cdot G\mu x}{x^3} F'''$$

$$\boxed{\frac{\partial^2 u}{\partial y^2} = \frac{2G\mu x}{x^3} F'''} \rightarrow (8)$$

$$v = -\frac{\partial \psi}{\partial x} = -\frac{\partial}{\partial x} \left(2G\mu x^{\frac{1}{4}} F(\xi, \eta) \right)$$

$$v = -2 \frac{\partial}{\partial x} \left(G\mu x^{\frac{1}{4}} F(\xi, \eta) \right)$$

$$v = -2 \left[G\mu x^{\frac{1}{4}} \left(\frac{\partial F}{\partial \xi} \cdot \frac{\partial \xi}{\partial x} + \frac{\partial F}{\partial \eta} \cdot \frac{\partial \eta}{\partial x} \right) + \frac{1}{4} \frac{G\mu x^{\frac{1}{4}}}{G\mu x} \cdot \frac{\partial (2G\mu x^{\frac{1}{4}})}{\partial x} \cdot F \right]$$

$$v = -2 \left[G\mu x^{\frac{1}{4}} \left(\frac{\partial F}{\partial \xi} + \frac{\eta}{2x} F' \right) + \frac{3G\mu x^{\frac{1}{4}}}{4x} F \right]$$

$$\boxed{v = -2 \left[G\mu x^{\frac{1}{4}} \left(\frac{\partial F}{\partial \xi} + \frac{\eta F'}{2x} \right) + \frac{3G\mu x^{\frac{1}{4}}}{4x} F \right]}$$

(7)

$$v \frac{\partial u}{\partial y} = -v \left[\frac{G_{rx} \frac{1}{4}}{x} \left(x \frac{\partial F}{\partial z} + \frac{\eta F'}{2} \right) + \frac{3}{4x} G_{rx} \frac{1}{4} F \right] \times \frac{v G_{rx} \frac{3}{4}}{x^2} F''$$

$$v \frac{\partial u}{\partial y} = -\frac{v^2 G_{rx}}{x^3} \left[3 \frac{\partial F}{\partial z} F'' + \frac{\eta F' F''}{2} + \frac{3}{4} F F'' \right]$$

$$\boxed{v \frac{\partial u}{\partial y} = -\frac{v^2 G_{rx}}{x^3} \left(3 F'' \frac{\partial F}{\partial z} + \frac{\eta F' F''}{2} + \frac{3}{4} F F'' \right)} \rightarrow (9)$$

$$\theta = \frac{T - T_{\infty}}{T_w - T_{\infty}} \Rightarrow \theta (T_w - T_{\infty}) = T - T_{\infty}$$

$$\Rightarrow \boxed{\theta \Delta T + T_{\infty} = T} = \theta \Delta T = T - T_{\infty} \rightarrow (10)$$

$$\phi = \frac{C - C_{\infty}}{C_w - C_{\infty}} \Rightarrow \phi (C_w - C_{\infty}) = C - C_{\infty}$$

$$\Rightarrow \boxed{\phi \Delta C = C - C_{\infty}} \rightarrow (11)$$

Put 7-11 in equation

Now $\frac{\partial u}{\partial y} =$
 put equation 7-11 in equation (2), we have

$$\frac{v^2 G_{rx}}{x^3} \left[\frac{1}{2} F''^2 + 3 F' \frac{\partial F'}{\partial z} + \frac{\eta}{2} F' F'' - 3 F'' \frac{\partial F}{\partial z} - \frac{\eta F' F''}{2} - \frac{3}{4} F F'' \right]$$

$$= v \cdot \frac{v \cdot G_{rx}}{x^3} F''' + g \beta_c \Delta T \theta + g \beta_c \Delta C \phi$$

$$u \frac{\partial T}{\partial x} = v \left(\frac{\partial T}{\partial x} + \frac{\partial T}{\partial y} \right) + \frac{\partial T}{\partial z}$$

$$u \frac{\partial T}{\partial x} = v \left(\frac{\partial T}{\partial x} + \frac{\partial T}{\partial y} \right) + \frac{\partial T}{\partial z}$$

$$\frac{\partial T}{\partial x} = \Delta T \left(\frac{\partial z}{\partial x} + \frac{\partial y}{\partial x} \right)$$

$$\frac{\partial T}{\partial x} = \Delta T \left(\frac{\partial z}{\partial x} + \frac{\partial y}{\partial x} \right)$$

$$\therefore T \text{ is constant}$$

$$\frac{\partial z}{\partial x} = 1$$

$$\therefore z = x$$

$$\frac{\partial T}{\partial x} = \Delta T \left(\frac{\partial z}{\partial x} + \frac{\partial y}{\partial x} \right) + \frac{\partial T}{\partial z}$$

Now we move towards to energy equation.

This is non-similar transformed momentum equation.

$$F_{III} + \frac{1}{3} F_{II} - \frac{1}{2} F_{I^2} + \phi = \frac{1}{3} \left(F_{I^2} \frac{\partial F}{\partial z} - F_{II} \frac{\partial F}{\partial z} \right)$$

$$\frac{1}{2} F_{I^2} + \frac{1}{3} F_{II} - \frac{1}{3} F_{II} \frac{\partial F}{\partial z} - \frac{1}{3} F_{II} \frac{\partial F}{\partial z} + \frac{1}{2} F_{I^2} - \frac{1}{3} F_{II} \frac{\partial F}{\partial z} + \phi = \frac{1}{3} \left(F_{I^2} \frac{\partial F}{\partial z} - F_{II} \frac{\partial F}{\partial z} \right)$$

(8)

$$\left(\frac{\partial T}{\partial x} = \frac{2 G r x^{\frac{1}{2}} \Delta T}{x^2} \left(3 F' \frac{\partial \theta}{\partial \xi} + \frac{\eta \theta'}{\xi} \right) \right) \rightarrow (13)$$

$$\frac{\partial T}{\partial y} = \Delta T \left(\frac{\partial \theta}{\partial \xi} \cdot \frac{\partial \xi}{\partial y} + \frac{\partial \theta}{\partial \eta} \cdot \frac{\partial \eta}{\partial y} \right) + \frac{\partial T}{\partial y}$$

$$\frac{\partial T}{\partial y} = \Delta T \theta' \cdot \frac{G r x^{\frac{1}{4}}}{x}$$

$$\frac{\partial T}{\partial y} = \frac{\Delta T G r x^{\frac{1}{4}}}{x} \theta'$$

$$\frac{\partial^2 T}{\partial y^2} = \frac{\Delta T G r x^{\frac{1}{4}}}{x} \theta'' \cdot \frac{G r x^{\frac{1}{4}}}{x}$$

$$\frac{\partial^2 T}{\partial y^2} = \frac{\Delta T G r x^{\frac{1}{2}}}{x^2} \theta'' \rightarrow (14)$$

$$\nu \frac{\partial T}{\partial y} = - \nu \left[G r x^{\frac{1}{4}} \left(\frac{\partial F}{\partial \xi} + \frac{\eta F'}{2x} \right) + \frac{3}{4} G r x^{\frac{1}{4}} F \right] \times \frac{\Delta T G r x^{\frac{1}{4}}}{x} \theta'$$

$$\nu \frac{\partial T}{\partial y} = - \frac{\nu}{x^2} G r x^{\frac{1}{2}} \left[\theta' x \frac{\partial F}{\partial \xi} + \frac{\eta F' \theta'}{2} + \frac{3}{4} F \theta' \right]$$

$$\nu \frac{\partial T}{\partial y} = - \frac{\nu G r x^{\frac{1}{2}}}{x^2} \left[3 \theta' \frac{\partial F}{\partial \xi} + \frac{\eta F' \theta'}{2} + \frac{3}{4} F \theta' \right] \rightarrow (15)$$

Put eq. (13) - (15) in Eq. (3), we obtain

$$\frac{2 \Gamma_{\infty} \frac{1}{2} \Delta T}{\alpha^2} \left(3 F' \frac{\partial \theta}{\partial \xi} + \frac{\eta \theta' F'}{2} \right) - \frac{2 \Gamma_{\infty} \frac{1}{2}}{\alpha^2} \left(3 \theta' \frac{\partial F}{\partial \xi} + \frac{\eta F' \theta'}{2} + \frac{3}{4} F \theta' \right)$$

$$= \frac{\Delta T \Gamma_{\infty} \frac{1}{2}}{\alpha^2} \theta'' + \frac{\alpha^2}{2 \Gamma_{\infty} \frac{1}{2} \Delta T}$$

$$3 \left(F' \frac{\partial \theta}{\partial \xi} - \theta' \frac{\partial F}{\partial \xi} \right) + \frac{\eta \theta' F'}{2} - \frac{\eta F' \theta'}{2} - \frac{3}{4} F \theta' = \frac{\alpha}{2} \theta''$$

$$3 \left(F' \frac{\partial \theta}{\partial \xi} - \theta' \frac{\partial F}{\partial \xi} \right) - \frac{3}{4} F \theta' = \frac{\alpha}{2} \theta''$$

$$3 \left(F' \frac{\partial \theta}{\partial \xi} - \theta' \frac{\partial F}{\partial \xi} \right) = \frac{1}{Pr} \theta'' + \frac{3}{4} F \theta'$$

$$\frac{1}{Pr} = \frac{\alpha}{2}$$

$$\therefore Pr = \frac{2}{\alpha}$$

$$\boxed{\frac{1}{Pr} \theta'' + \frac{3}{4} F \theta' = 3 \left(F' \frac{\partial \theta}{\partial \xi} - \theta' \frac{\partial F}{\partial \xi} \right)} \rightarrow (16)$$

This is non-similar transformed form of energy equation.

Now, we transform mass concentration equation.

$$\phi = \frac{C - C_{\infty}}{C_w - C_{\infty}}$$

$$\Delta C = C_w - C_{\infty}$$

$$\phi \Delta C = C - C_{\infty}$$

$$\phi \Delta C + C_{\infty} = C$$

$$\frac{\partial \phi}{\partial \eta} = \Delta C \left(\frac{\partial \phi}{\partial \xi} + \frac{\eta}{\xi} \frac{\partial \phi}{\partial \eta} + \frac{\eta}{\xi} \frac{\partial \phi}{\partial \eta} \right) + \frac{\partial \phi}{\partial \eta}$$

$$\frac{\partial \phi}{\partial x} = \Delta C \left(\frac{\partial \phi}{\partial z} + \frac{\eta}{2x} \phi' \right) \quad (11)$$

$$\frac{\partial \phi}{\partial x} \cong \Delta C \left(\frac{\partial \phi}{\partial z} + \frac{\eta \phi'}{2x} \right)$$

$$u \frac{\partial \phi}{\partial x} = \frac{v G_{m,2}^{1/2}}{x} F' \Delta C \left(\frac{\partial \phi}{\partial z} + \frac{\eta \phi'}{2x} \right)$$

$$= \frac{v \Delta C}{x^2} G_{m,2}^{1/2} \left(F' \frac{\partial \phi}{\partial z} + \eta \phi' F' \right)$$

$$u \frac{\partial \phi}{\partial x} = \frac{v \Delta C}{x^2} G_{m,2}^{1/2} \left(F' \frac{\partial \phi}{\partial z} + \eta \phi' F' \right) \rightarrow (17) \quad \text{--- } z = x$$

$$\frac{\partial \phi}{\partial y} = \Delta C \left(\frac{\partial \phi}{\partial z} \cdot \frac{\partial z}{\partial y} + \frac{\partial \phi}{\partial x} \cdot \frac{\eta}{2x} \right) + \frac{\partial \phi}{\partial y}$$

$$\frac{\partial \phi}{\partial y} = \Delta C \left(\frac{\partial \phi}{\partial z} + \frac{\eta}{2x} \phi' \right)$$

$$\frac{\partial \phi}{\partial y} \cong \Delta C \phi' \cdot \frac{G_{m,2}^{1/4}}{x}$$

$$\frac{\partial \phi}{\partial y} = \frac{\Delta C G_{m,2}^{1/4}}{x} \phi'$$

$$\frac{\partial^2 \phi}{\partial y^2} = \frac{\Delta C \Gamma_{rn}^{1/4}}{\alpha} \phi' \quad \text{--- (17)}$$

$$\frac{\partial^2 \phi}{\partial y^2} = \frac{\Delta C \Gamma_{rn}^{1/2}}{\alpha^2} \phi'' \quad \rightarrow \text{(18)}$$

$$\nu \frac{\partial \phi}{\partial y} = - \frac{2 \Gamma_{rn}^{1/2}}{\alpha^2} \left[\frac{3}{2} \phi' \frac{\partial F}{\partial \xi} + \frac{\eta F \phi'}{2} + \frac{3}{4} F \phi' \right] \quad \rightarrow \text{(19)}$$

Put eqs (17) & (18) into eq. (4), we get

$$\frac{2 \Gamma_{rn}^{1/2} \Delta C}{\alpha^2} \left(\frac{3}{2} F \phi' \frac{\partial \phi}{\partial \xi} + \frac{\eta F \phi'}{2} \right) - \frac{2 \Gamma_{rn}^{1/2}}{\alpha^2} \left(\frac{3}{2} \phi' \frac{\partial F}{\partial \xi} + \frac{\eta F \phi'}{2} + \frac{3}{4} F \phi' \right)$$

$$= D_m \frac{\Delta C \Gamma_{rn}^{1/2}}{\alpha^2} \phi''$$

$$\frac{3}{2} \left(F \phi' \frac{\partial \phi}{\partial \xi} - \phi' \frac{\partial F}{\partial \xi} \right) + \frac{\eta F \phi'}{2} - \frac{\eta F \phi'}{2} - \frac{3}{4} F \phi' = \frac{D_m}{\nu} \frac{\Delta C \Gamma_{rn}^{1/2}}{\alpha^2} \frac{\alpha^2}{\Gamma_{rn}^{1/2} \Delta C} \phi''$$

$$\frac{3}{2} \left(F \phi' \frac{\partial \phi}{\partial \xi} - \phi' \frac{\partial F}{\partial \xi} \right) - \frac{3}{4} F \phi' = \frac{1}{Sc} \phi''$$

$$\frac{1}{Sc} \phi'' + \frac{3}{4} F \phi' = \frac{3}{2} \left(F \phi' \frac{\partial \phi}{\partial \xi} - \phi' \frac{\partial F}{\partial \xi} \right)$$

where Pr , Sc are Prandtl number; $Sc = \frac{\nu}{D_m}$
and Schmidt number, respectively.

B-c)

(B)

Transformed boundary conditions are

$$F(\xi, 0) = F'(\xi, 0) = 0, \quad \theta(\xi, 0) = 1, \quad \phi(\xi, 0) = 1, \quad \text{at } \eta = 0$$
$$F(\xi, \infty) = 0, \quad \theta(\xi, \infty) = 0, \quad \phi(\xi, \infty) = 0, \quad \text{as } \eta \rightarrow \infty.$$

Here to note that:

F is dimensionless Stream function.

x coordinate measuring distance along the plate.

y coordinate measuring distance normal to plate.

α effective thermal diffusivity (m^2/sec).

θ is dimensionless temperature.

ϕ is dimensionless mass concentration.

η is similarity variable

ξ scaled streamwise coordinate

